

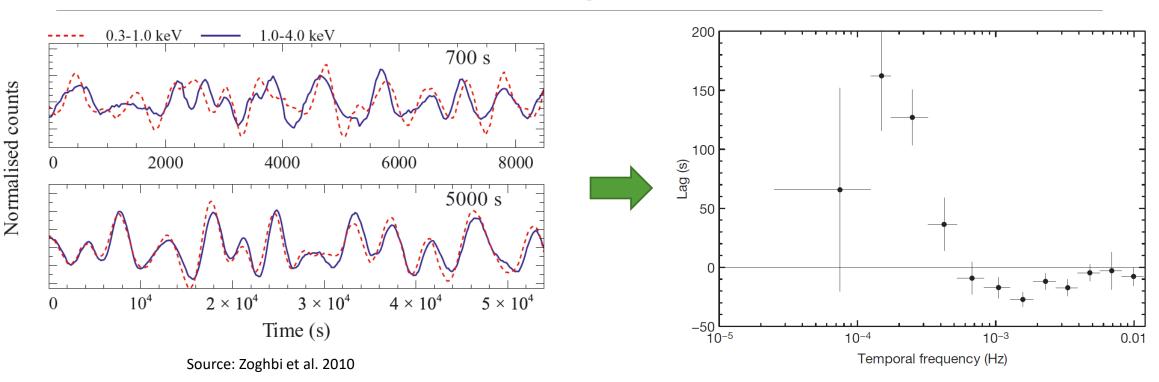
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A Theoretical Fourier-Transformation Model for the Formation of X-ray Time Lags from Black Hole Accretion Disks

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Source: Fabian et al. 2009





Definition of Time Lags

Definition of Time Lags

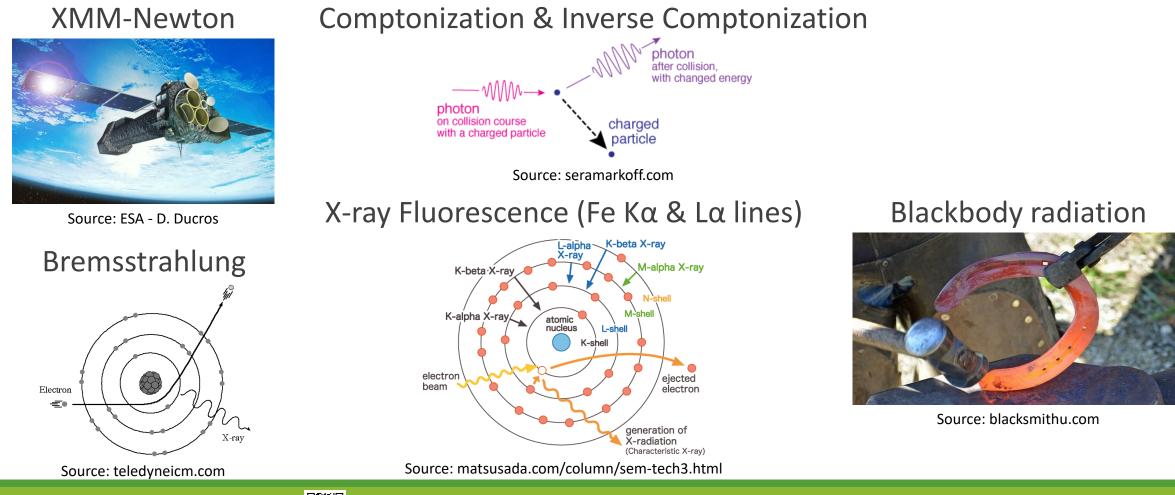
$$\begin{split} & \widehat{\mathfrak{s}}(t) = \mathscr{F}(\epsilon_{\mathrm{s}}, t) \\ & \mathfrak{h}(t) = \mathscr{F}(\epsilon_{\mathrm{h}}, t) \\ & \widehat{\mathfrak{F}}(\omega) = \widetilde{\mathscr{F}}(\epsilon_{\mathrm{s}}, t) = \int_{-\infty}^{+\infty} \mathfrak{s}(t) e^{i\omega t} dt \\ & \mathfrak{H}(\omega) = \widetilde{\mathscr{F}}(\epsilon_{\mathrm{h}}, t) = \int_{-\infty}^{+\infty} \mathfrak{h}(t) e^{i\omega t} dt \\ & \underbrace{\mathfrak{C}(\omega) = \mathfrak{S}^{*}(\omega)\mathfrak{H}(\omega)}_{\phi(\omega) = \operatorname{Arg}[\mathfrak{S}^{*}(\omega)\mathfrak{H}(\omega)]} \longrightarrow \delta t(\omega) = \frac{\phi(\omega)}{\omega} \end{split}$$

