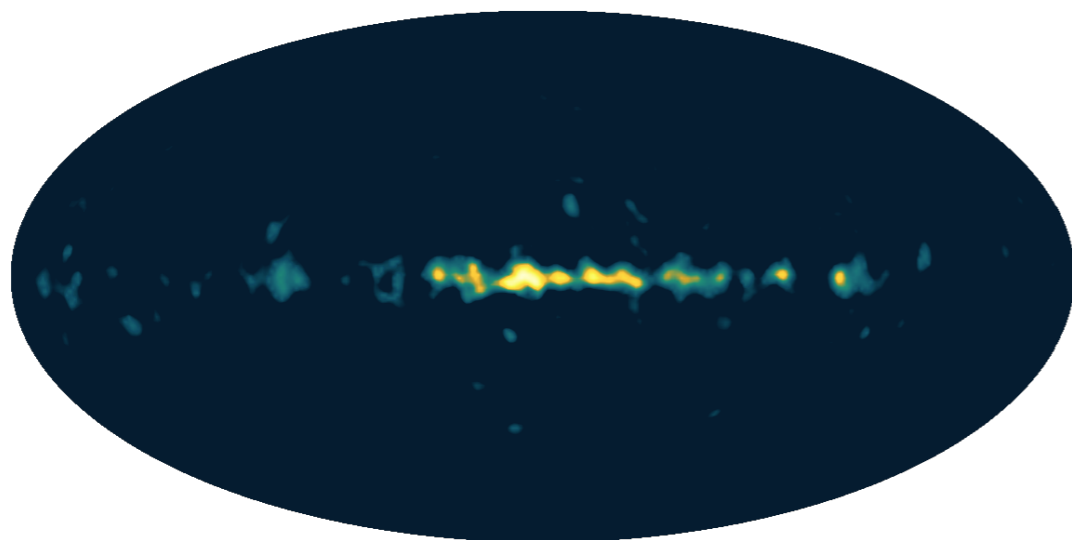


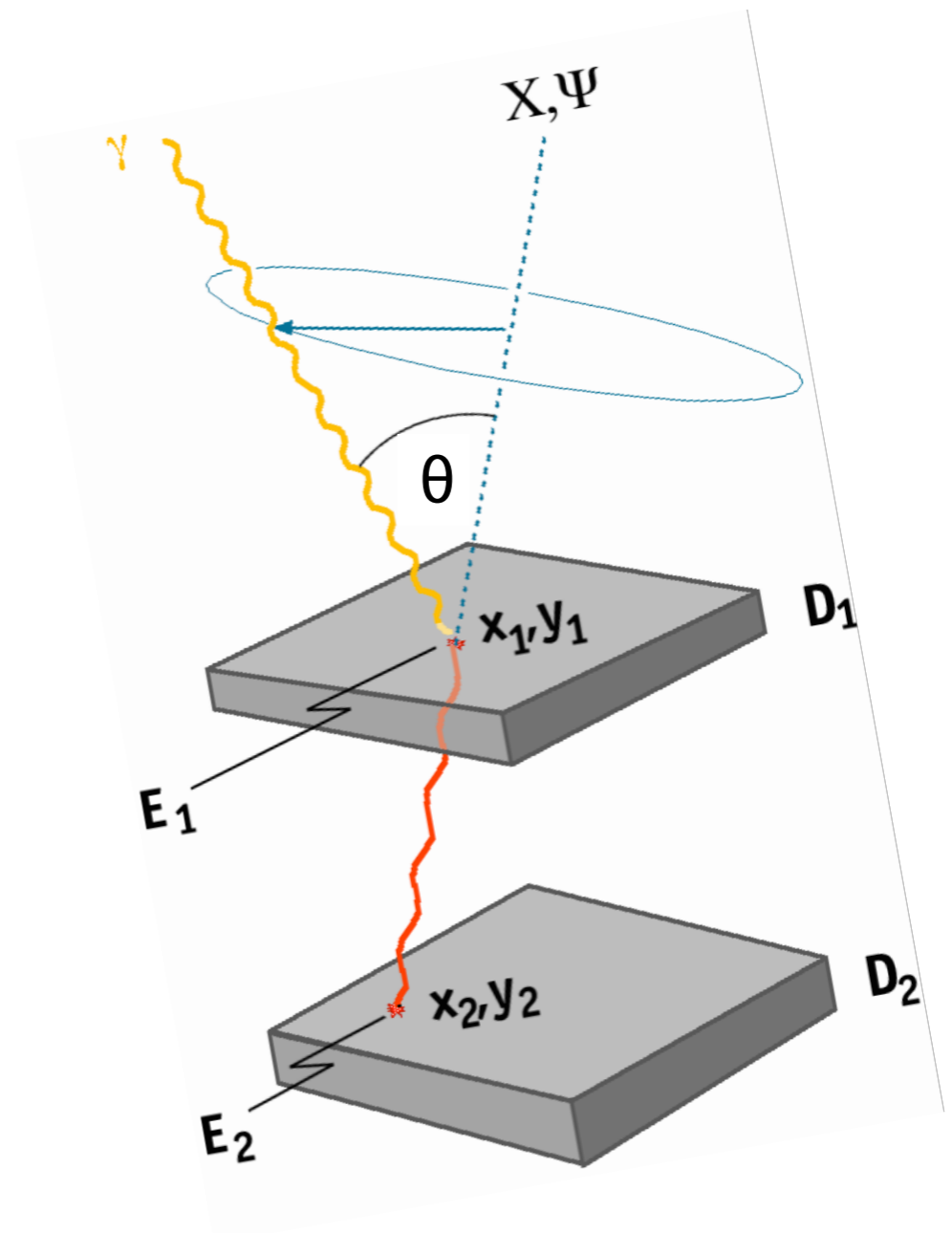
# Compton telescopes: Design considerations, etc

J. Eric Grove

U.S. Naval Research Laboratory

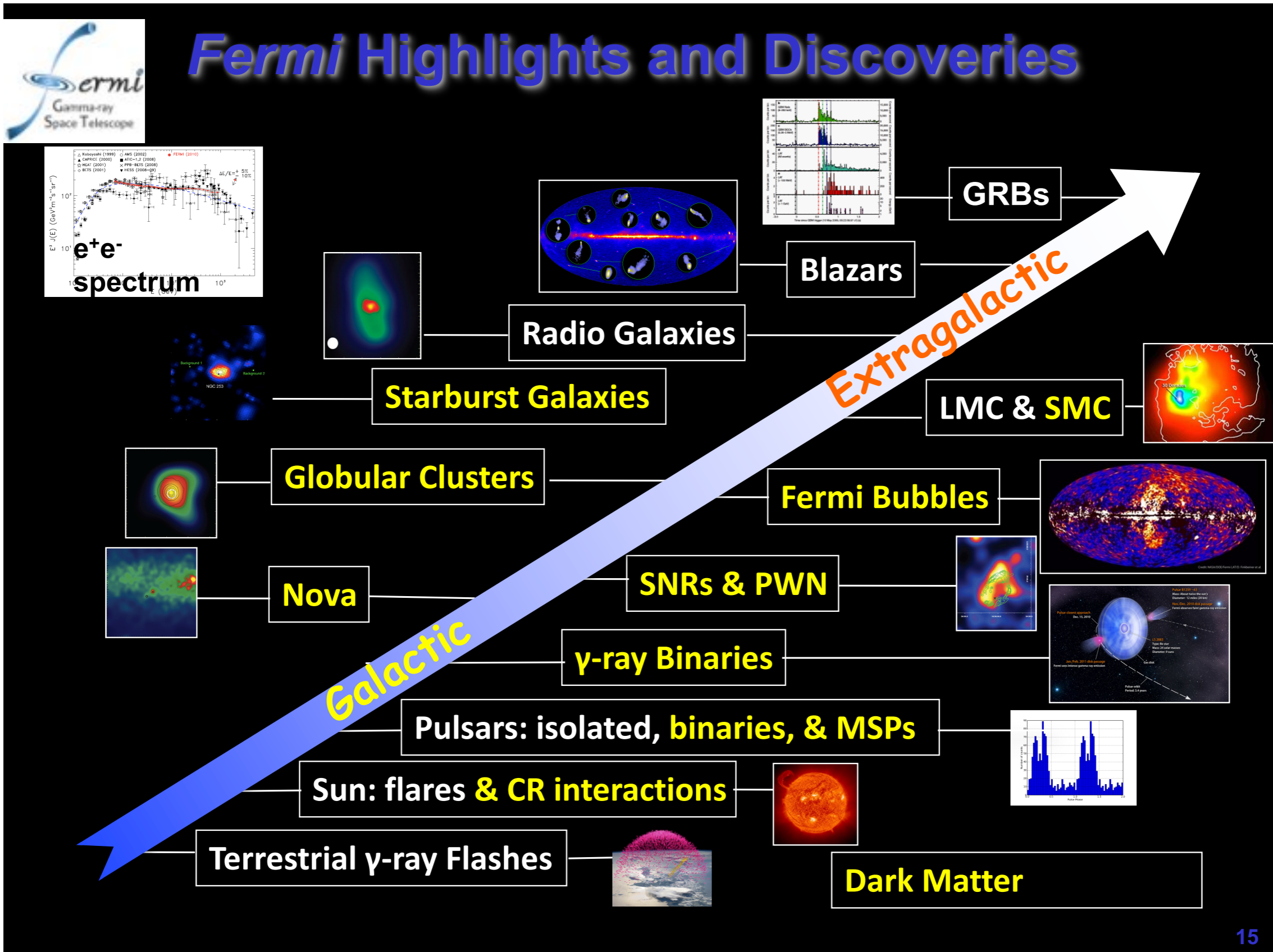


first  $^{26}\text{Al}$  all-sky map





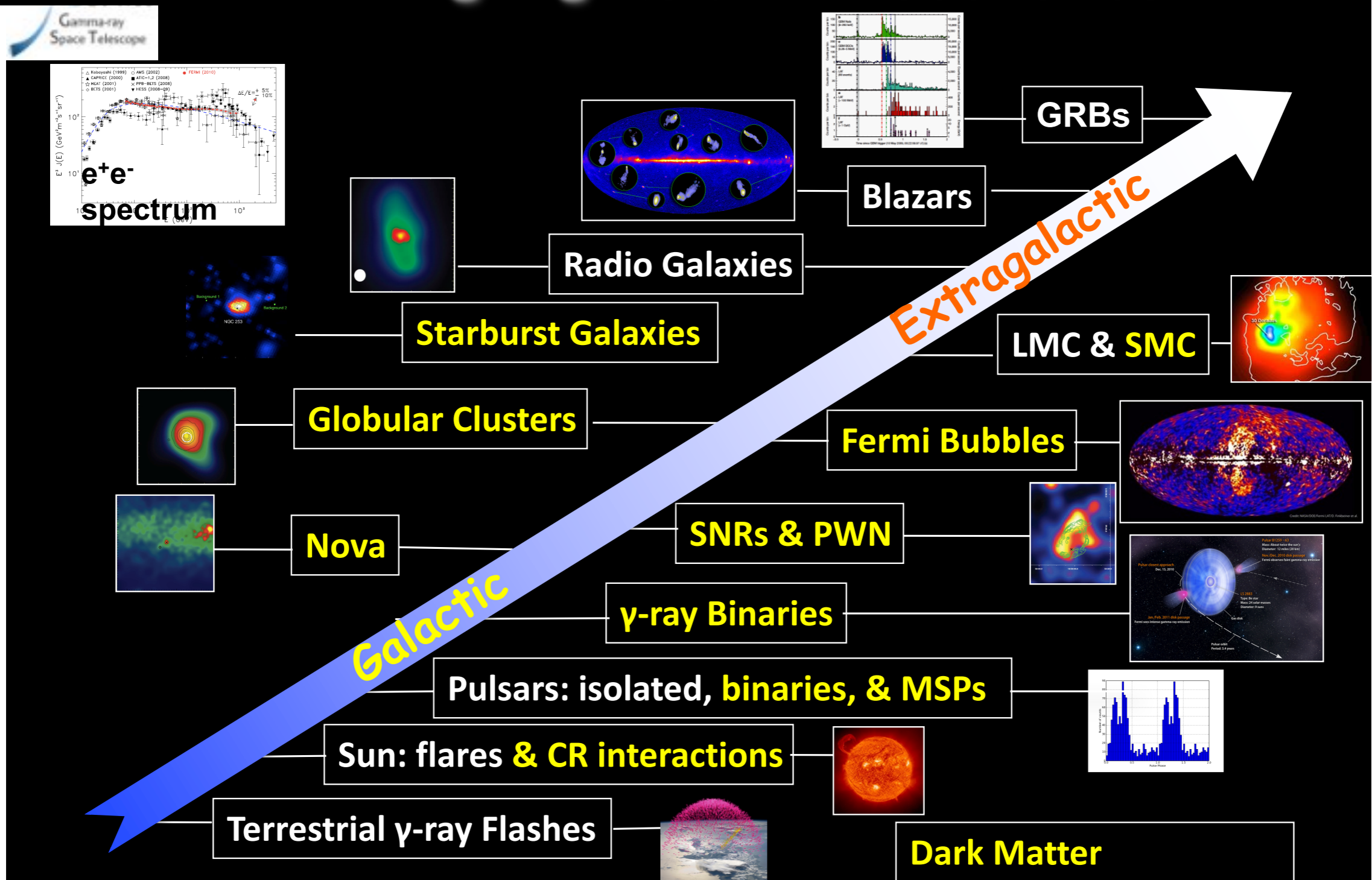
# Why observe near 1 MeV?





# Why observe near 1 MeV?

## Future MeV Telescope Highlights and Discoveries





# Why observe near 1 MeV?

## Future MeV Telescope Highlights and Discoveries

### *All this, and nuclear astrophysics too!*

*Studying the life cycle of matter, the synthesis and distribution of the elements*

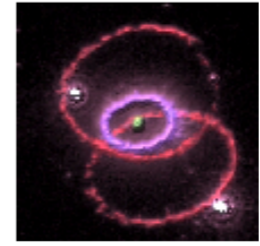
#### Giants ( $^{26}\text{Al}$ )

- Convective shell burning



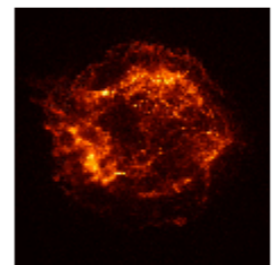
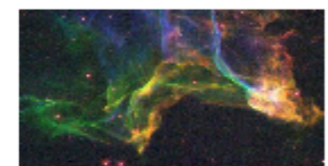
#### Supernovae ( $^{56}\text{Ni}$ , $^{57}\text{Ni}$ , $^{44}\text{Ti}$ )

- Understand the explosion mechanism
- Abundances of synthesized elements
- Structure & dynamics of ejecta



#### Stars/Sun (many)

- Ion acceleration
- Ambient abundances
- Accelerated abundances

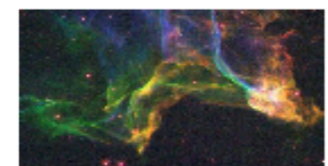


#### Remnants ( $^{44}\text{Ti}$ , $^{26}\text{Al}$ , $^{106}\text{Sn}$ )

- Structure & dynamics
- Ejection of heavy elements
- Compact object formation
- Discovery of young remnants
- CR acceleration

#### Interstellar Medium ( $^{26}\text{Al}$ , $^{60}\text{Fe}$ ,...)

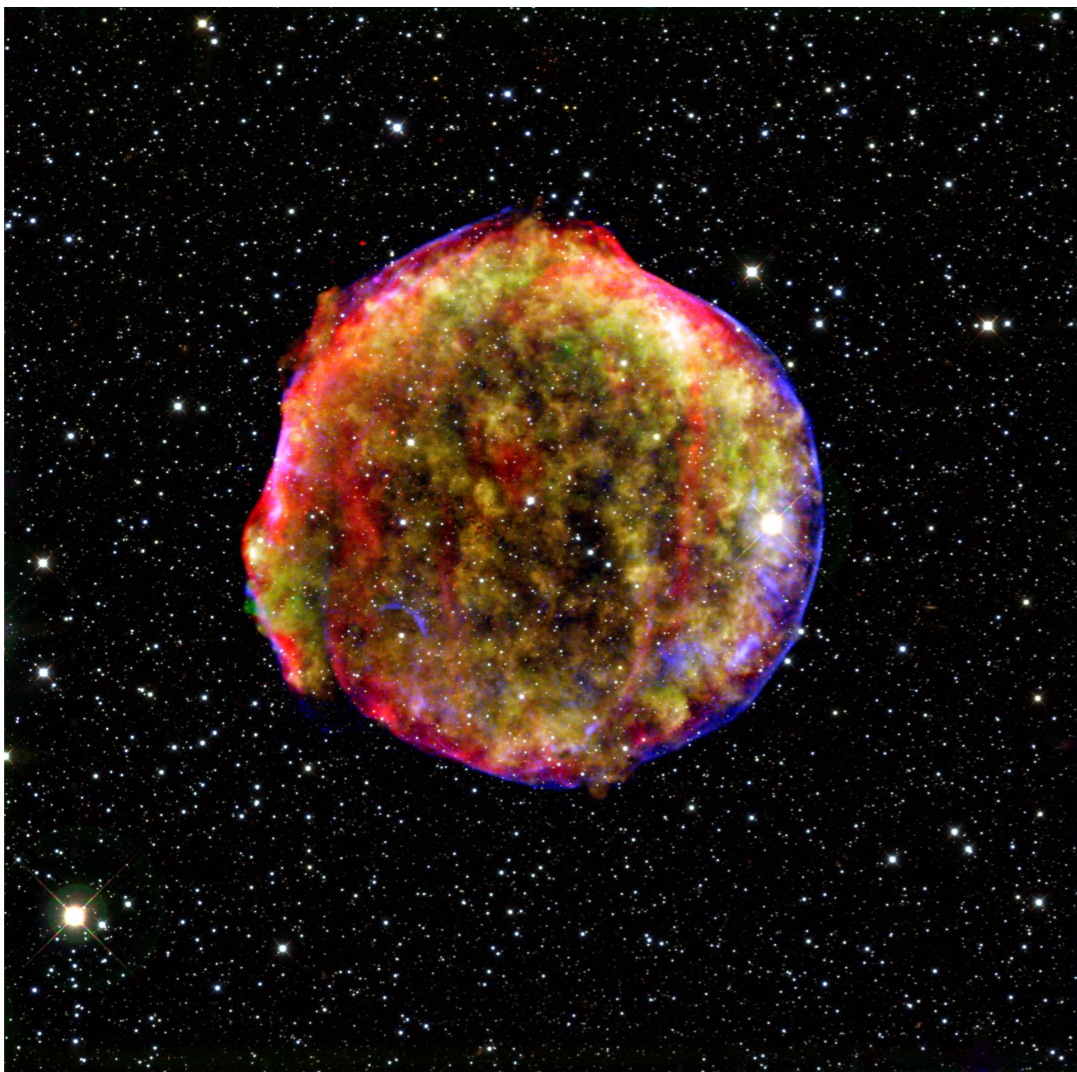
- Galactic history of nucleosynthesis
- Formation, evolution and death of massive stars
- CR interactions, chemical composition