$$\begin{split} \mathfrak{h}(t + \Delta t) &= \mathfrak{s}(t) \\ \mathfrak{H}(\omega) &= \int_{-\infty}^{+\infty} \mathfrak{h}(t) e^{i\omega t} dt \\ \mathfrak{S}(\omega) &= \int_{-\infty}^{+\infty} \mathfrak{h}(t + \Delta t) e^{i\omega t} dt \\ \mathfrak{S}(\omega) &= e^{-i\omega\Delta t} \int_{-\infty}^{+\infty} \mathfrak{h}(t + \Delta t) e^{i\omega(t + \Delta t)} d(t + \Delta t) \\ &= e^{-i\omega\Delta t} \mathfrak{H}(\omega) \\ &= e^{i\omega\Delta t} |\mathfrak{H}(\omega)|^2 \\ \phi(\omega) &= \operatorname{Arg}[e^{i\omega\Delta t} |\mathfrak{H}(\omega)|^2] = \omega\Delta t \\ \delta t(\omega) &= \frac{\omega\Delta t}{\omega} = \Delta t \end{split}$$





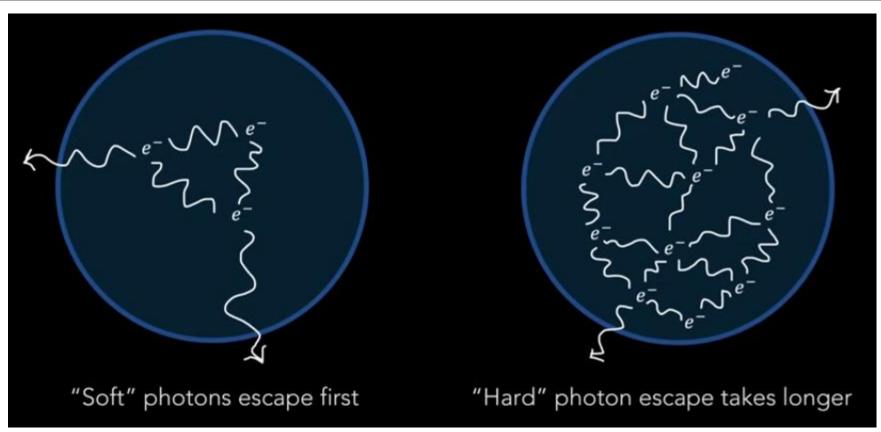








Previous Model

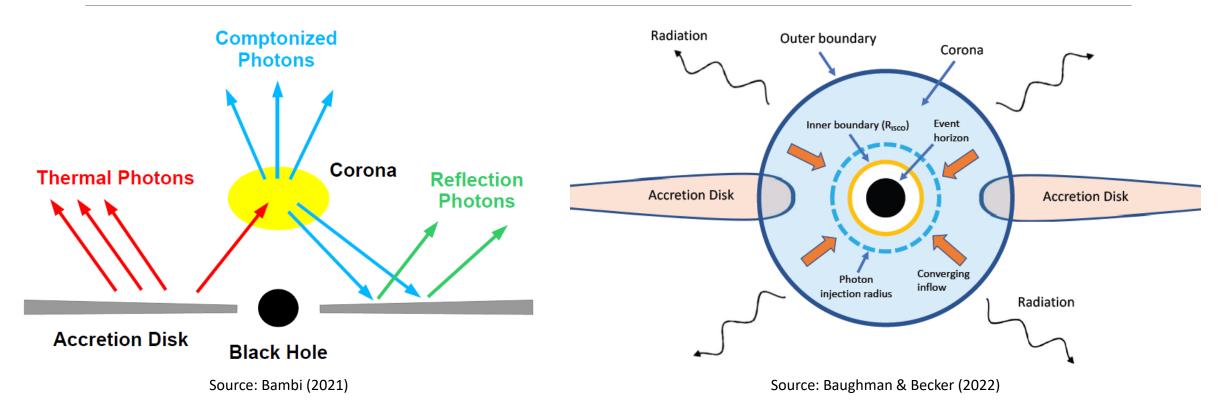


Source: Baughman & Becker (2022)





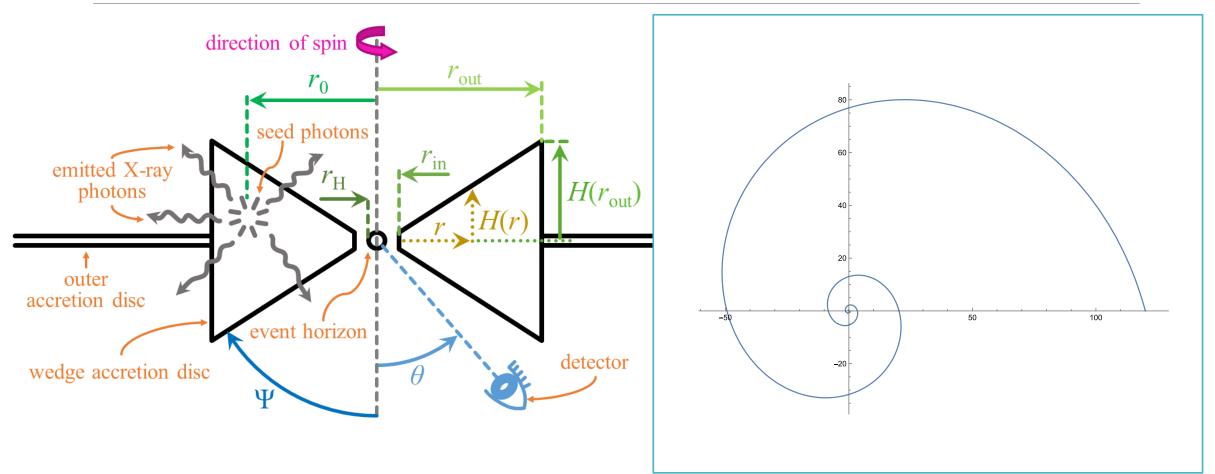
















Model Derivation

Radiation distribution function $f(\vec{r}, \epsilon, t)$

Vertical average