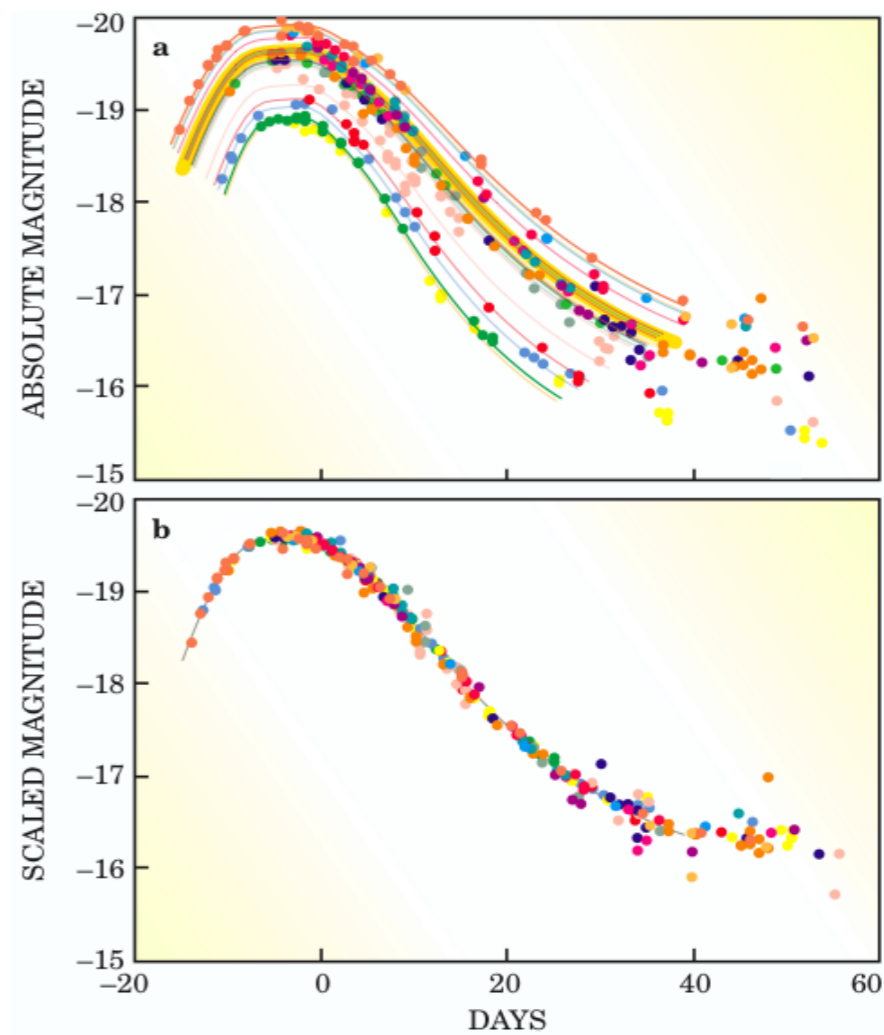


## Supernovae Type Ia and MeV astronomy

- SN Ia have been crucial in understanding the structure of the universe



Tycho SNR (1572 CE)  
Warren et al. 2005

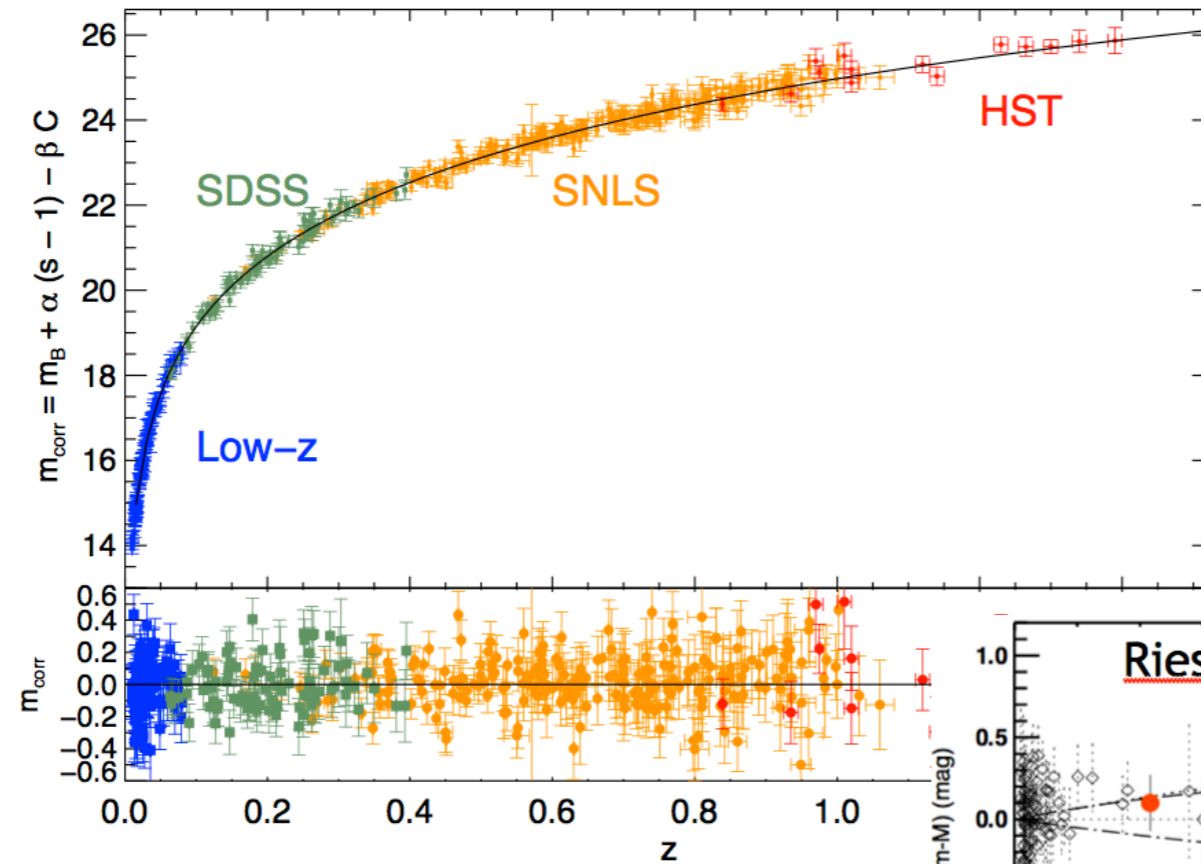


Perlmutter 1999 Physics Today

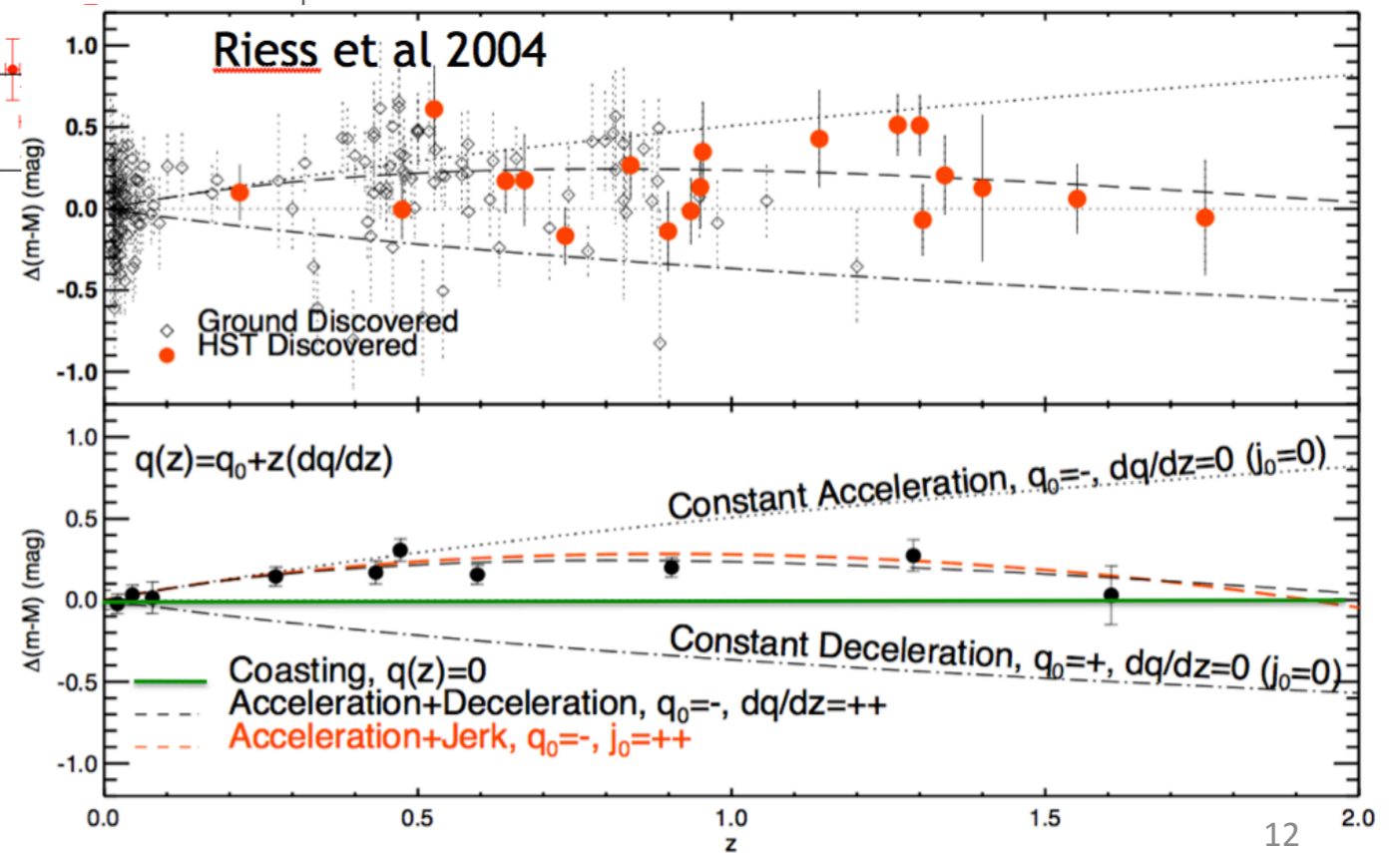


# Supernova Type Ia cosmology

Conley et al. 2011



Assuming a one-parameter calibration of light curves, we use SN Ia with dramatic implications for Dark Energy. We have to understand absolute luminosities to  $\sim 0.1$  mag to high  $z$  and low  $Z$ .



Grove - MeV Astronomy



## ... yet we do not understand the nearby SN Ia well

### Type Ia supernova scenario

- Nuclear ignition in a White Dwarf (WD), with a large fraction of the mass burned to nuclear statistical equilibrium, centered around  $^{56}\text{Ni}$ . The entire star is disrupted in the explosion.

### Explains:

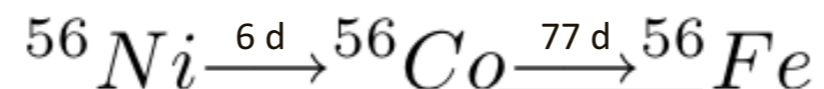
- Hydrogen deficiency
- Late-time Fe-group emission
- High velocities,  $\sim 10^4$  km/s
- Presence in elliptical galaxies

### Outstanding Issues:

- Nature of progenitor systems, composition & mass of WD
- Ignition location(s)/conditions
- Nuclear flame propagation
- Effect of rotation, **B** on explosion
- Nature of observed luminosity/duration correlations, and impact of metallicity

(Timmes et al. 2003, Jackson et al. 2010)

$^{56}\text{Ni}$  is key to understanding these explosions....



158 keV  
812 keV  
etc.

847 keV  
1238 keV  
etc.

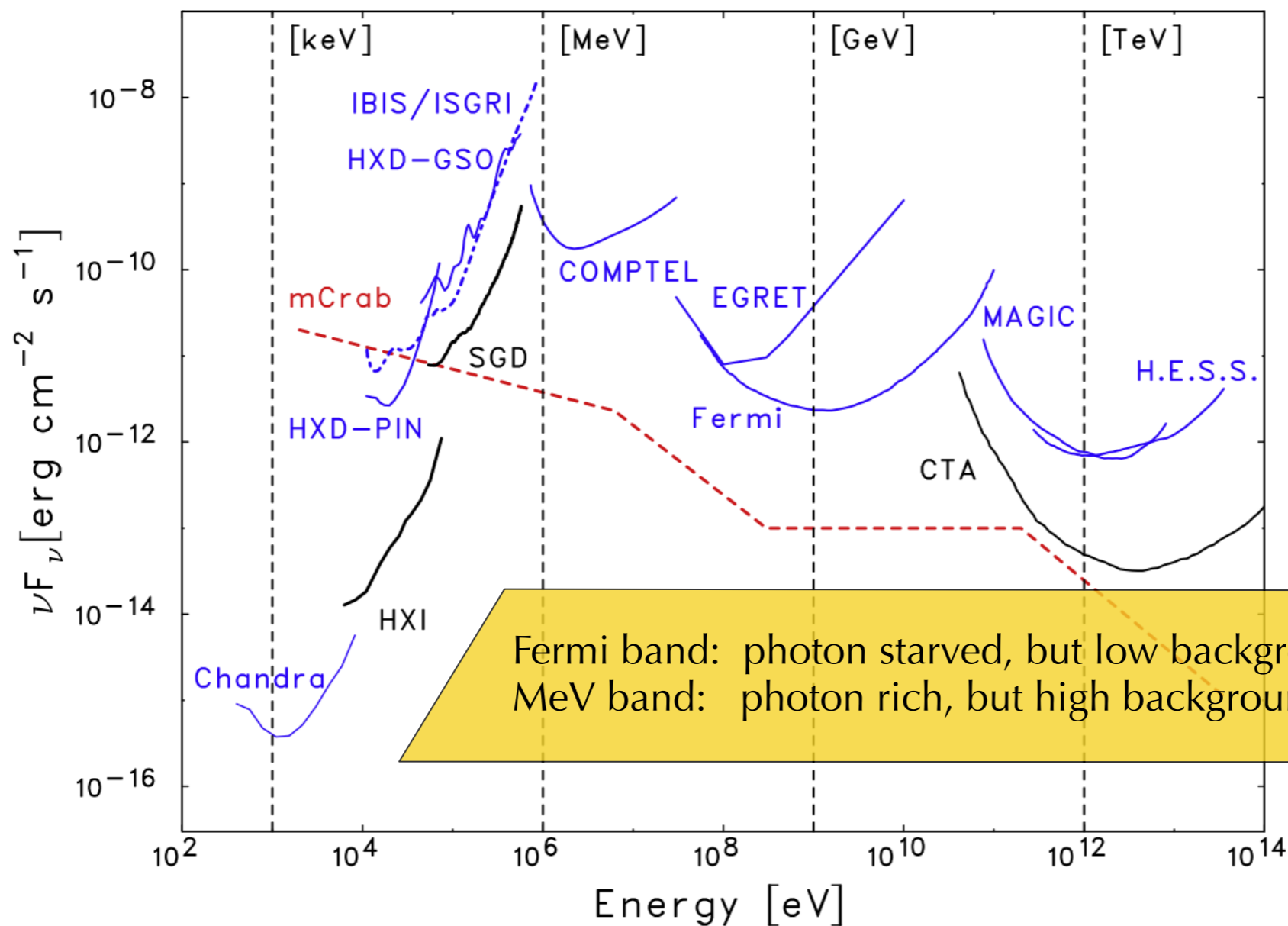
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13



# Why hasn't this been done already?

Rich science at ~1 MeV is hampered, so far, by poor sensitivity



Why so difficult?

- No focusing; imaging difficult
- Small interaction probability
- High background

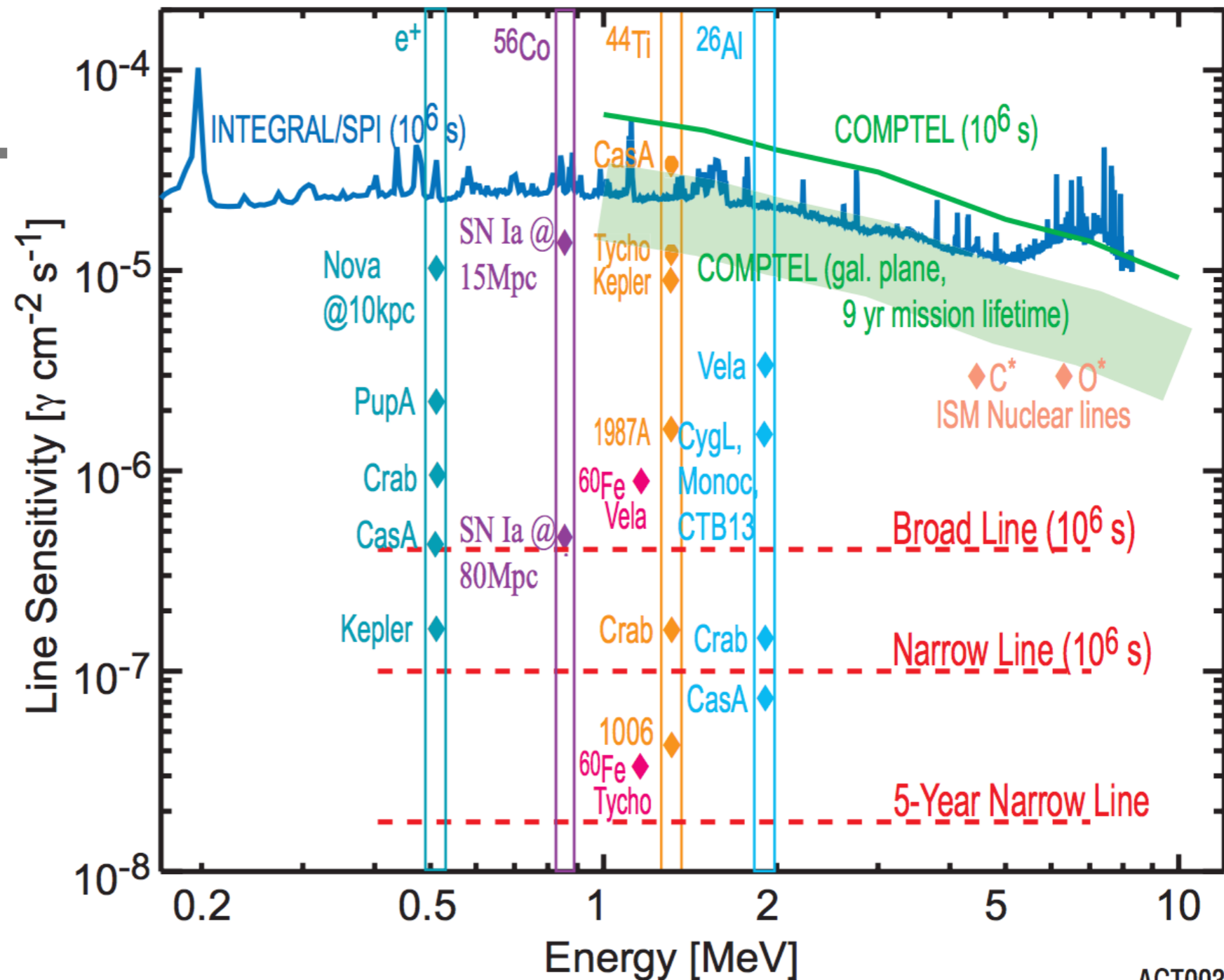
Differential continuum sensitivities of X-ray and  $\gamma$ -ray instruments for isolated point sources

Takahashi et al. 2013





From Advanced Compton Telescope mission concept study (Boggs 2007)



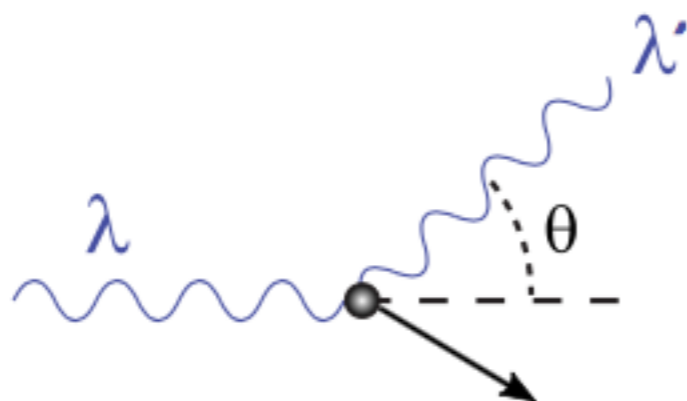
ACT003

- Narrow and broad line sensitivity needed for nuclear astrophysics
  - Compare current/historical mission sensitivities to theoretical predictions
    - COMPTTEL and Integral/SPI:  $\sim \text{few} \times 10^{-5} \text{ ph cm}^{-2} \text{ s}^{-1}$  in  $10^6$  sec
    - Want  $< \text{few} \times 10^{-7} \text{ ph cm}^{-2} \text{ s}^{-1}$  in  $10^6$  sec
  - Goal
    - Time dependence and many sources



# Interaction with matter

- Photon interaction cross section
  - Photoelectric absorption
  - Compton scattering
  - Pair production
- Compton scattering
  - Interaction with bound electrons



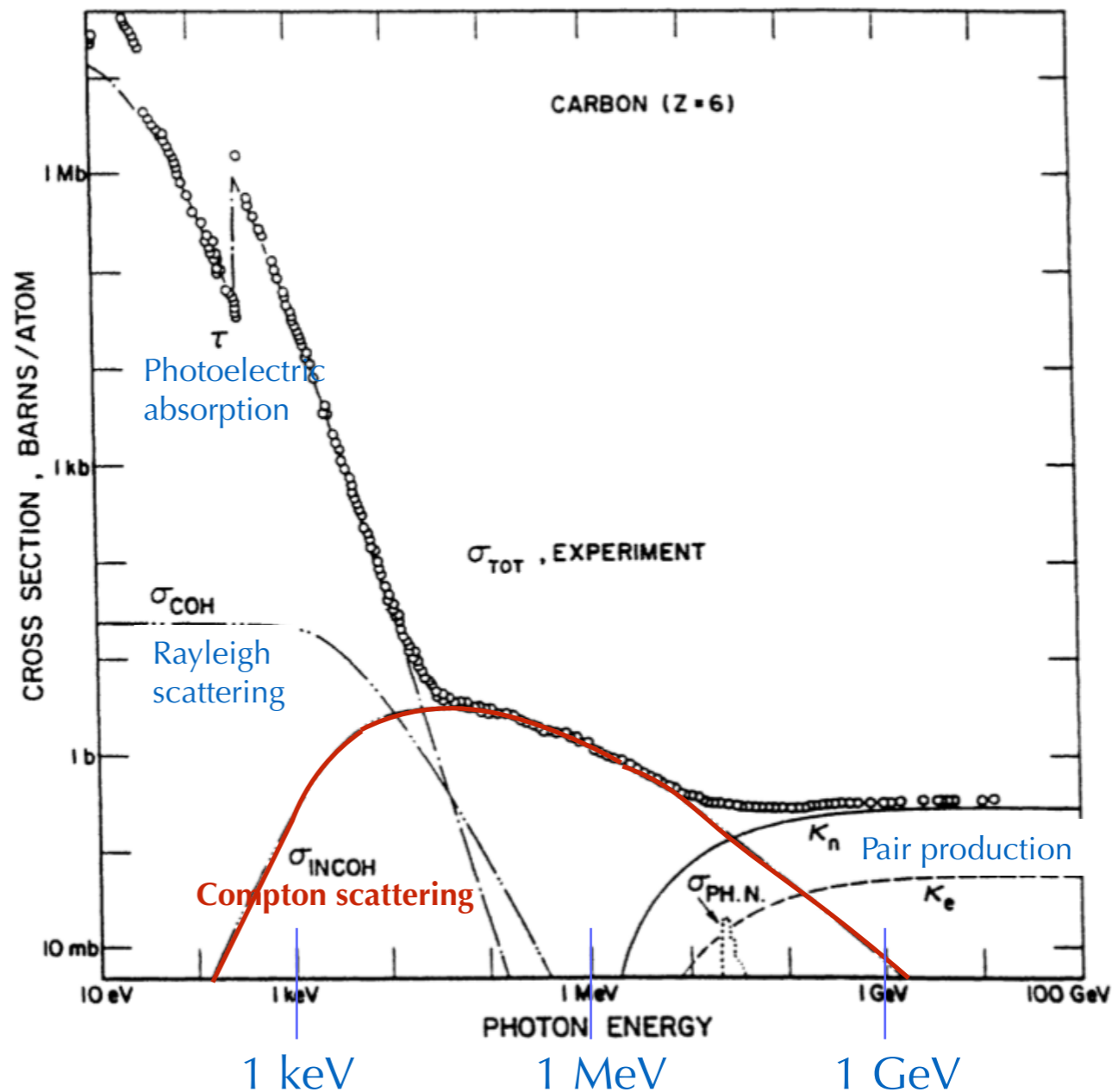
Credit: Wikipedia

$$\lambda' - \lambda = \frac{h}{m_e c} (1 - \cos \theta)$$

- Scattering angle and photon energy loss are directly related

Review of photon interaction cross section data

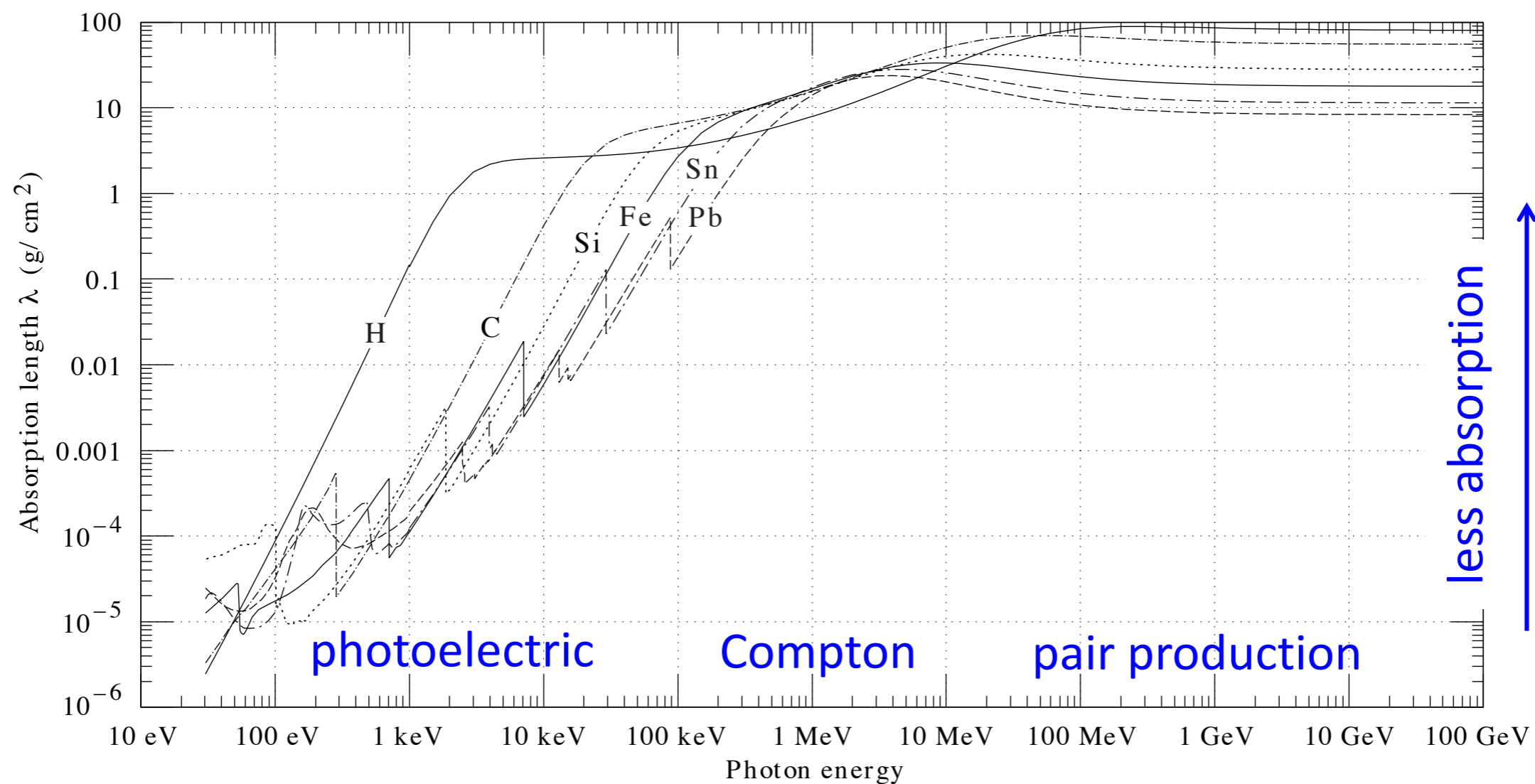
Hubbell 1999





# Interaction with matter

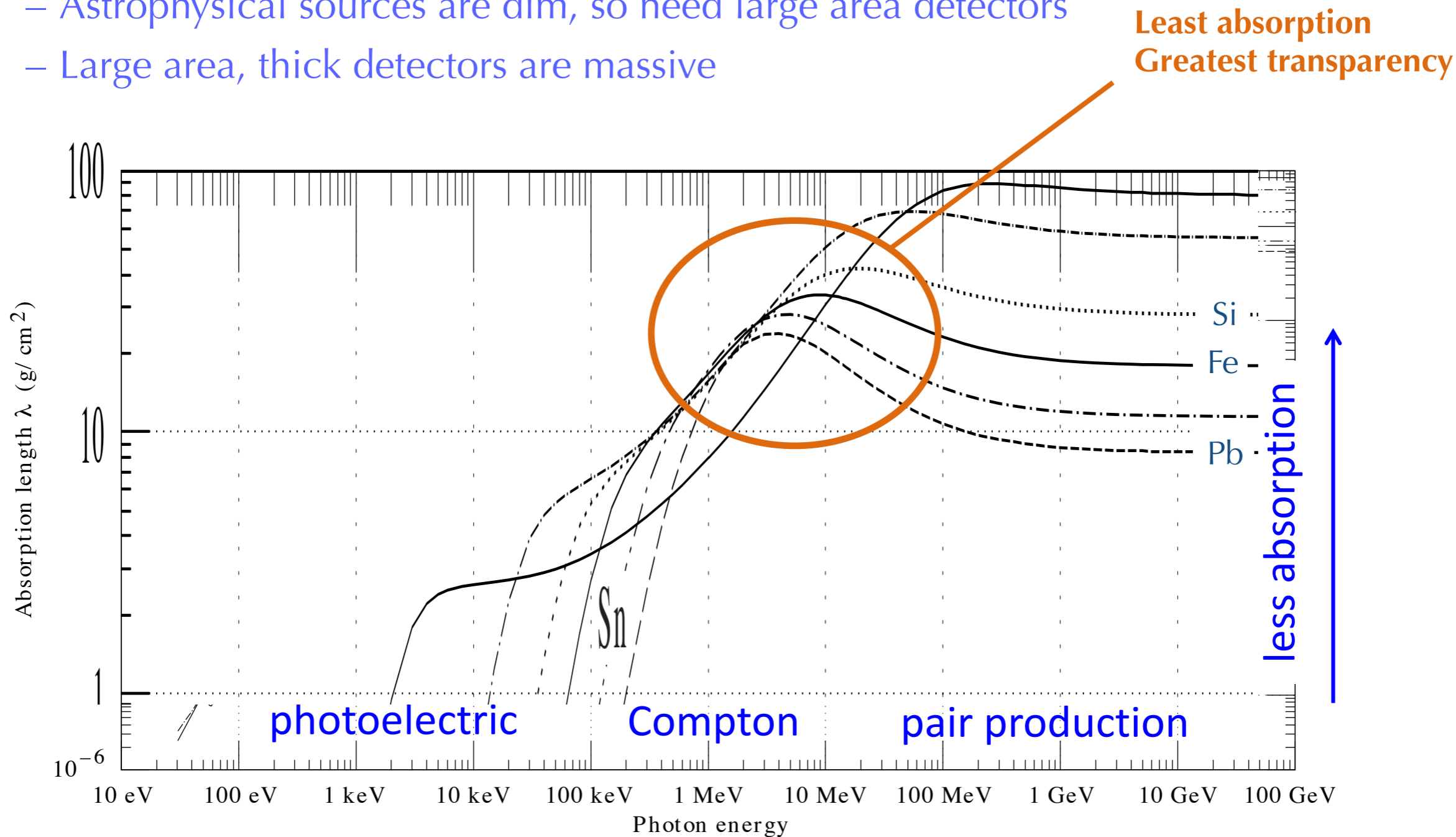
- Detector materials are most transparent in the MeV band!
  - Need thick detectors
  - Astrophysical sources are dim, so need large area detectors
  - Large area, thick detectors are massive





# Interaction with matter

- Detector materials are most transparent in the MeV band!
  - Need thick detectors
  - Astrophysical sources are dim, so need large area detectors
  - Large area, thick detectors are massive



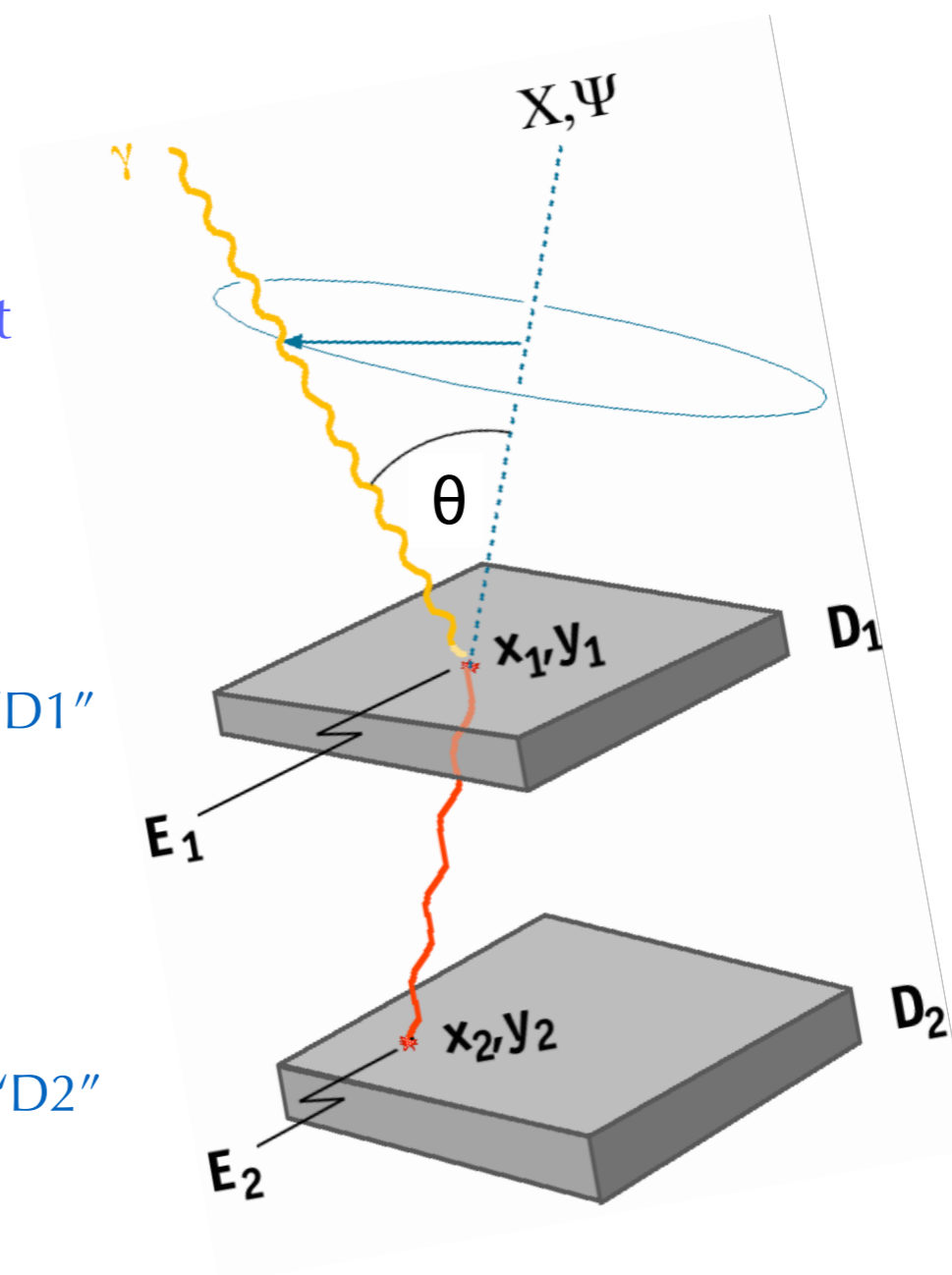


# Basic Compton telescope

- Two components
  - One detector to scatter the  $\gamma$  ray
  - Another detector to absorb the scattered  $\gamma$
- Use Compton formula to reconstruct incident photon direction
  - Need to measure position and energy deposition in both detectors

Scattering detector: "D1"

Absorbing detector: "D2"



- What happens if scattered  $\gamma$  ray isn't fully absorbed? How would we know it's fully absorbed?



# Compton telescope principles

- Compton event reconstruction

- Incident energy = sum of measured energies  $E_1, E_2$
- Incident direction from Compton formula

$$\cos\theta = 1 - m_e c^2 \left( \frac{1}{E_2} - \frac{1}{E_1 + E_2} \right)$$

- Incident direction reconstructs to cone
  - Annulus on sky
- Width of cone (“angular response measure”) is function of E and position resolution

- dE term:

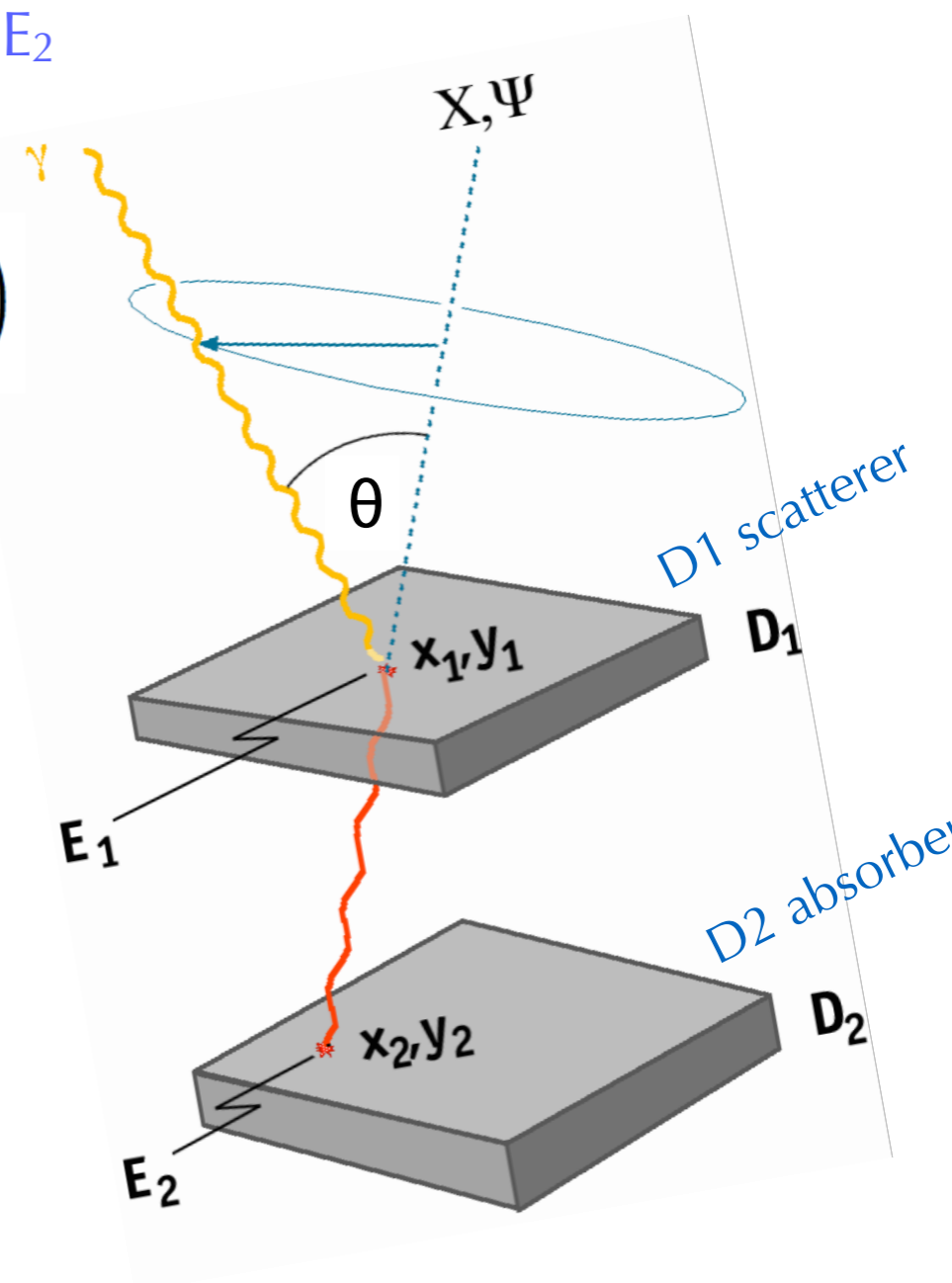
$$\Delta\theta = \frac{m_e c^2}{\sin\theta} \left[ \frac{\Delta E_1^2}{(E_1 + E_2)^4} + \Delta E_2^2 \left( \frac{1}{E_2^2} - \frac{1}{(E_1 + E_2)^2} \right)^2 \right]^{1/2}$$

- Must have good dE and good dx

- Angular resolution depends on spectral performance
  - e.g. good dE: Si, Ge, LaBr<sub>3</sub>, Srl<sub>2</sub>

- Must fully absorb in D2

- Reconstructed E and direction will be wrong if energy escapes
- (Not entirely trivial to find such a detector material; e.g. need lots of Si, Ge)





# Compton telescope principles

- Compton event reconstruction

- Width of cone (“angular response measure”) is function of E and position resolution

- dE term:

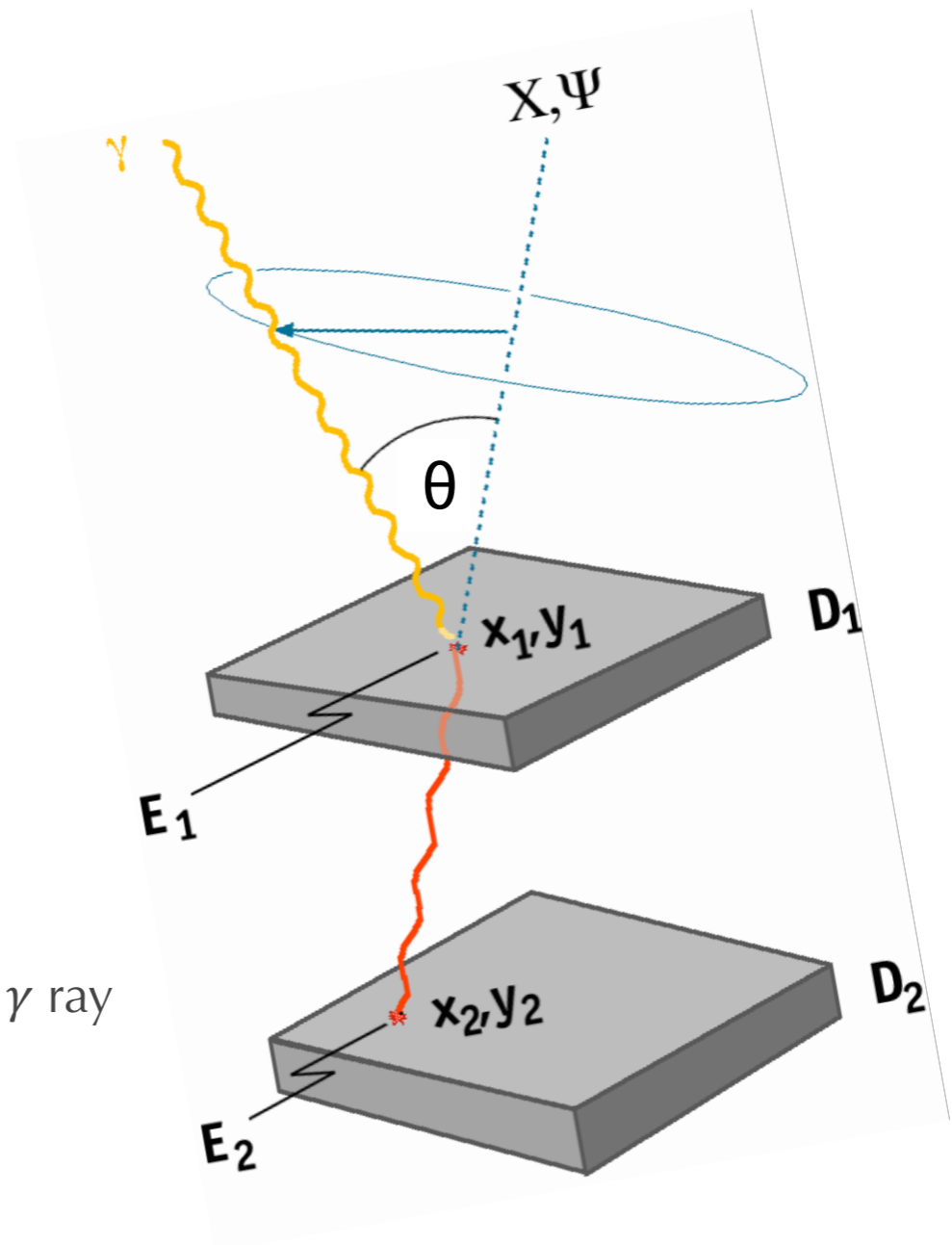
$$\Delta\theta = \frac{m_e c^2}{\sin\theta} \left[ \frac{\Delta E_1^2}{(E_1 + E_2)^4} + \Delta E_2^2 \left( \frac{1}{E_2^2} - \frac{1}{(E_1 + E_2)^2} \right)^2 \right]^{1/2}$$

- dx term:

- » Separating D1 and D2 increases lever arm

- Design trade

- Separating D1 and D2
  - Improves angular resolution
  - Decreases effective area
    - » Decreases efficiency for detecting scattered  $\gamma$  ray

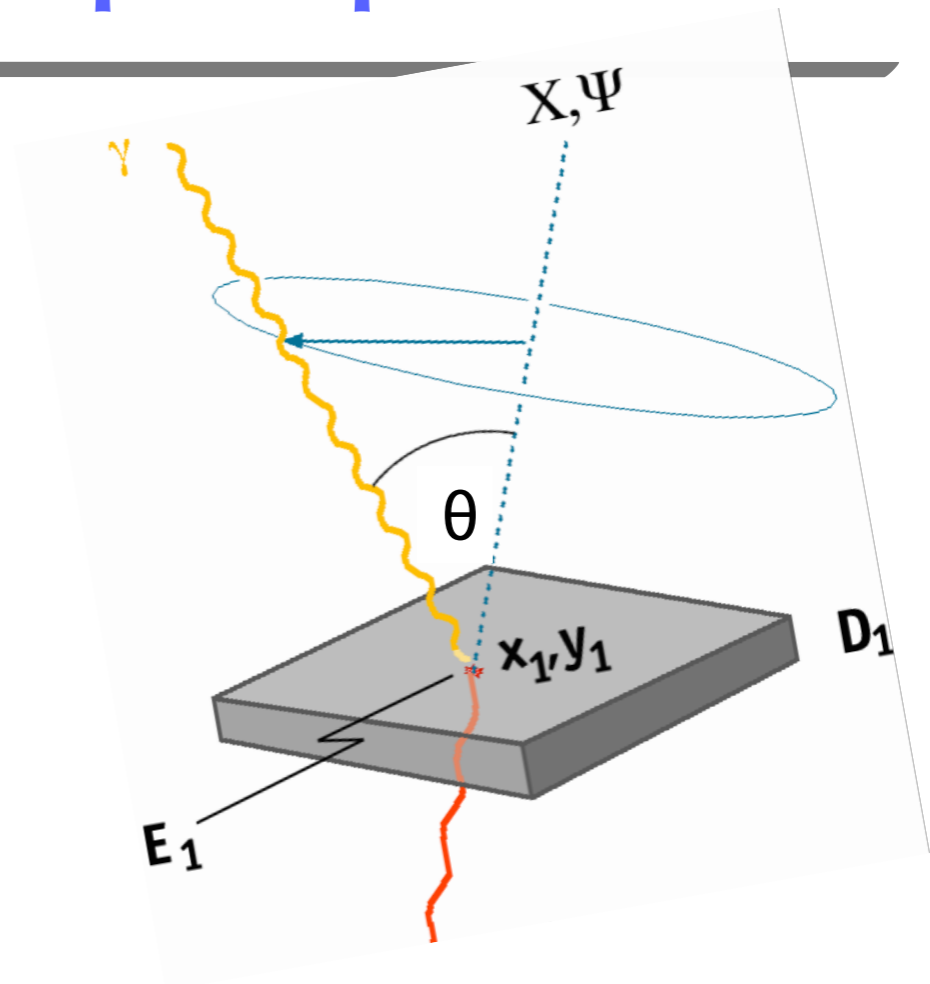


What's the right choice?



# Compton telescope principles

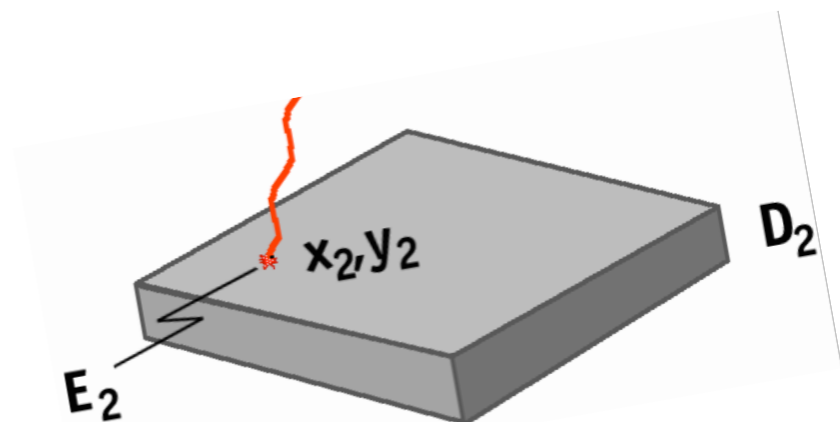
- Compton event reconstruction
  - Width of cone (“angular response measure”) is function of E and position resolution
    - dE term:
$$\Delta\theta = \frac{m_e c^2}{\sin\theta} \left[ \frac{\Delta E_1^2}{(E_1 + E_2)^4} + \Delta E_2^2 \left( \frac{1}{E_2^2} - \frac{1}{(E_1 + E_2)^2} \right)^2 \right]^{1/2}$$
    - dx term:
      - » Separating D1 and D2 increases lever arm



## – Design trade

- Separating D1 and D2
  - Improves angular resolution
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What’s the right choice?

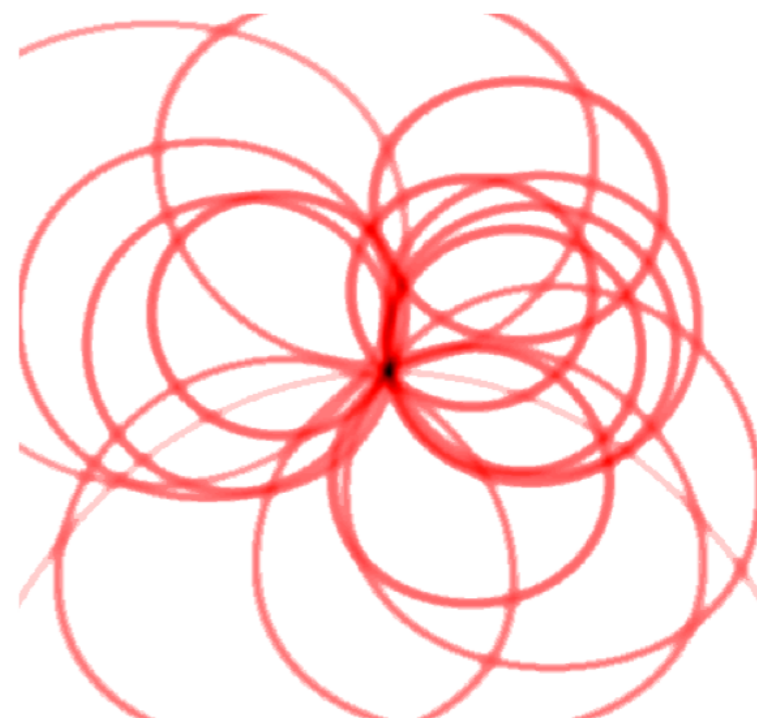






# Compton imaging

- Image
  - Each photon = ring
  - Intersection of many rings
- Issue
  - Source confusion
  - Rejection of sky background
    - Especially the bright atmosphere of Earth
  - Complicated PSF
- Mitigation
  - Best possible E and position resolution
    - Keep them well matched
  - Get more information about the scatter in D1: track the recoil electron

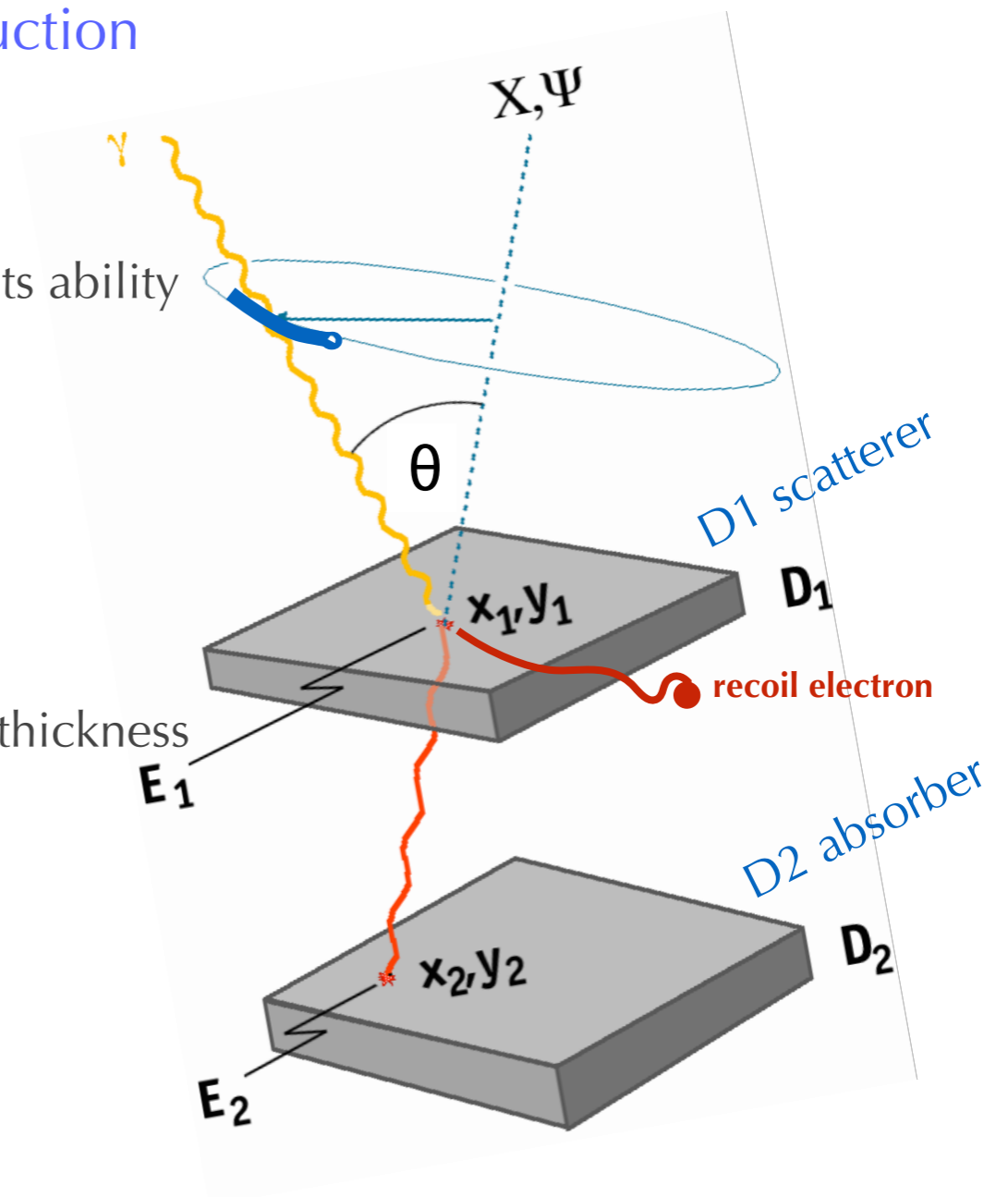


- The origin of a single not-tracked event can be restricted to the so called “event circle”.
- The photon originated at the point of all overlap.



# Electron tracking

- Measuring recoil electron track
  - Allows almost complete interaction reconstruction
  - Reduces Compton ring to an arc
    - Multiple coulomb scattering of recoil electron limits ability to measure initial recoil direction
      - Want low-density D1 to minimize MCS
      - But that minimizes Compton probability
    - But it raises low E threshold for instrument
      - Recoil electron must travel through finite detector thickness
      - e.g. range(500 keV e) < 1 mm of Si
- Note: Doppler broadening
  - Momentum of bound electron is important
    - Significant contribution below ~1 MeV
    - Whether electron tracking or not
  - Sets fundamental limit on measurement uncertainty
  - Low-Z materials make better scatterers





# Monte Carlo example: cubic meter of Si

- Point source at 847 keV
  - $10^4$  shots into  $1 \times 1 \times 1$  m Si cube: ~2 days observing SNIa at 15 Mpc
  - Require 2 interactions, no electron tracking: ~8500 reconstructed  $\gamma$  rays

Perfect detector, only Doppler broadening

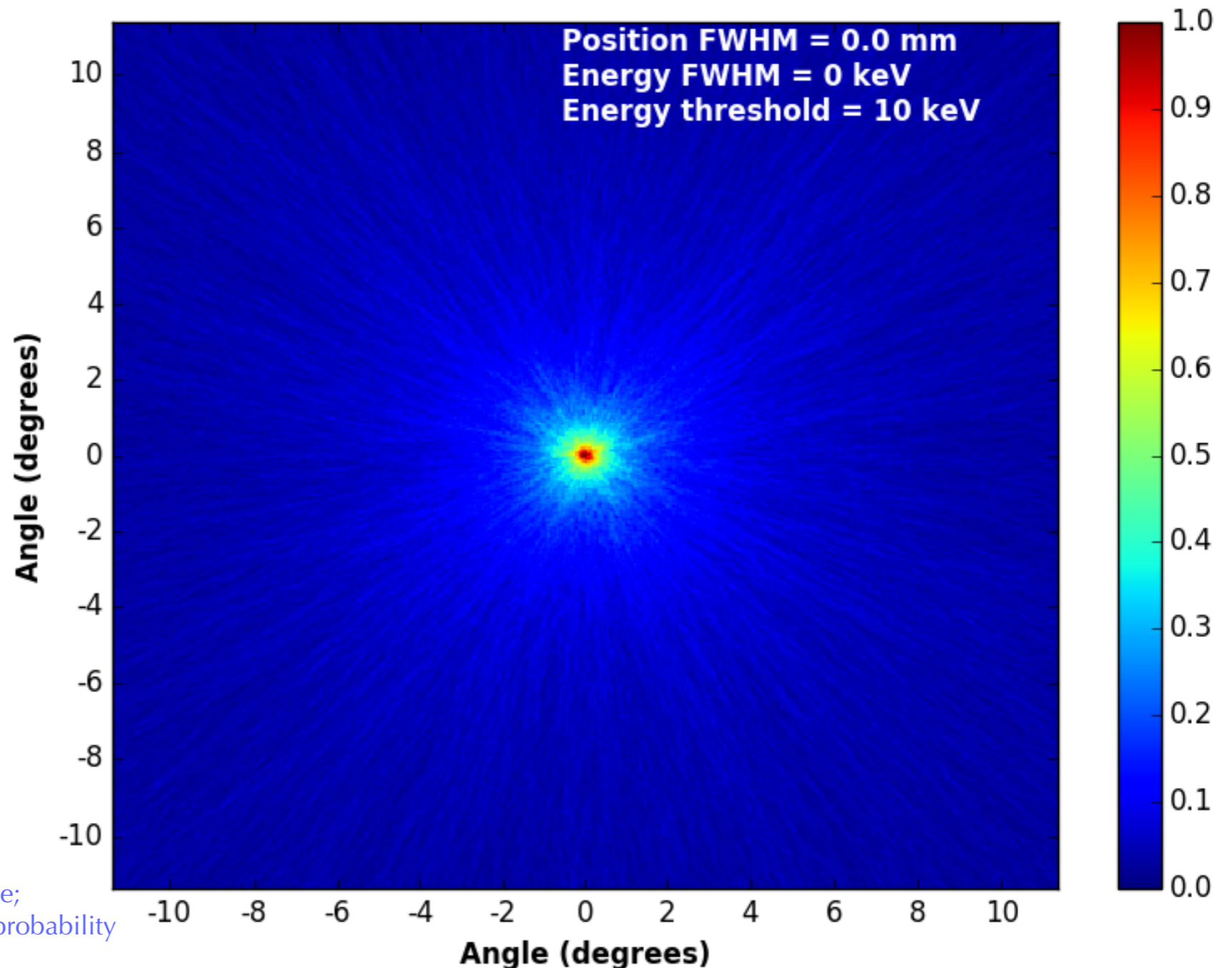
- Individual rings (almost) visible

“Realistic” detector parameters

- Rings are wide
- Rings are blurred over large area

*Note to Compton experts:*

*This looks worse than expected because solid Si puts interactions too close together (reduces lever arm). Vertical gaps between Si wafers are a Good Thing™*



Color scale is relative;  
fraction of peak probability



# Monte Carlo example: cubic meter of Si

- Point source at 847 keV
  - $10^4$  shots into  $1 \times 1 \times 1$  m Si cube:  $\sim 2$  days observing SNIa at 15 Mpc
  - Require 2 interactions, no electron tracking:  $\sim 8500$  reconstructed  $\gamma$  rays