- + introduce Green's function
- + Fourier transform
- + introduce dimensionless parameters

Fourier-Transform of Green's function $F_{\rm G}(\tau, \chi, q)$

 $-i\tilde{\omega}\frac{\hat{u}\hat{b}\tilde{r}^{3/2}}{\dot{m}}F_{\rm G} - \hat{u}\tilde{r}^{-1/2}\frac{\partial F_{\rm G}}{\partial \tau} = -\frac{\hat{u}^{2}\hat{b}}{2\dot{m}}\chi\frac{\partial F_{\rm G}}{\partial \chi} + \frac{2\hat{u}\hat{b}\tilde{r}^{1/2}}{3\dot{m}}\frac{\partial F_{\rm G}}{\partial \tau} + \frac{1}{3}\frac{\partial^{2}F_{\rm G}}{\partial \tau^{2}} - \frac{\hat{u}\tilde{r}^{1/2}}{\dot{m}}\operatorname{Min}\left(\frac{\hat{u}\tilde{r}^{1/2}}{\dot{m}},1\right)F_{\rm G} + \frac{\Theta}{\chi^{2}}\frac{\partial}{\partial\chi}\left[\chi^{4}\left(F_{\rm G} + \frac{\partial F_{\rm G}}{\partial\chi}\right)\right] + \frac{N_{0}\delta(\tau - \tau_{0})\delta(\chi - \chi_{0})e^{i\tilde{\omega}q_{0}}}{4\pi cR_{g}^{2}\tilde{r}_{0}^{2}\hat{b}\Theta^{3}(m_{e}c^{2})^{3}\chi_{0}^{2}}$

Apply separation of variables

$$F_{\rm G}(\tau, \tau_0, \chi, \chi_0, \tilde{\omega}) = \sum_{n=0}^{\infty} C_n(\tau_0, \chi_0) \mathcal{G}_n(\tau, \tilde{\omega}) \mathcal{H}_n(\chi, \tilde{\omega})$$

Spatial-dependent equation

$$\frac{d^2\mathcal{G}}{d\tau^2} + \left[2\frac{\hat{u}\hat{b}}{\dot{m}}\tilde{r}^{1/2} + 3\hat{u}\tilde{r}^{-1/2}\right]\frac{d\mathcal{G}}{d\tau} + \left[i\frac{\hat{u}\hat{b}\tilde{\omega}}{\dot{m}}\tilde{r}^{3/2} - \frac{\hat{u}\tilde{r}^{1/2}}{\dot{m}}\right]$$
$$\operatorname{Min}\left(\frac{\hat{u}\tilde{r}^{1/2}}{\dot{m}}, 1\right) + \frac{\hat{u}^2\hat{b}}{2\dot{m}}\lambda\right]3\mathcal{G} = 0$$

Energy-dependent equation $\frac{1}{\chi^2} \frac{d}{d\chi} \left[\chi^4 \left(\mathcal{H} + \frac{d\mathcal{H}}{d\chi} \right) \right] - \frac{\hat{u}^2 \hat{b}}{2\dot{m}\Theta} \chi \frac{d\mathcal{H}}{d\chi} - \frac{\hat{u}^2 \hat{b}}{2\dot{m}\Theta} \lambda \mathcal{H} = 0$

Solve for $C_n(\tau_0, \chi_0)$, $\mathcal{G}_n(\tau, \tilde{\omega})$, and $\mathcal{H}_n(\chi, \tilde{\omega})$ to obtain the Green's function $F_{\mathbf{G}}$





Model Derivation

Fourier transform of X-ray light curve is given by

•
$$\tilde{\mathcal{F}}_{side}(\chi, \tilde{\omega}) = R_g \hat{b} \left(\frac{\chi k_{\rm B} T_e R_g \tilde{r}_{out}}{D} \right)^2 F_{\rm G}(\chi, \tau_{out}, \tilde{\omega})$$

• $\tilde{\mathcal{F}}_{face}(\chi, \tilde{\omega}) = \frac{\chi^2 \Theta^2 (m_e c^2)^2 \hat{b} R_g^3}{D^2} \int_0^{\tau_{out}} \frac{\hat{u} \tilde{r}^{5/2}}{\dot{m}} \operatorname{Min} \left(\frac{\hat{u} \tilde{r}^{1/2}}{\dot{m}}, 1 \right) F_{\rm G}(\chi, \tau, \tilde{\omega}) d\tau$

$$\tilde{\mathcal{F}}(\chi,\tilde{\omega}) = \tilde{\mathcal{F}}_{\rm side}(\chi,\tilde{\omega})\sin\theta + \tilde{\mathcal{F}}_{\rm face}(\chi,\tilde{\omega})\cos\theta = \tilde{\mathcal{F}}(\epsilon,\omega)$$

• Phase lag
$$\phi(\omega) = \operatorname{Arg}[\mathcal{S}^*(\omega)\mathcal{H}(\omega)]$$