Perfect detector, only Doppler broadening

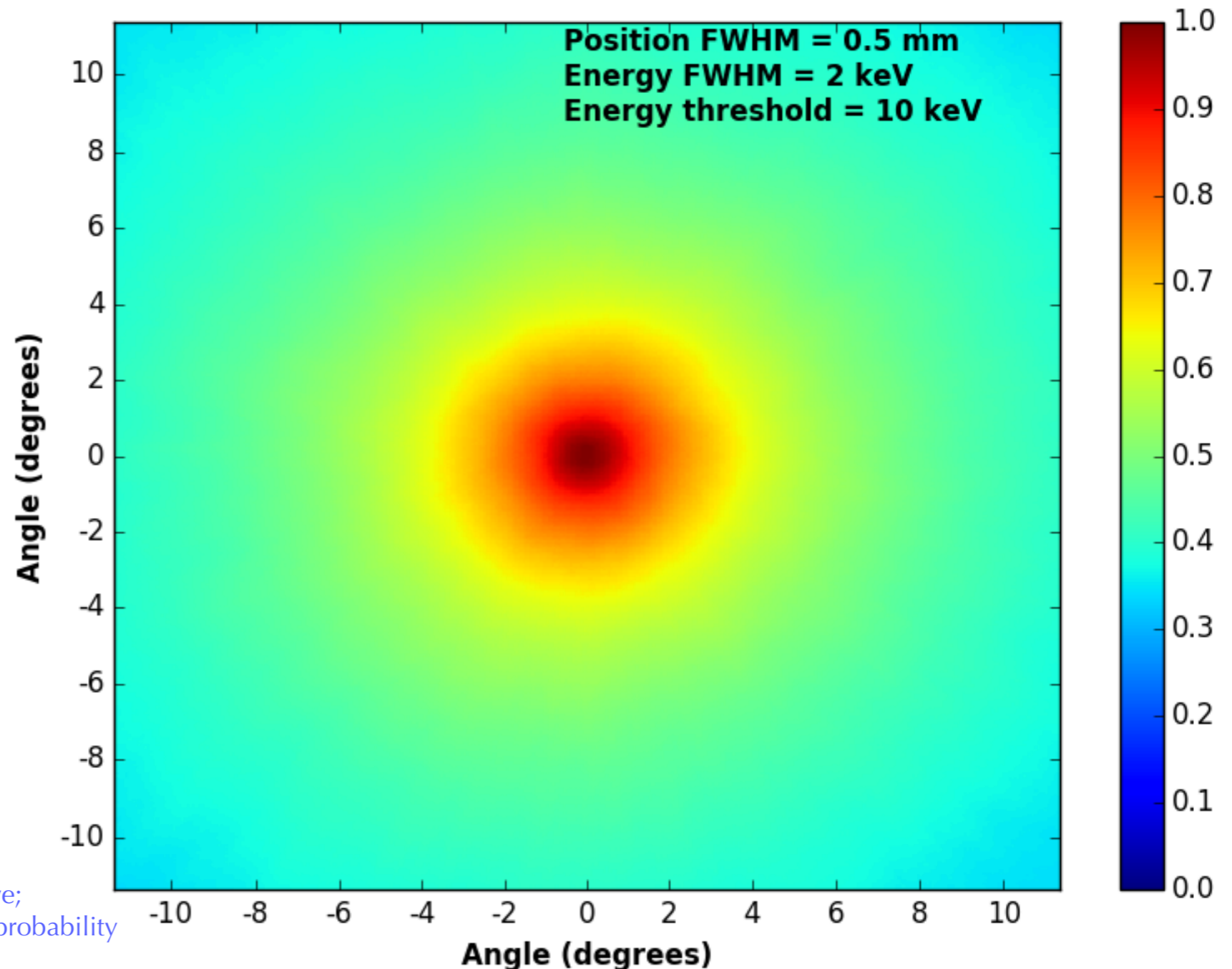
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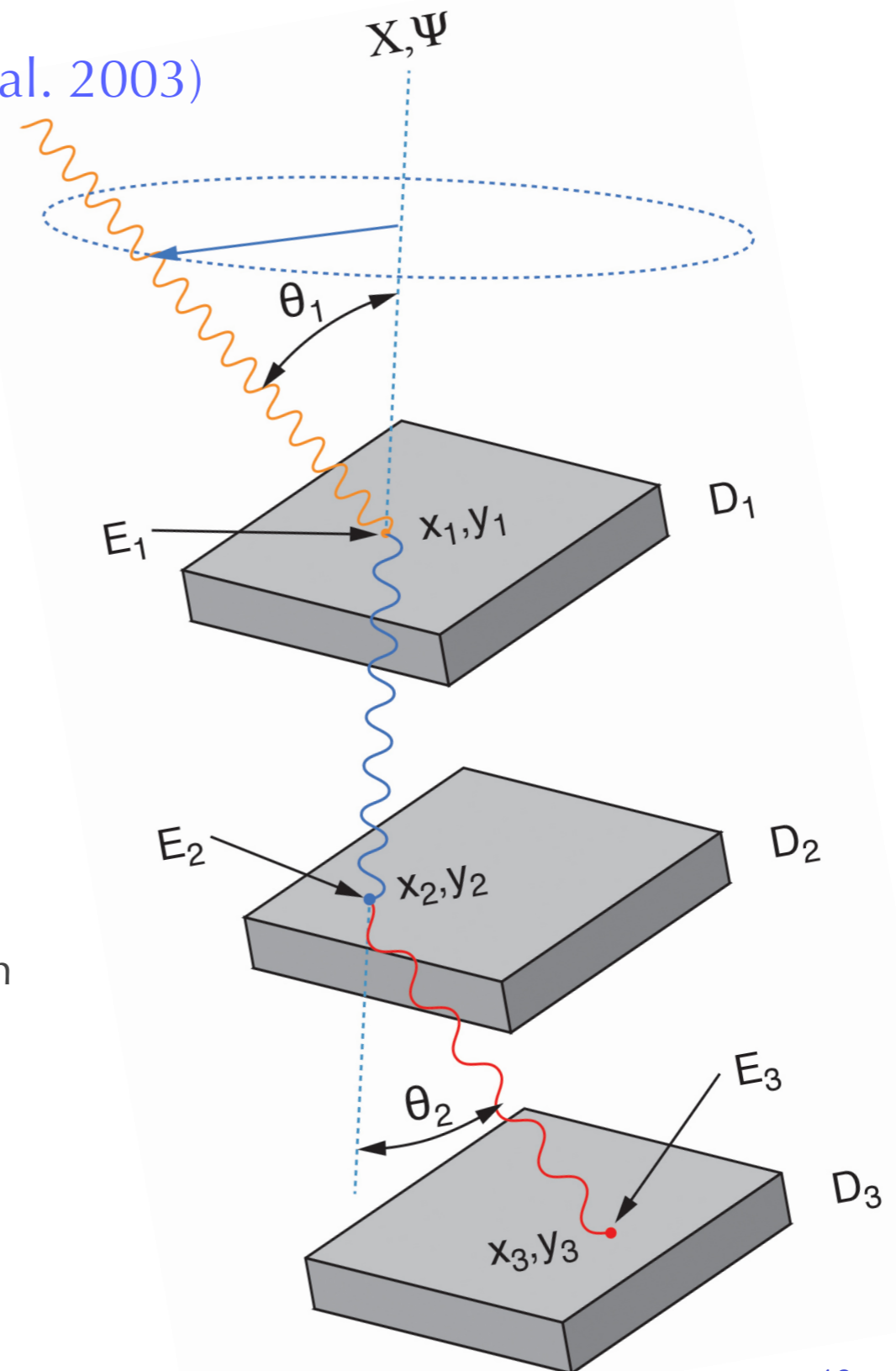


# Multiple Compton scatters

- But wait... we can be more clever
  - No need for full absorption (patent: Kurfess et al. 2003)
  - If  $\gamma$  ray scatters at least twice (i.e. at least 3 interaction points in telescope), can measure incident energy without totally absorbing the incident photon

$$E_0 = E_1 + \frac{1}{2}E_2 + \frac{1}{2} \left[ E_2^2 + \frac{4m_e c^2 E_2}{1 - \cos\theta_2} \right]^{1/2}$$

- Thick, high-Z D2 absorber is not required!
  - Imagine a large array of Si detectors....
  - “Easy” to get excellent position and energy resolution





# Multiple-Compton technique

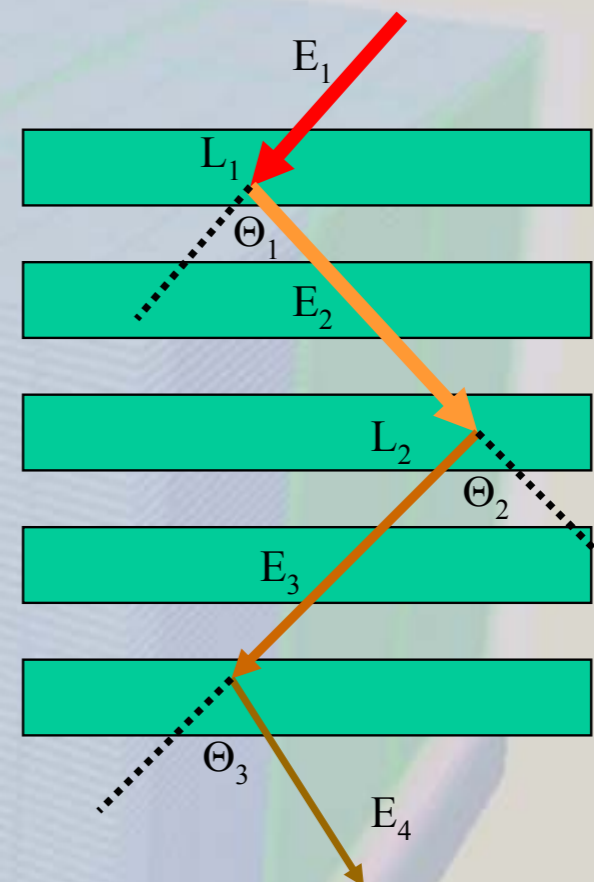
## Three Gamma Interaction Technique

$\Theta_2$  is computed from geometry

$$E_1 = L_1 + \frac{L_2 + \left[ L_2^2 + \frac{4m_e c^2 L_2}{1 - \cos \Theta_2} \right]^{\frac{1}{2}}}{2}$$

above is derived from:

$$\begin{cases} \cos \Theta_2 = 1 - m_e c^2 \left( \frac{1}{E_3} - \frac{1}{E_2} \right) \\ L_2 = E_2 - E_3 \end{cases}$$



- **Unknown source:** 3 interactions required to determine energy,  $E_1$
- **Known source:** 2 interactions required to determine energy,  $E_1$
- Does **not** require total energy absorption
- Efficient Compton telescope, even if using *silicon* detectors
- **Ordering algorithm is essential**

See N1-1: Wulf et al. for prototype results



# Multiple-Compton technique

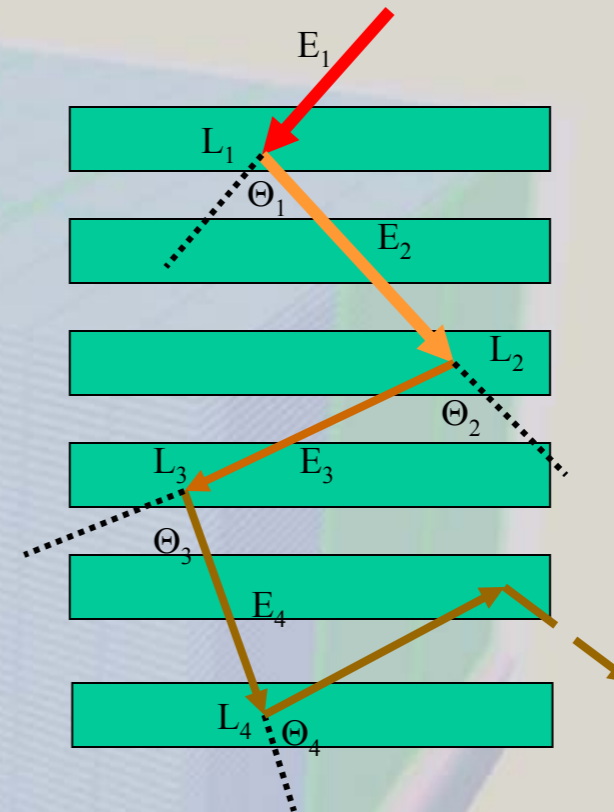
- Warning: order of scatters is essential, and number of scatters can be large

## Multiple Estimates of Incident Gamma Ray Energy

$$E_1 = L_1 + \frac{L_2 + \left[ L_2^2 + \frac{4m_e c^2 L_2}{1 - \cos \Theta_2} \right]^{\frac{1}{2}}}{2} ;$$

$$E_1 = L_1 + L_2 + \frac{L_3 + \left[ L_3^2 + \frac{4m_e c^2 L_3}{1 - \cos \Theta_3} \right]^{\frac{1}{2}}}{2}$$

$$E_1 = L_1 + L_2 + L_3 + \frac{L_4 + \left[ L_4^2 + \frac{4m_e c^2 L_4}{1 - \cos \Theta_4} \right]^{\frac{1}{2}}}{2}$$



- Ordering algorithm is essential
- N hits result in N! possible sequences
- N interactions provide N-2 estimates of  $E_1$
- Sequence with the most consistent estimates of  $E_1$  is accepted
- Currently accept only events with 4 to 8 hits, and fully absorbed events with 3 hits
- In the future: check Klein-Nishina and absorption probabilities; electron tracking

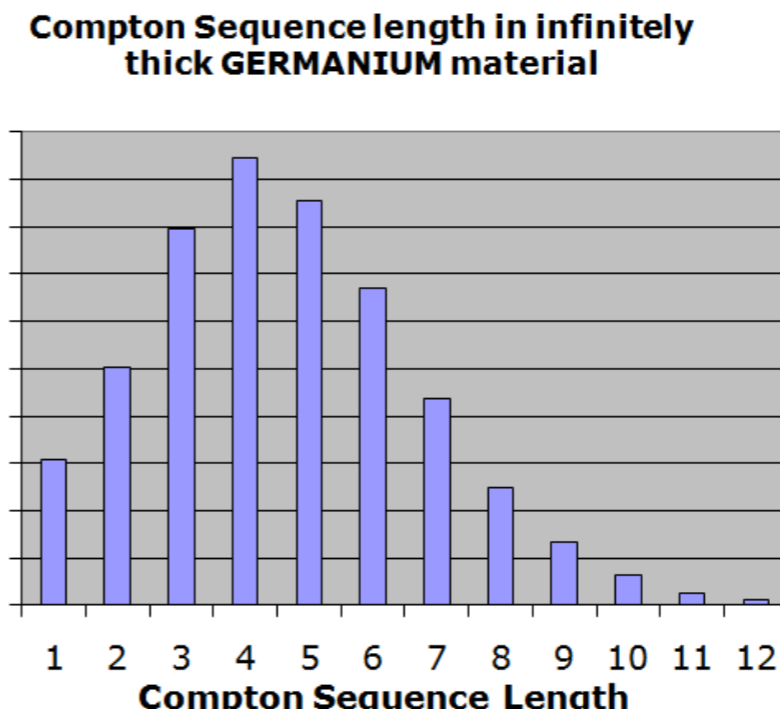
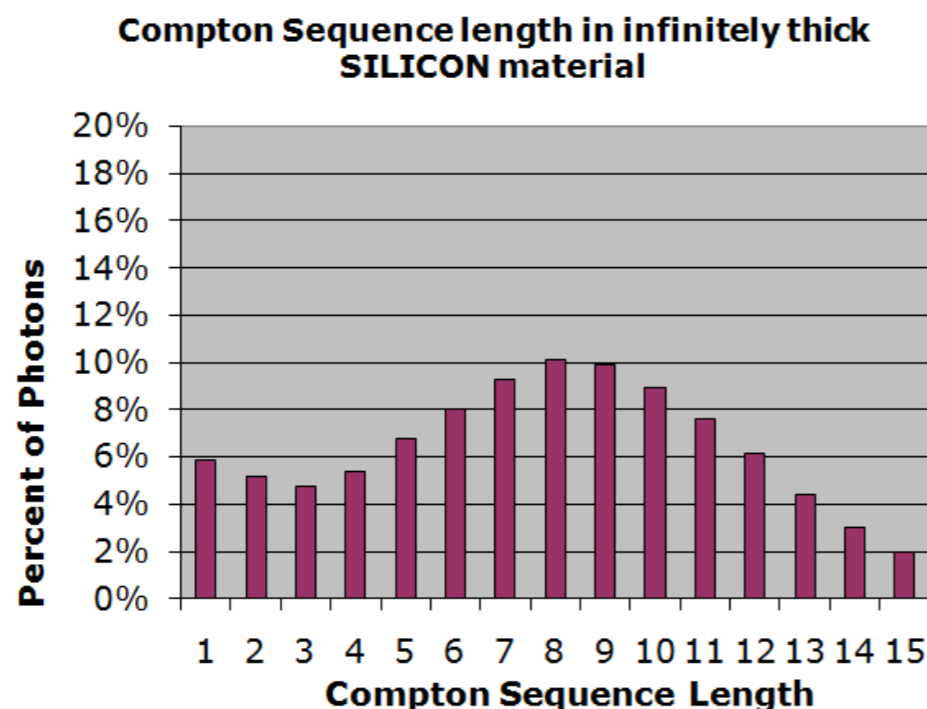




# How common are multiple Comptons?

- Very, so don't throw them out
  - Maybe Si-only Compton telescope isn't optimal, but needs higher Z somewhere

asCi design considerations - material



Ge :  $4 \pm 2$  interactions needed to transform full energy (75% of photons)

Si :  $8 + 3 - 2$  interactions needed to transform full energy (75% of photons)

Ge : provides sufficient number of interactions (algorithms require  $\geq 2$ ) while providing enough stopping power to prevent too many interactions (makes reconstruction impossible, since they increase with  $n!$ ) and increase the chance of the full photon energy being deposited.





# Si ACT (and variants)

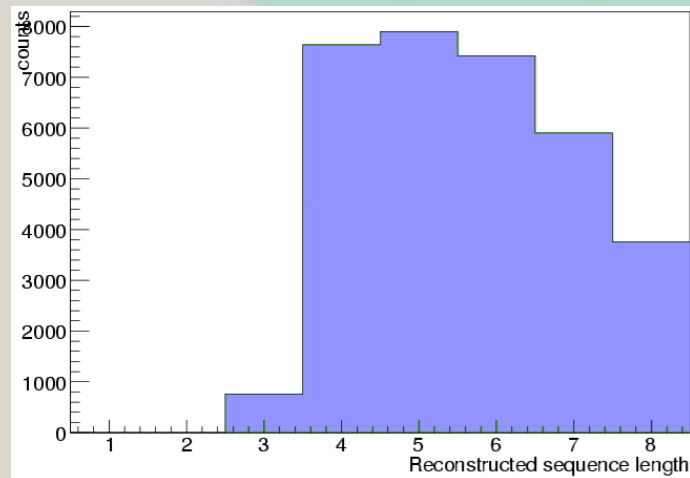
## Mass model modifications



Compton sequence length for: Horizon cut  $92.5^\circ$ ;  $E = 847 \pm 22.75$  keV; 1 mil events

**Main Model**  
64 x 3mm Si

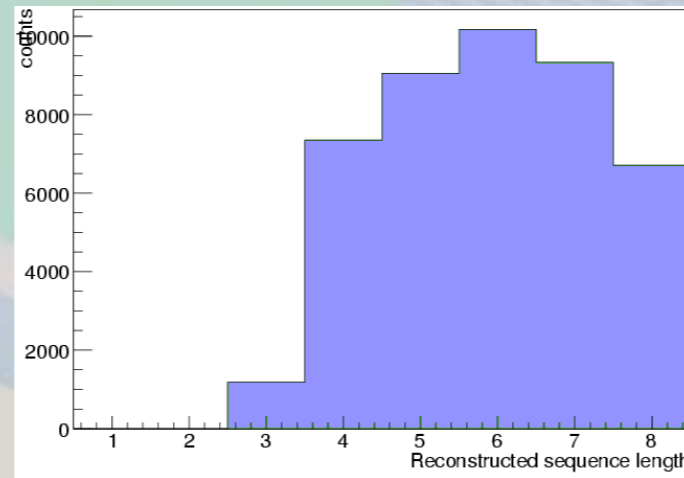
**Baseline**



**Eff. Area  $860 \text{ cm}^2$**   
**FWHM Ang.Res:  $1.44^\circ$**

**Thicker Si (Wafer Bonding\*)**  
25 x 7.5mm Si

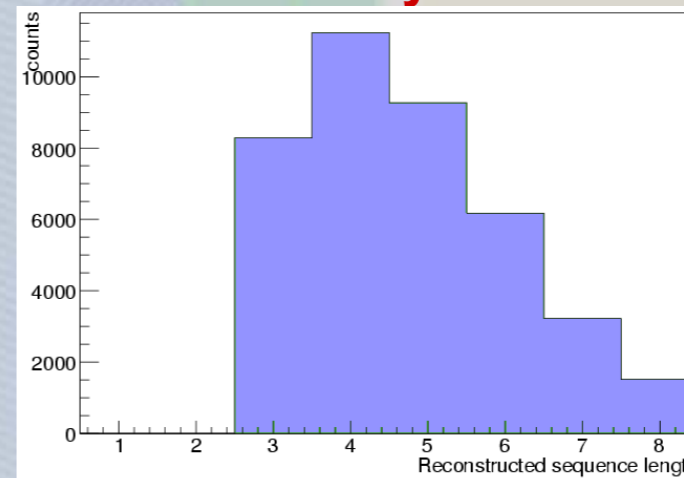
**Less dead material:**  
**longer sequences are recoverable**



**Eff. Area  $1266 \text{ cm}^2$**   
**FWHM Ang.Res:  $1.9^\circ$**

**Si interleaved with CZT**  
(5mm CZT + 12 x 3mm Si) x 4

**High-Z material interleaved:**  
**no need for longer sequences, because more events are fully absorbed**



**Eff. Area  $1073 \text{ cm}^2$**   
**FWHM Ang.Res:  $1.78^\circ$**

**\*see N35-62: Philips et al. for wafer bonding**





# Si ACT (and variants)

## Mass model modifications

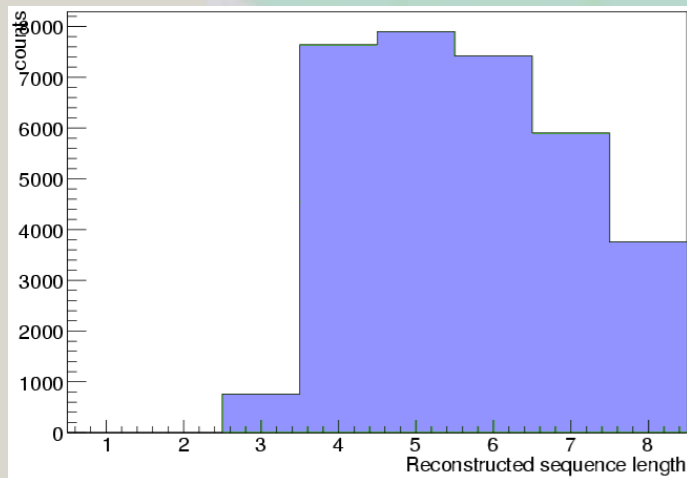


Compton sequence length for: Horizon cut  $92.5^\circ$ ;  $E = 847 \pm 22.75$  keV