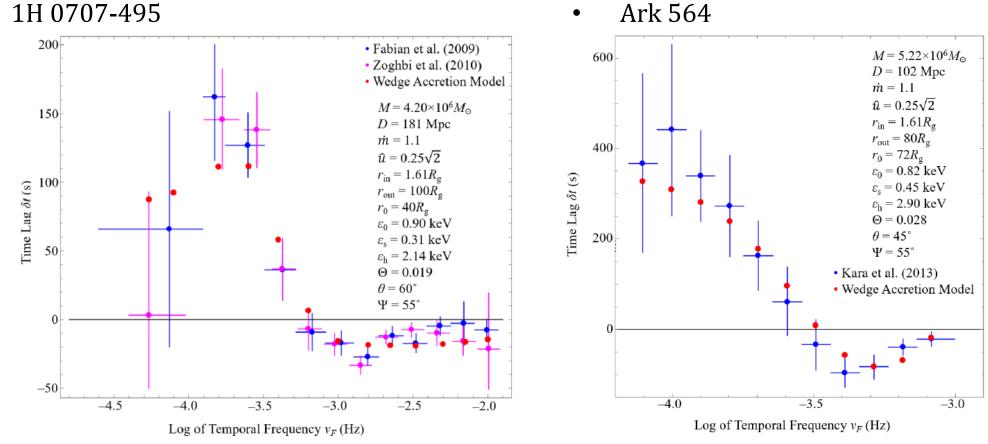
• Time lag $\delta t(\omega) = \frac{\phi(\omega)}{\omega}$





Applications

1H 0707-495 ٠



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Discussion

- Local thermal instability \rightarrow bremsstrahlung emission flash \rightarrow Fe L-line and K-line photons
- $k_{\rm B}T_e$ = 9.71 keV in 1H 0707-495 and $k_{\rm B}T_e$ = 14.31 keV in Ark 564
- $\epsilon \sim$ 0.7 keV for Fe L-line photons and $\epsilon \sim$ 6.4 keV for Fe K-line photons
- Fe L-line photons produce time lags in the observational window (\sim 0.3 4 keV)
- Fe K-line photons are outside the observational window but may illuminate the outer cooler disk, creating the reverberation feature noted by Kara et al. (2013)





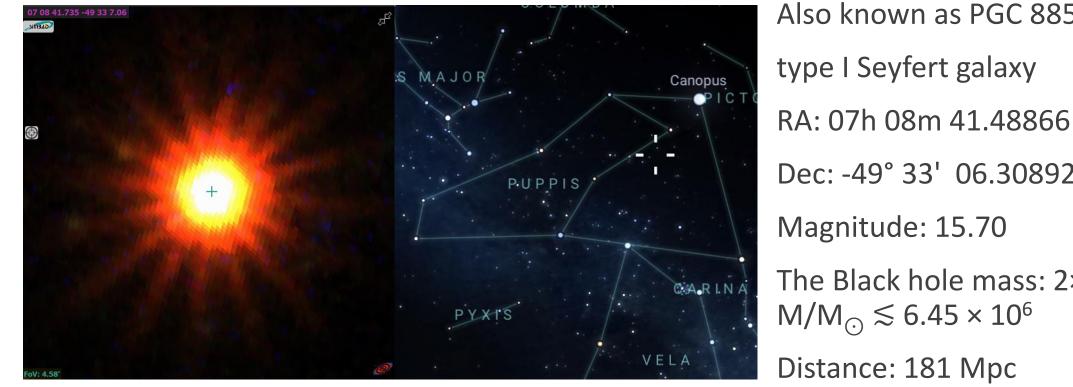
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Questions?

THANK YOU

Application to 1H 0707-495



Source: simbad.u-strasbg.fr XMM-Newton Catalog

Source: stellarium-web.org

Also known as PGC 88588 RA: 07h 08m 41.488661470s Dec: -49° 33' 06.308921700" The Black hole mass: $2 \times 10^6 \lesssim$





Application to Ark 564



Source: simbad.u-strasbg.fr XMM-Newton Catalog

Source: stellarium-web.org