Insert high Z, high resolution det  
Truncates scatter sequence  
Preserves dE, ARM

**Main Model**  
64 x 3mm Si

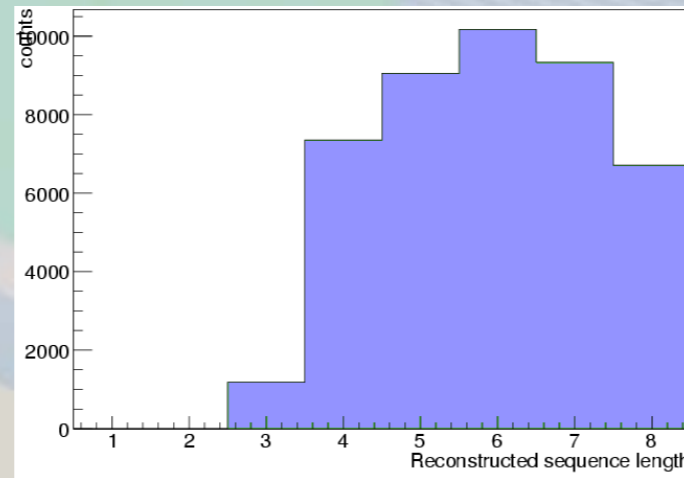
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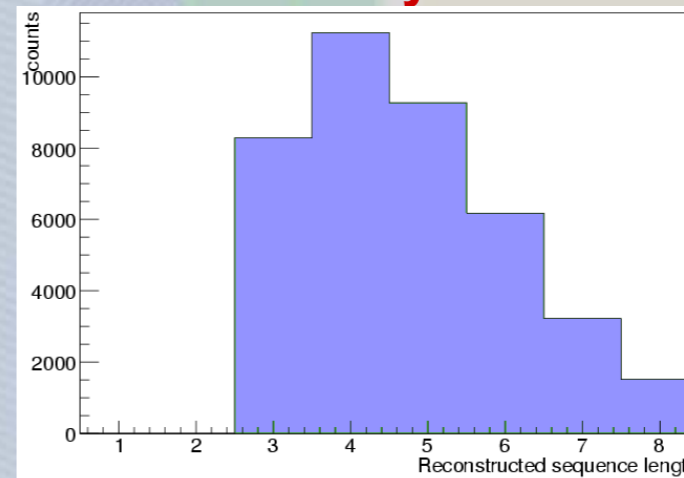
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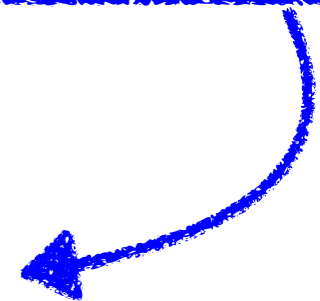
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**FWHM Ang.Res:  $1.78^\circ$**



\*see N35-62: Philips et al. for wafer bonding





# “Requirements” for high-res Compton tele

---

- Low Z scatterer
  - Minimizes Doppler broadening (most important below MeV)
  - Minimizes MCS of recoil electron, if tracking
- High Z absorber
  - Good stopping power to absorb scattered gamma (and minimize multi-Compton)
- High efficiency
  - Proper scatterer and absorber to give highest possible efficiency
  - Compact (as possible) to maximize geometric cross section for interaction
- Excellent energy resolution
  - Well matched with  $d^3x$
- Fine position resolution
  - Well matched with  $dE$ 
    - Thumb:  $\sim 1$  mm and  $\sim$ few keV are commensurate
- Low-power electronics
  - Preserve intrinsic  $dE$ ,  $d^3x$  of detectors while staying within power budget
- Minimal passive mass within detection volume
  - Interactions can be missed in passive material, and kill Compton performance
  - Minimize structural supports, co-located electronics
  - Minimize detector guard rings



# Polarimetry

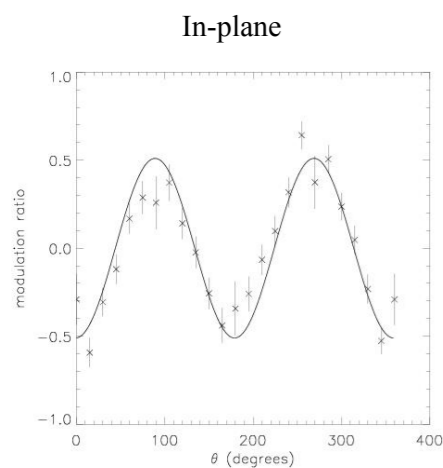
- Compton telescope is good polarimeter
  - Compton scatter preferentially in direction perpendicular to polarization vector
  - Measure intrinsic polarization of  $\gamma$ -ray source by measuring modulation in scatter angles in detector

## Polarization response

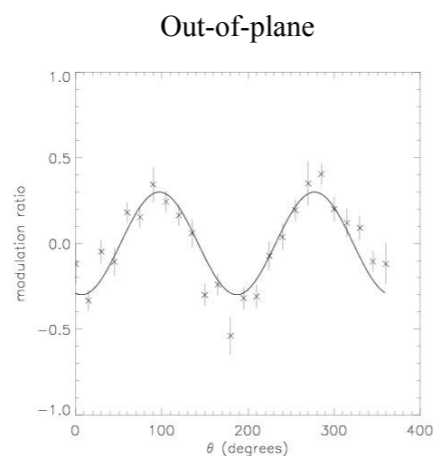
Crab-like source (ref: Jourdain & Roques 2009)

Energy range (MeV)	Selections	Modulation $\mu_{100}$	Source ( $s^{-1}$ )	Atmosph. bgd ( $s^{-1}$ )	CGB ( $s^{-1}$ )	Cosmic-ray induced bgd ( $s^{-1}$ )	MDP $_{3\sigma}$ (c)
0.2 – 2	2+ events without e- tracking $\theta_{EHC}=20^\circ, \theta_{ARM}=3.5^\circ$	0.305	28.3	15.0	61.4	7.0 (a)	0.37%
3 – 10	3+ events with e- tracking $\theta_{EHC}=20^\circ, \theta_{ARM}=1.5^\circ$	0.124	0.13	0.36	0.10	0.37 (b)	19.2%

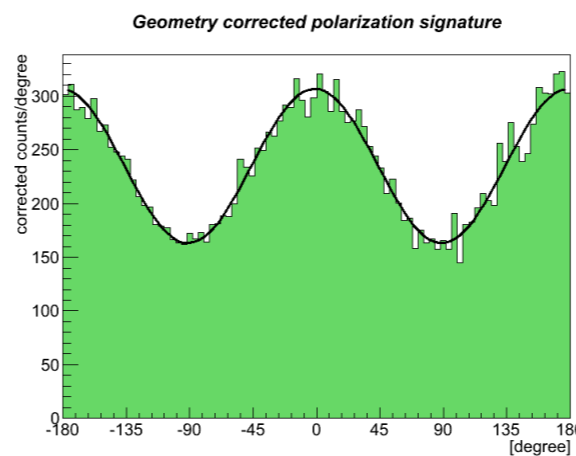
## Modulation ratios for 2-layer instrument



- High polarization ratio.
- Short lever arm.
- High geometric efficiency for thick detectors (strip pitch < thickness).
- Data more difficult to process.



- Lower polarization ratio.
- Longer lever arm.
- Efficiency rises as  $\sim N^2$ .
- Data simpler to process.



Polarigramme for a Crab-like source on axis in the range 0.2 – 2 MeV, yielding a modulation  $\mu_{100} = 0.305$

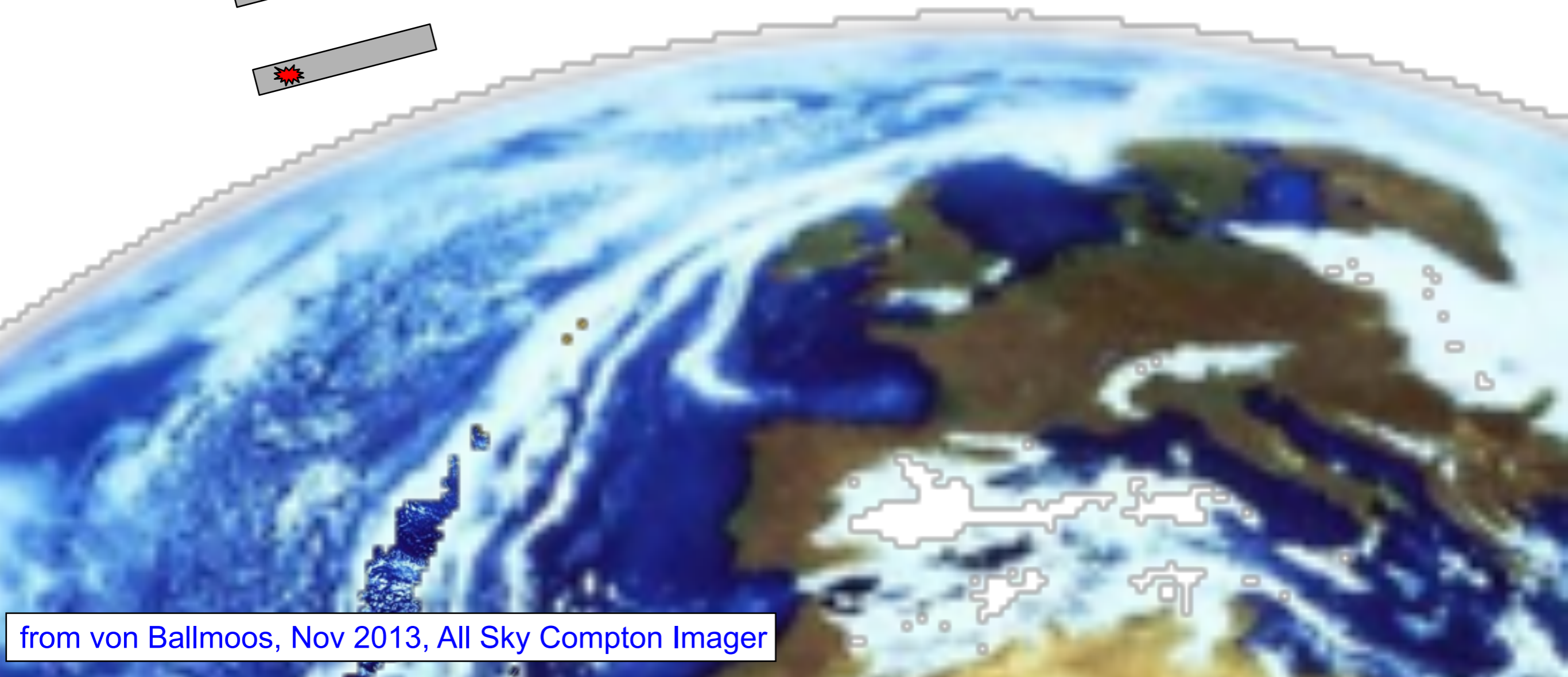
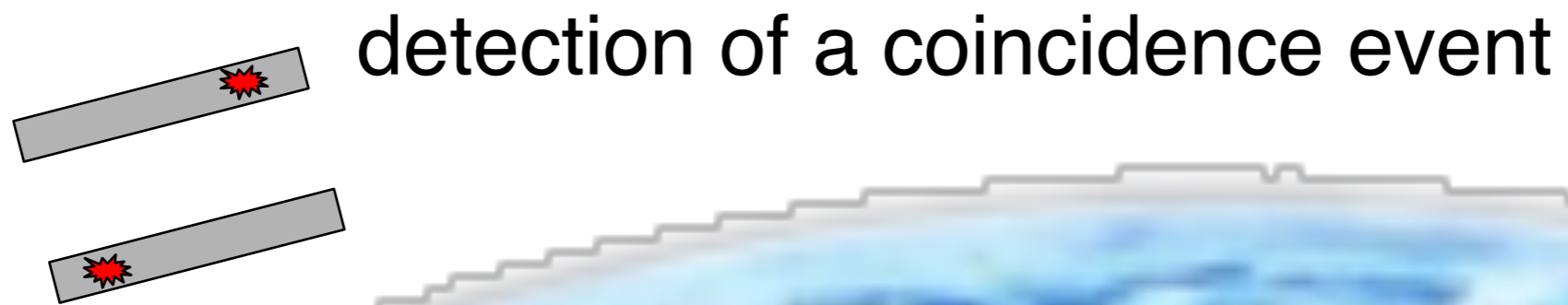
(a) Activation from both primary and secondary (i.e. semi-trapped) protons; (b) Activation from primary and secondary protons + prompt reactions from primary protons, and secondary protons and leptons; (c)  $3\sigma$  minimum detectable polarization for  $T_{obs} = 10^6$  s

- Minimum detectable polarization:

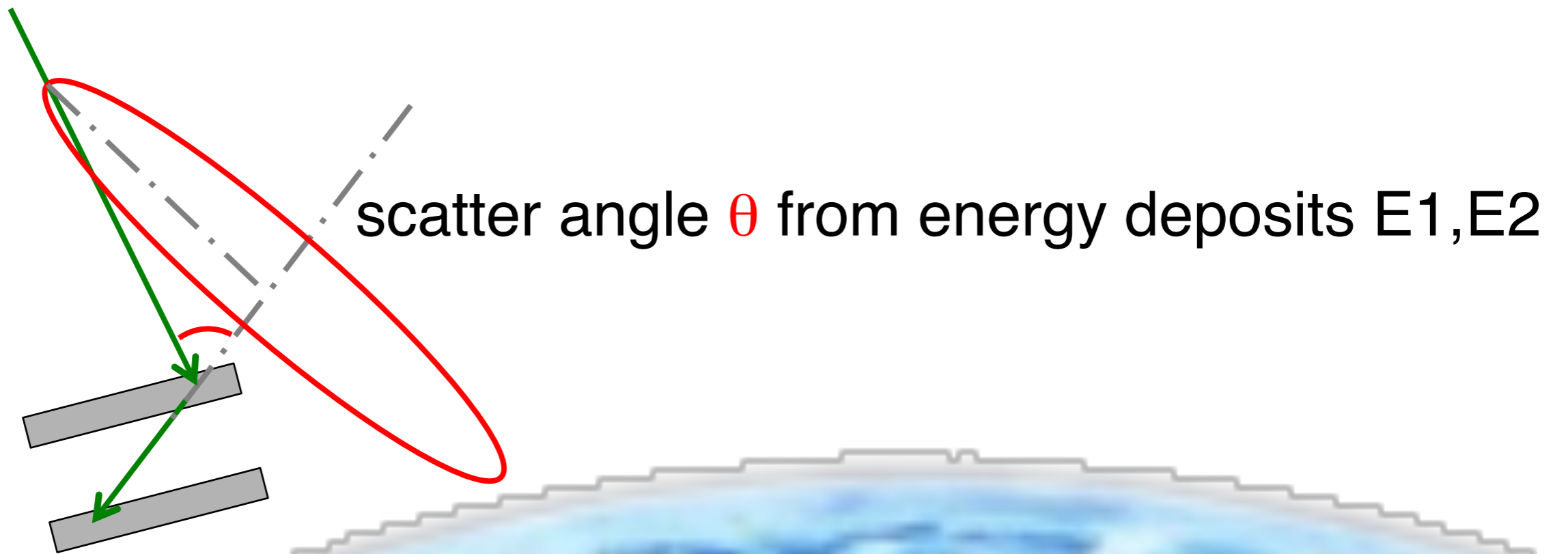
$$MDP_{3\sigma} = \frac{3\sqrt{C_S + B}}{\mu_{100} C_S \sqrt{T_{obs}}}$$

where  $B$  and  $C_S$  are the background and source count rates and  $\mu_{100}$  the modulation

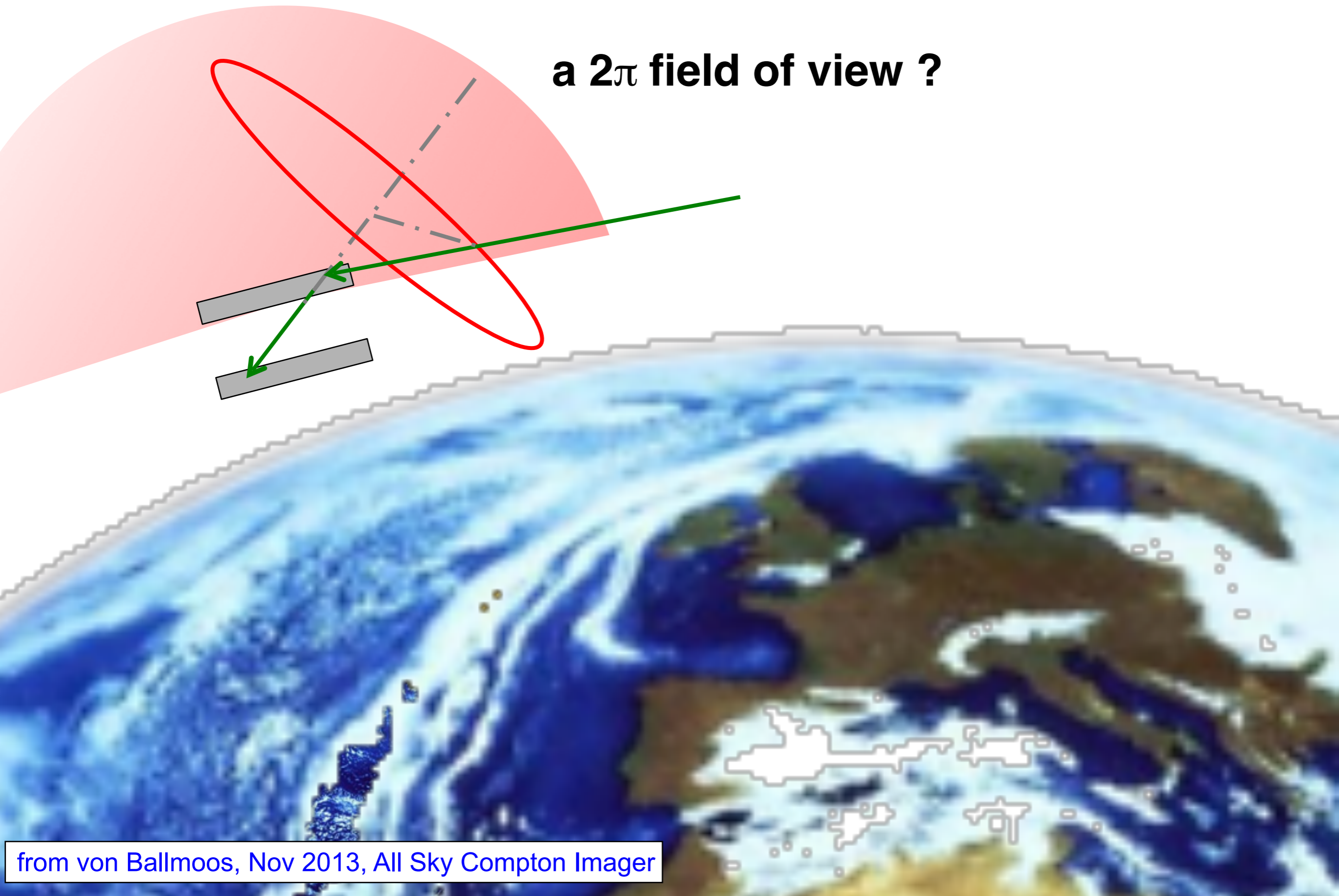
# Compton telescopes have wide fields



# Compton telescopes have wide fields

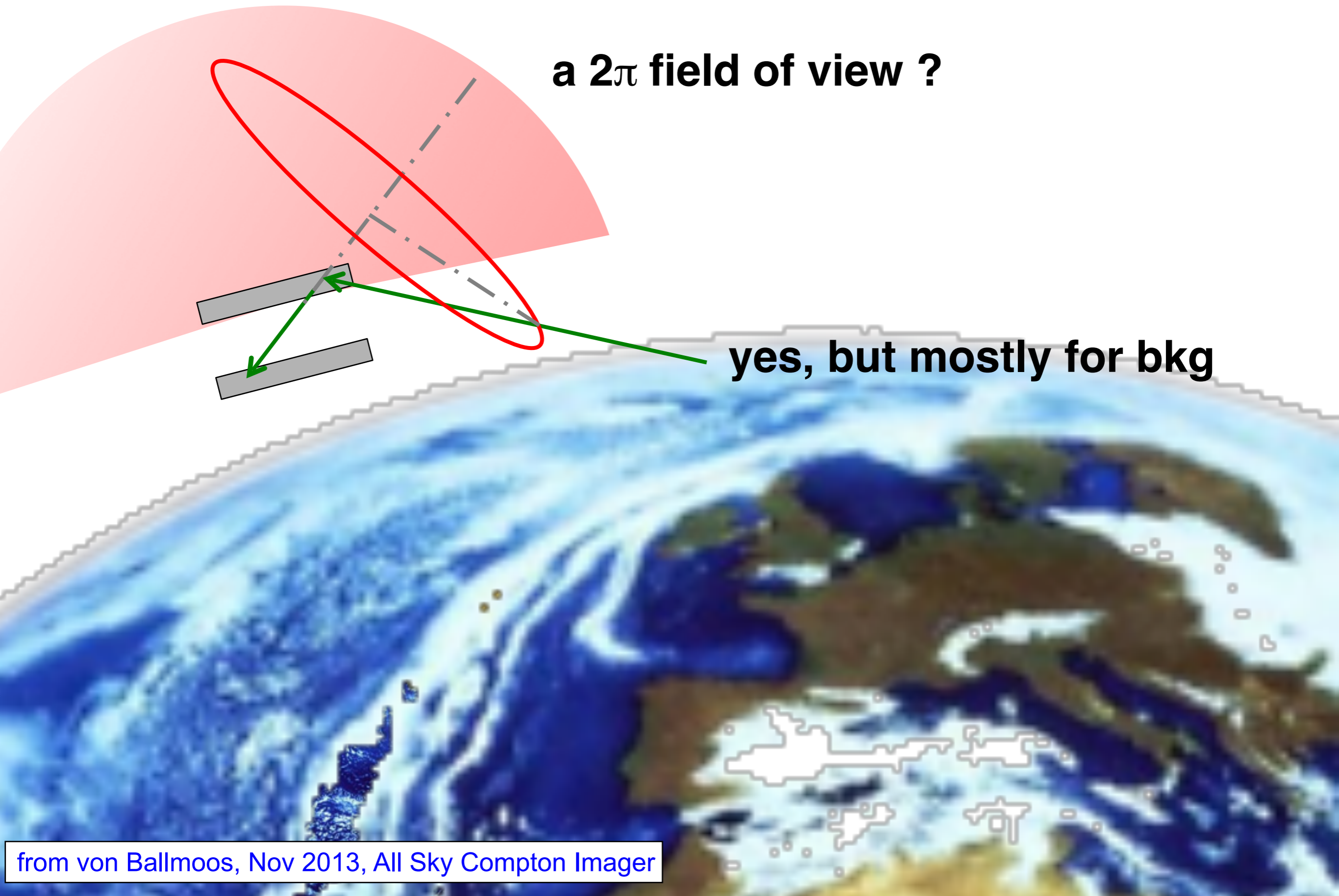


# Compton telescopes have wide fields



**a  $2\pi$  field of view ?**

# Compton telescopes have wide fields



**a  $2\pi$  field of view ?**

**yes, but mostly for bkg**



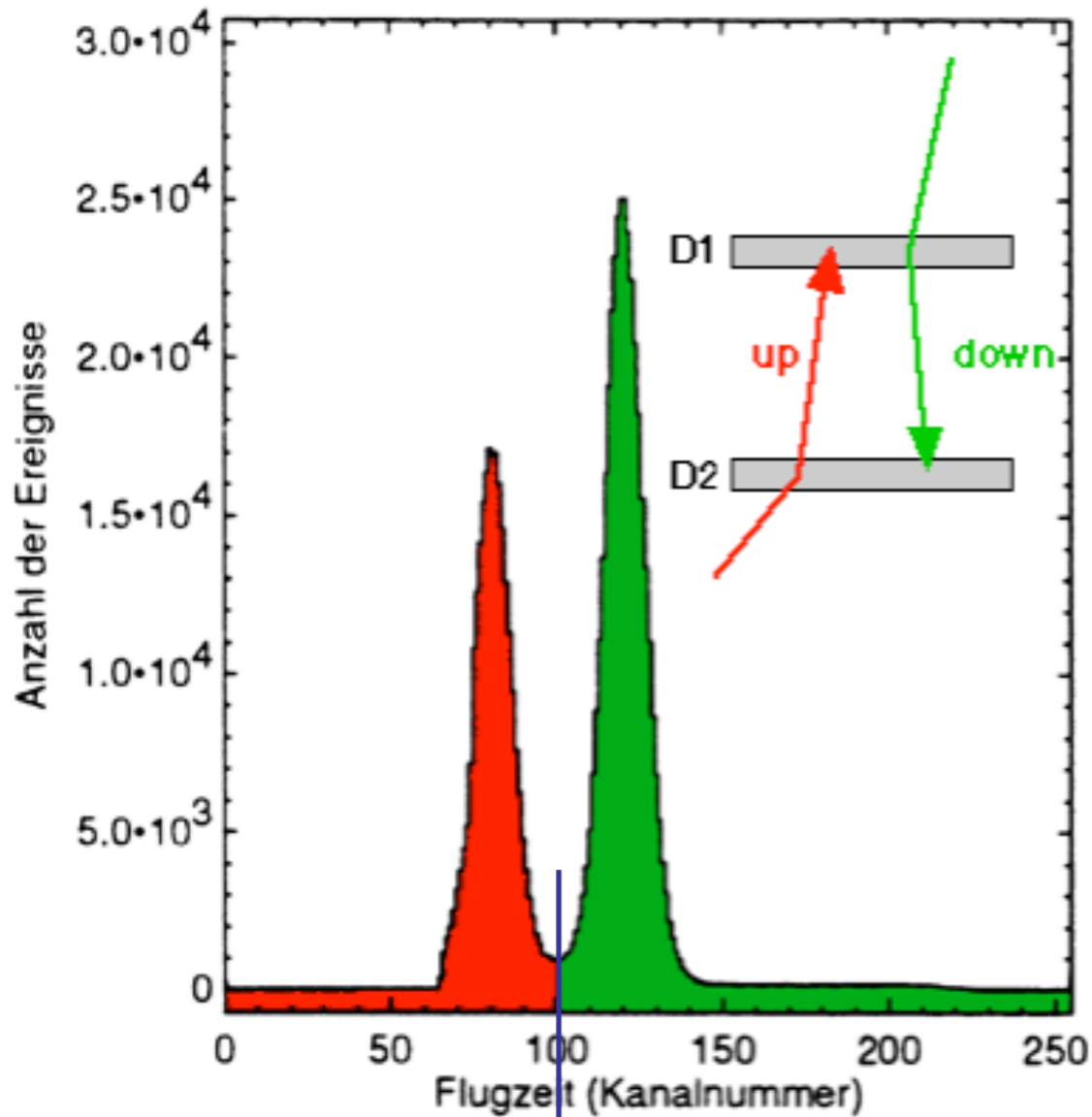
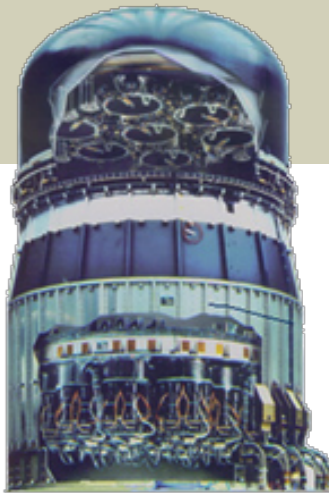


# Backgrounds

---

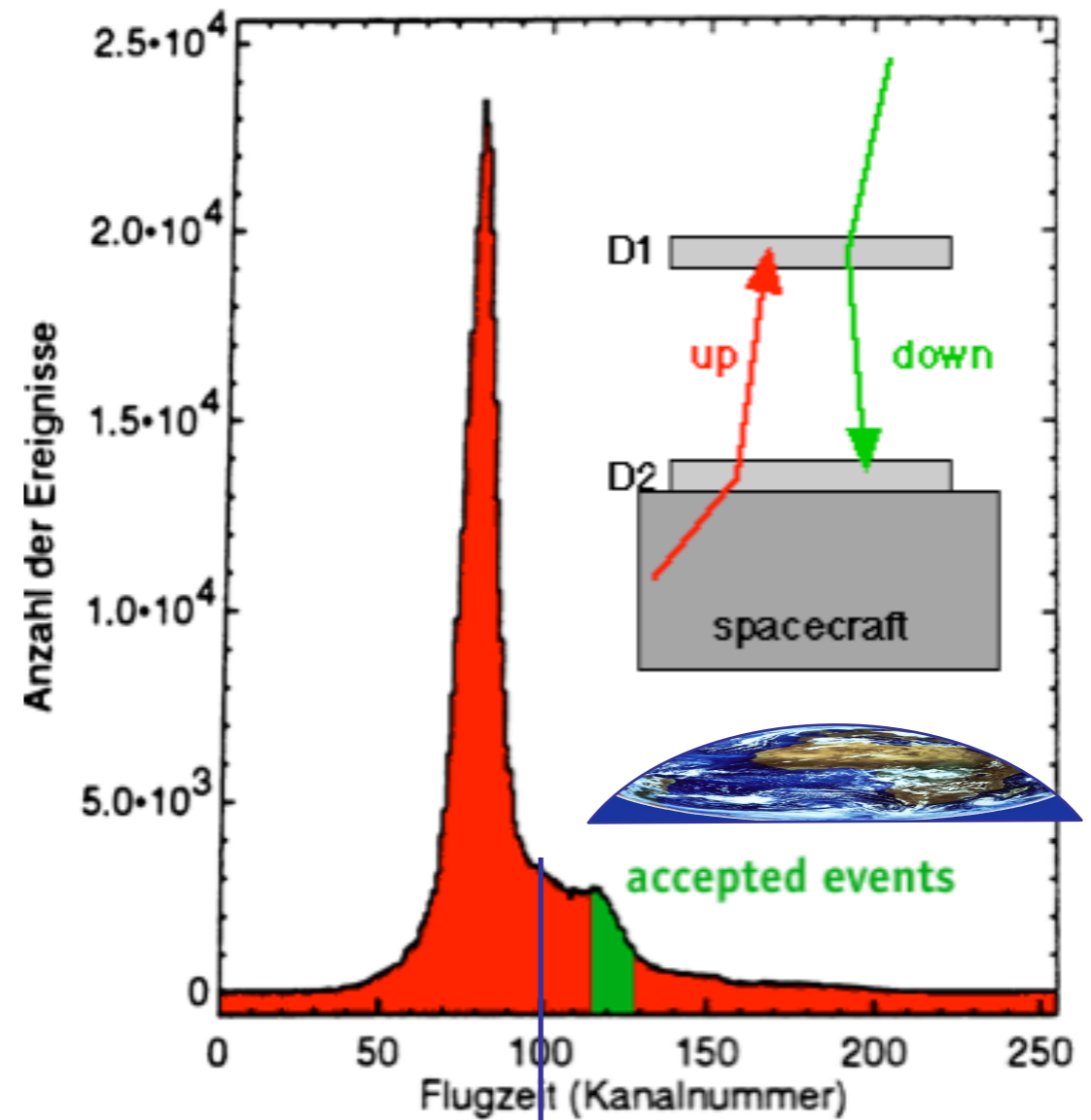
- Sources of background
  - Same as LAT
    - CR primaries, trapped particles, particle albedo
    - Prompt CR secondaries
      - Atmo gammas and local gammas
        - » Atmosphere is brightest source
        - » Beware your spacecraft, the pressure vessel on your gas TPC, etc
  - And below 10 MeV, beware radioactivities
    - Self-activity and CR-induced activation
- Mitigating the backgrounds
  - Fight the bkg
    - Shielding
      - Passive
      - Active anti-coincidence shielding
    - Bkg discrimination
      - Pattern recognition
      - Pulse shape discrimination
      - Time-of-flight
  - Avoid the bkg
    - Optimal orbit
    - Minimize passive material
    - Choose low-bkg materials

# Time of Flight coincidence - COMPTEL data



<- upward | downward ->  
**COMPTEL calibration data**

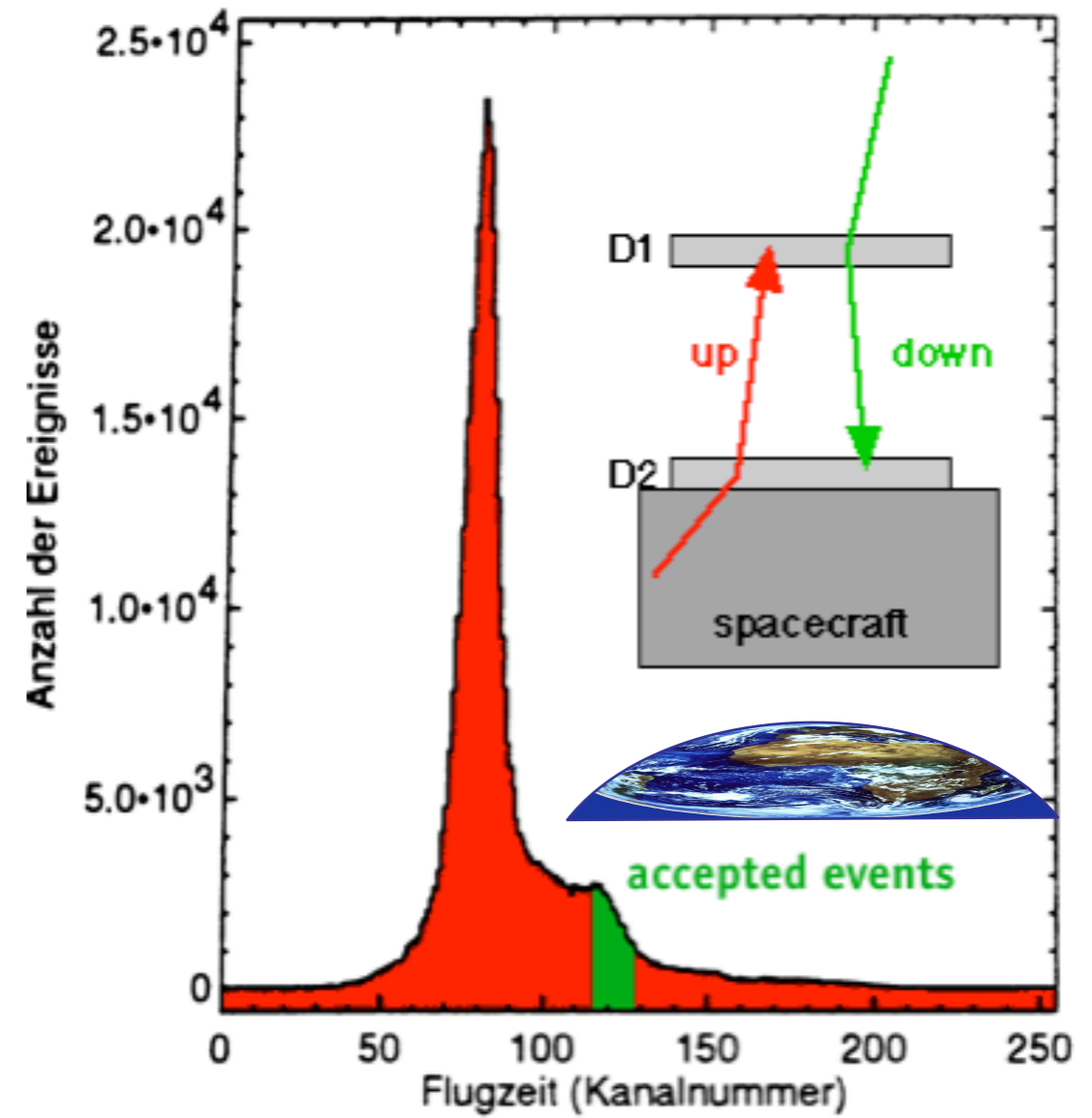
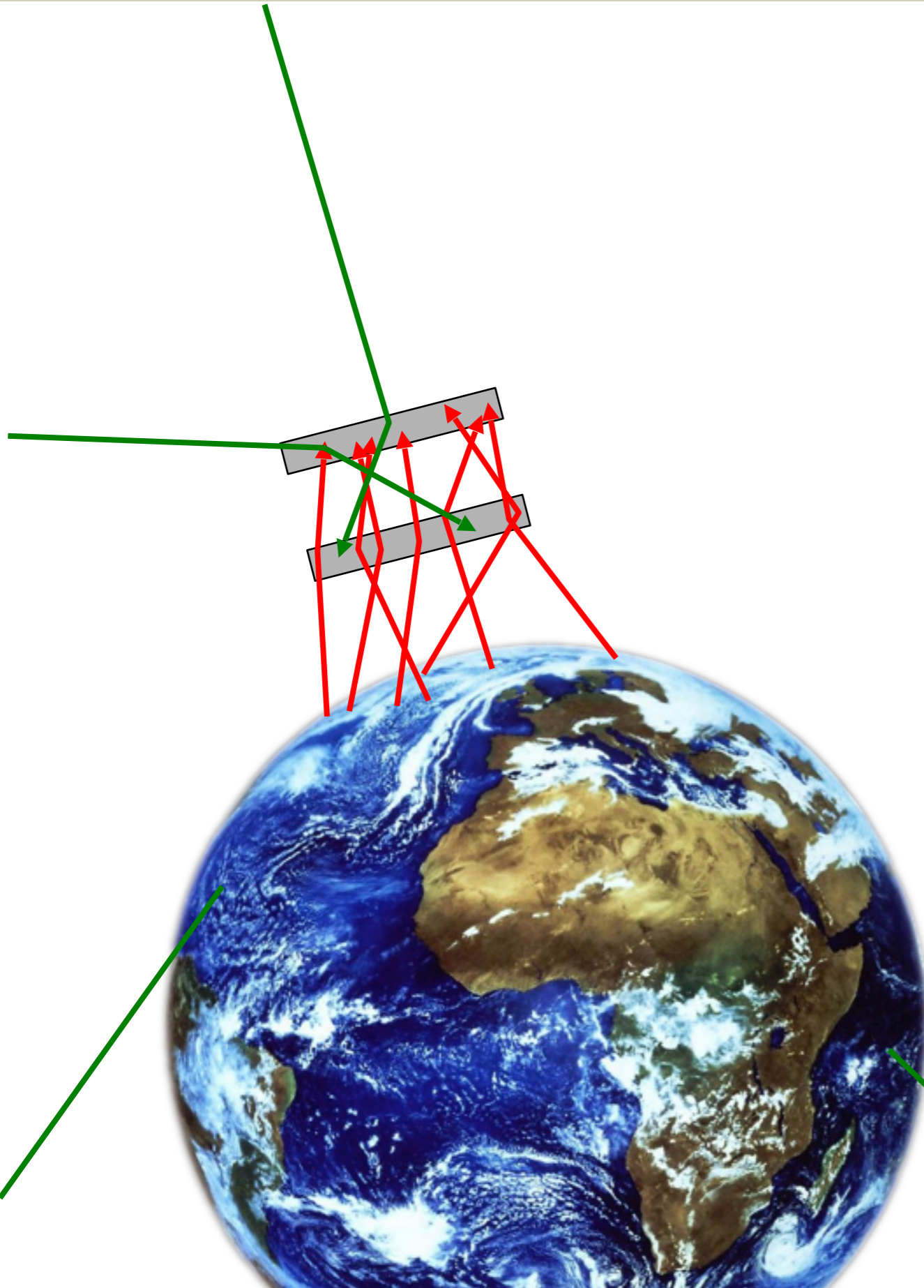
channel width : 0.25 ns  
 distance D1-D2 : 1.5 m  $\approx$  5 ns)



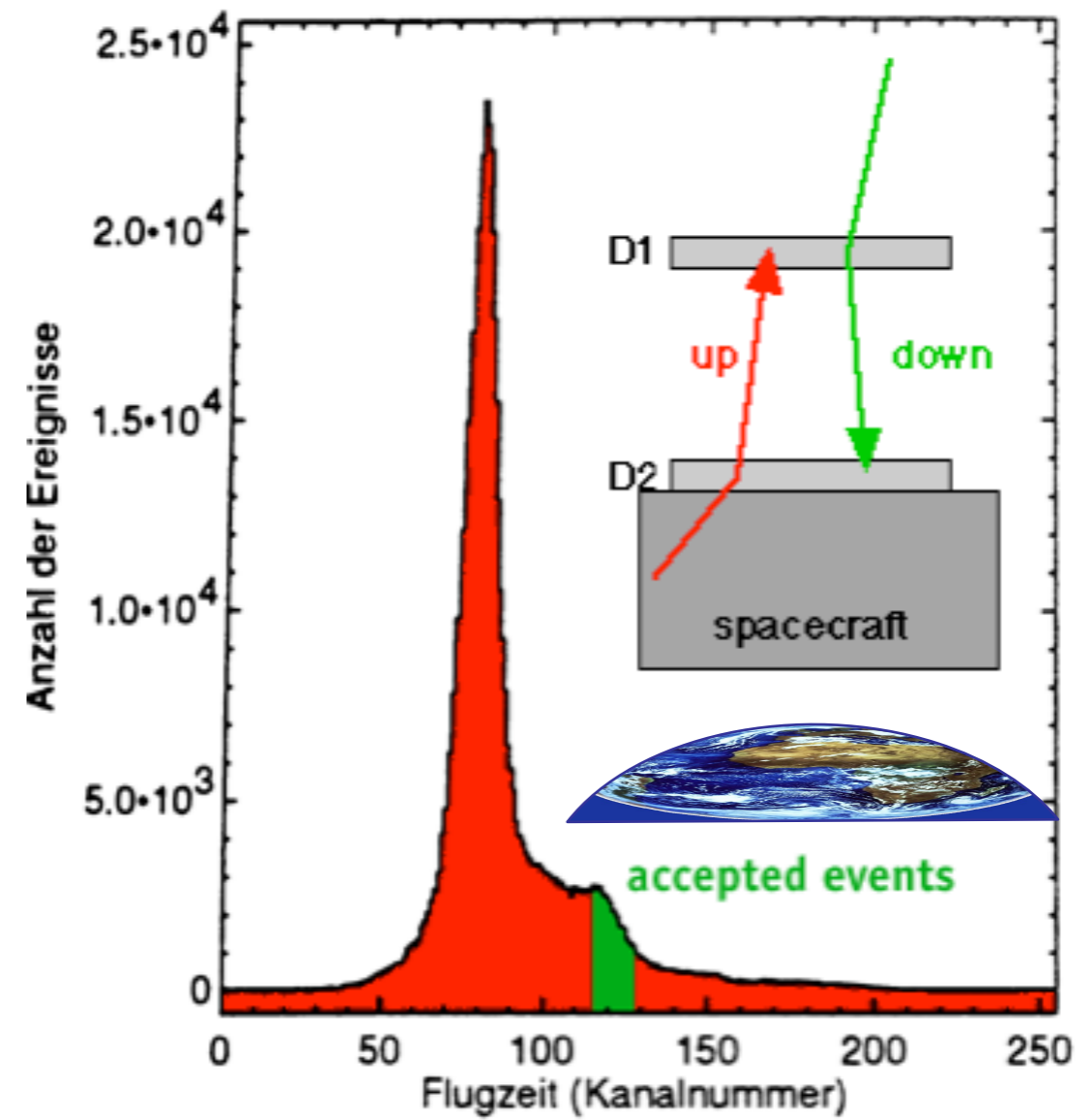
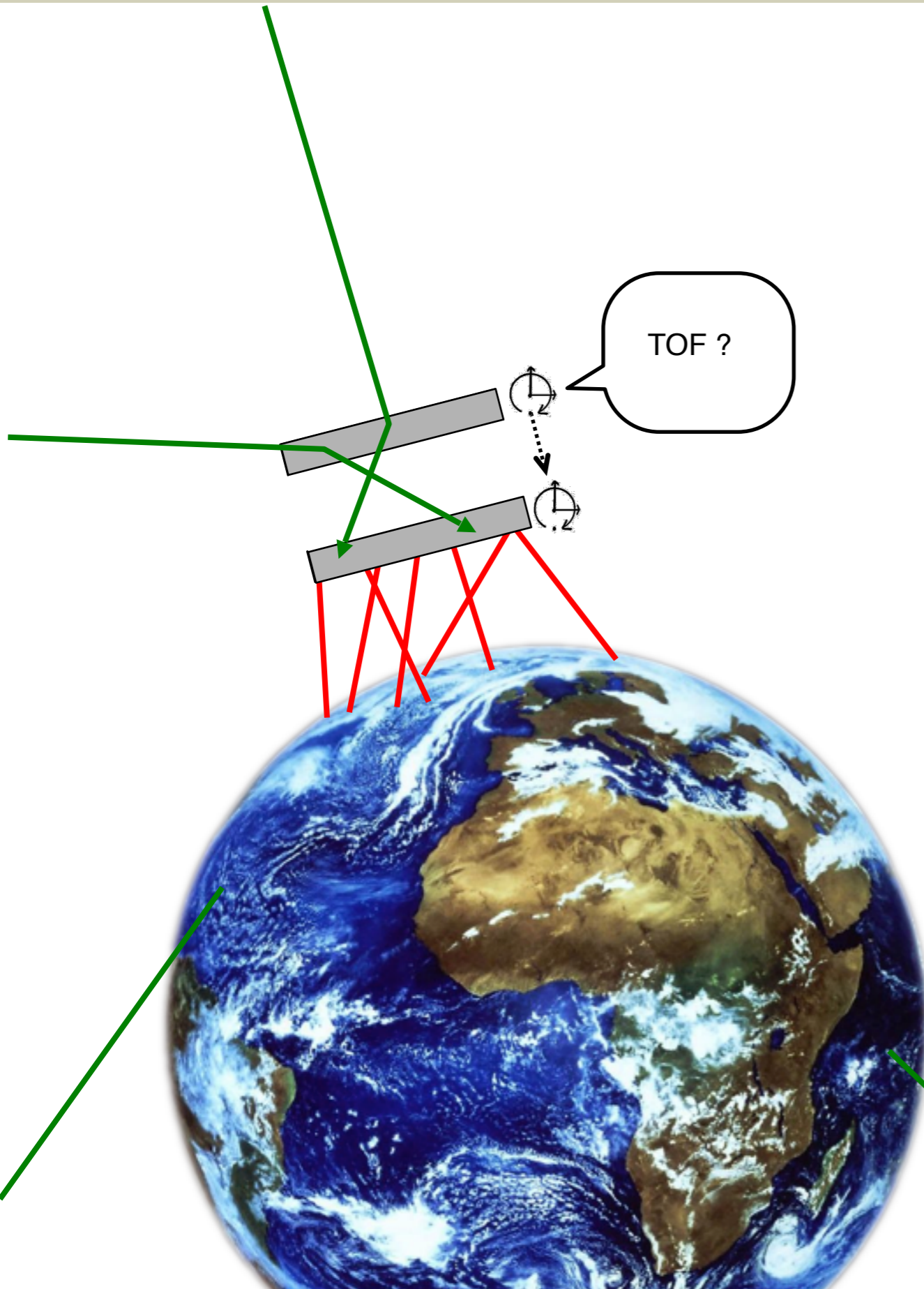
<- upward | downward ->  
**COMPTEL flight data**

channel width : 0.25 ns  
 "upward bkg" from spacecraft and Earth

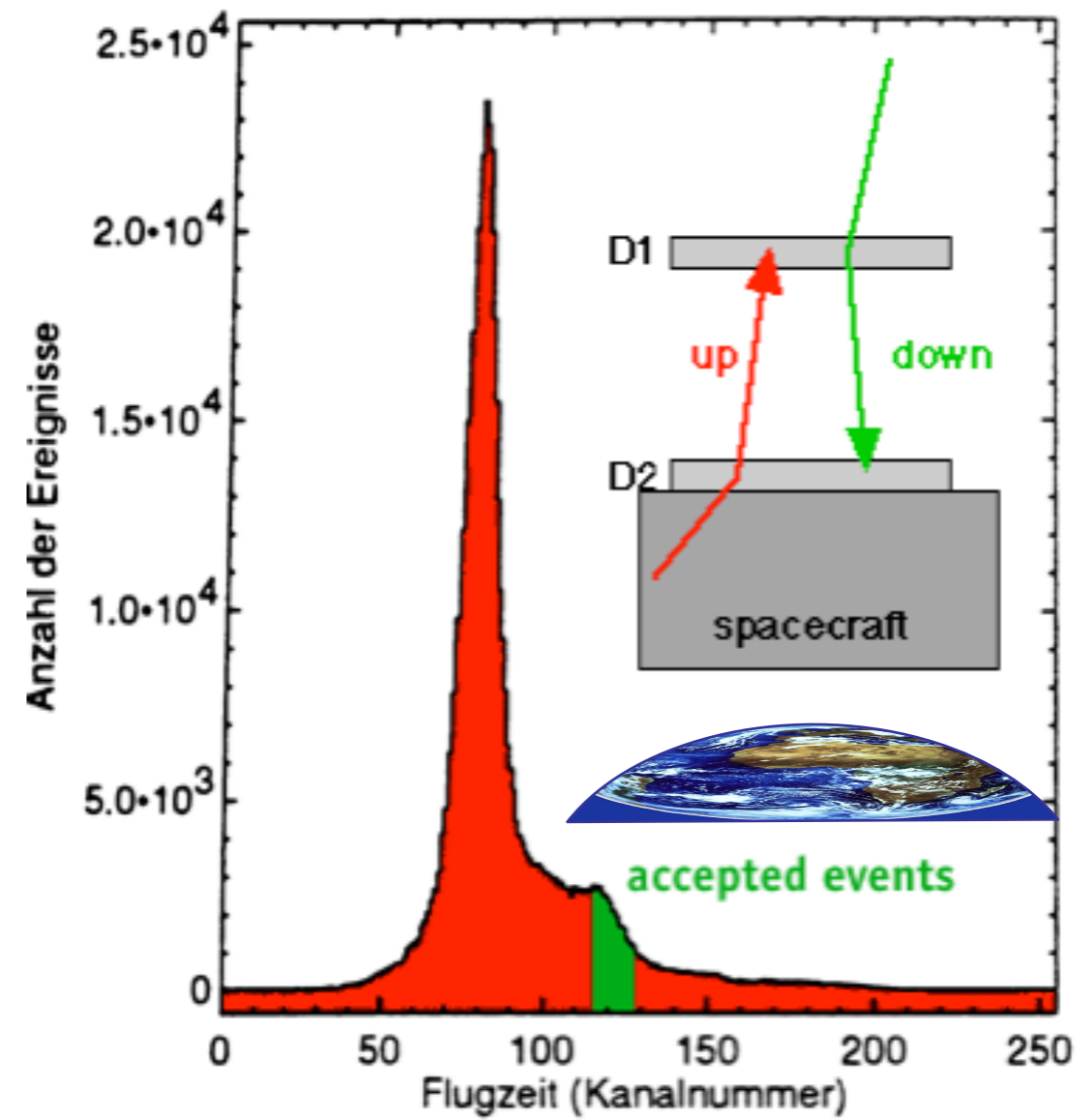
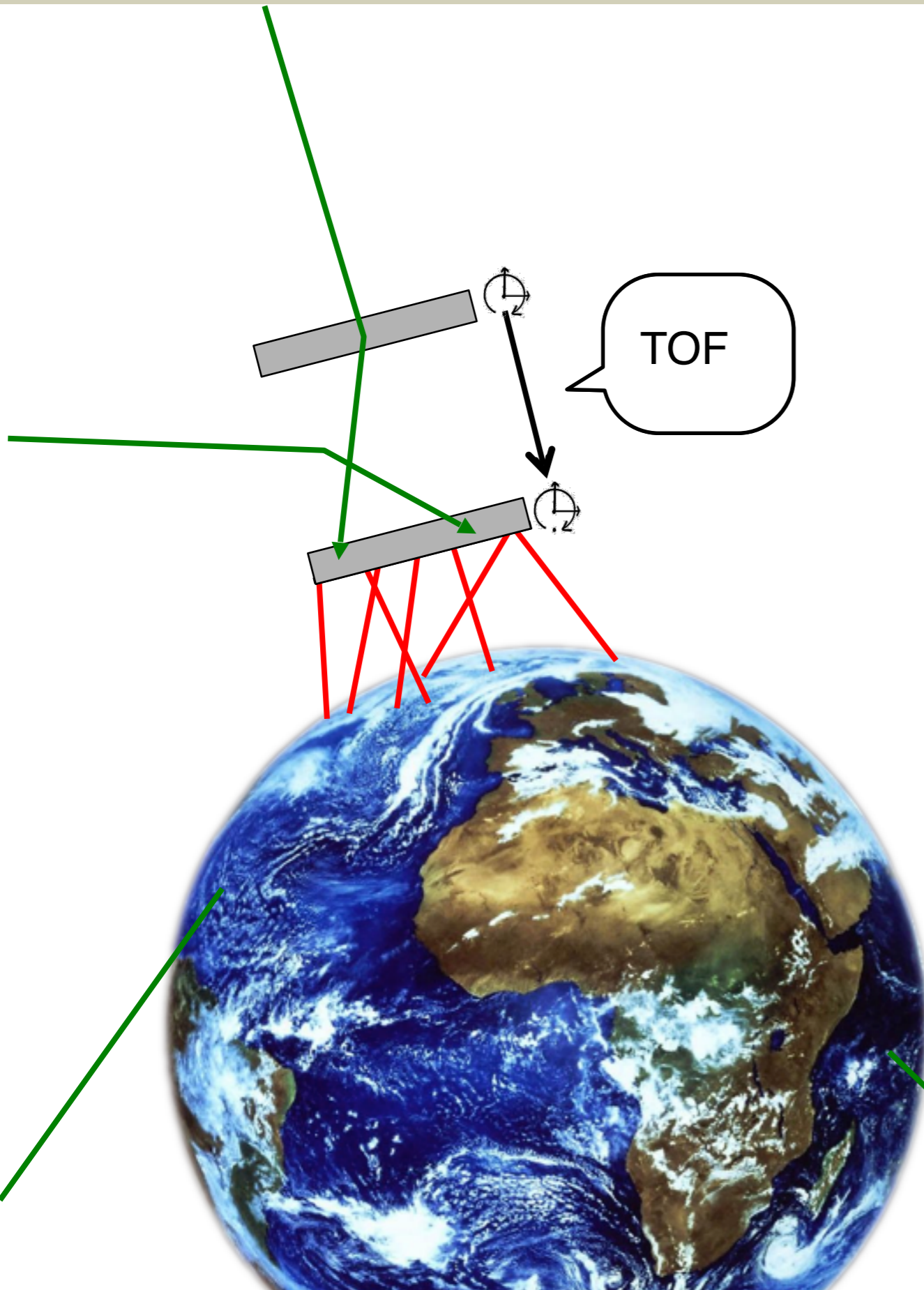
# option A : time-of-flight electronics



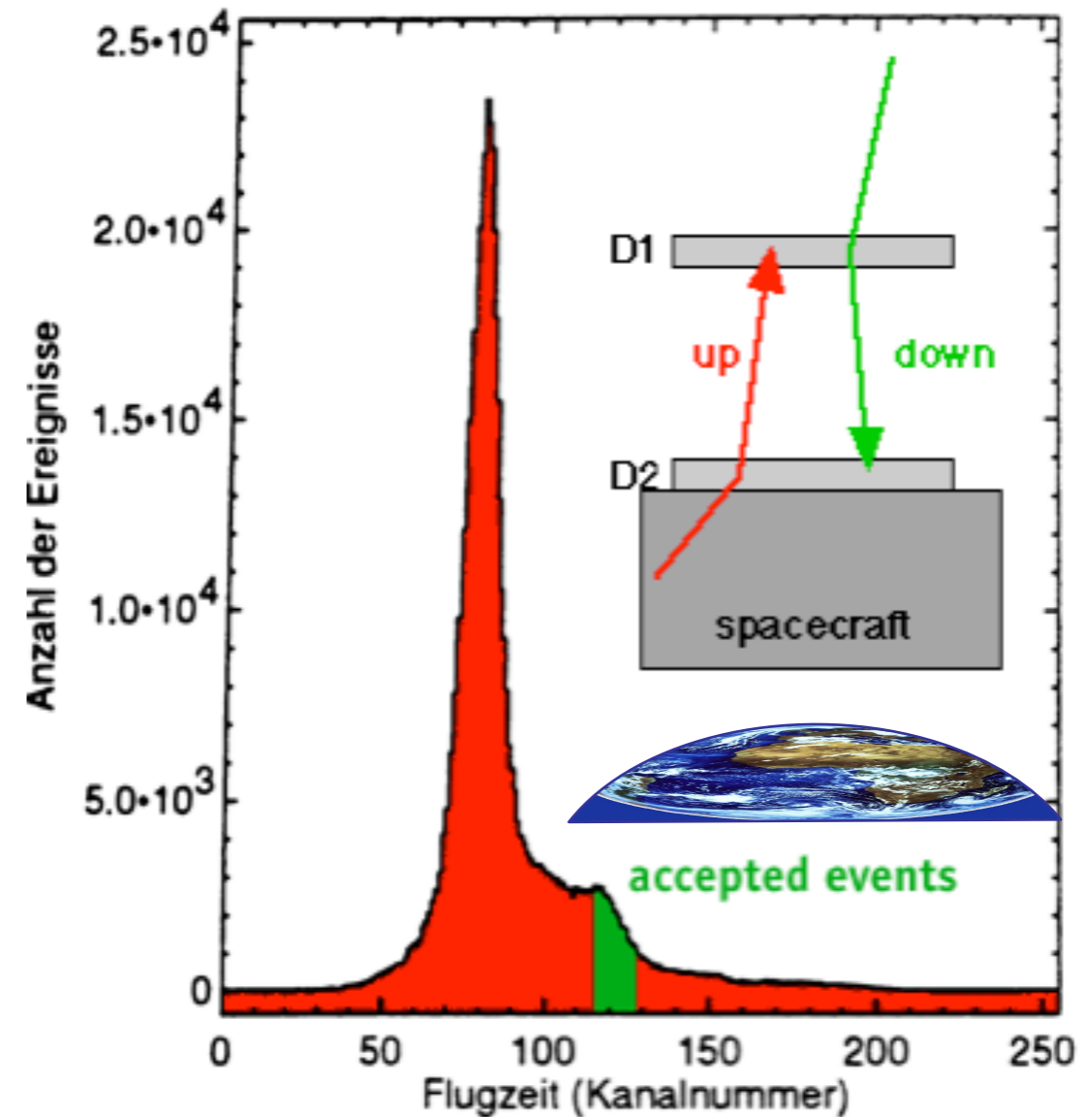
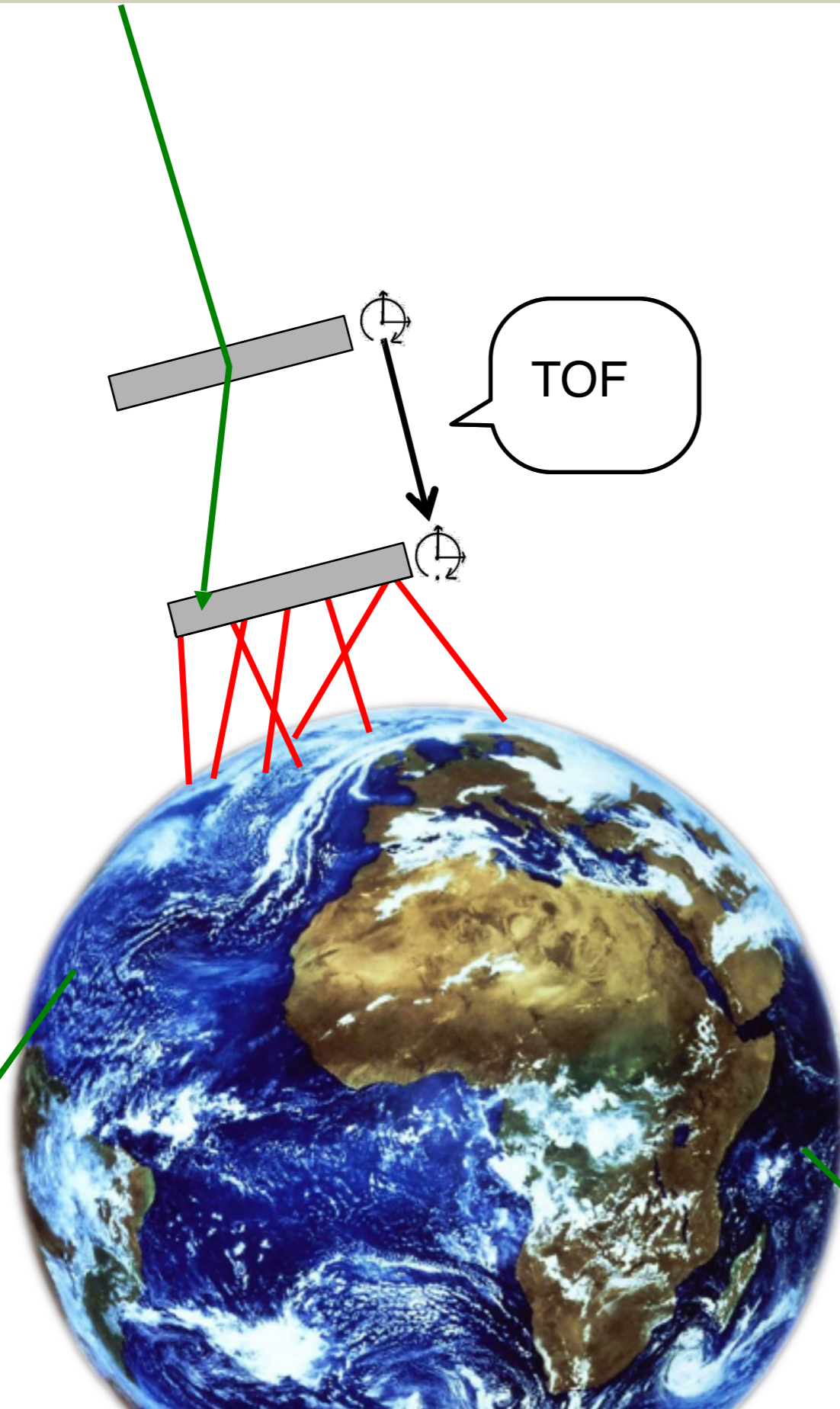
# option A : time-of-flight electronics



# option A : time-of-flight electronics

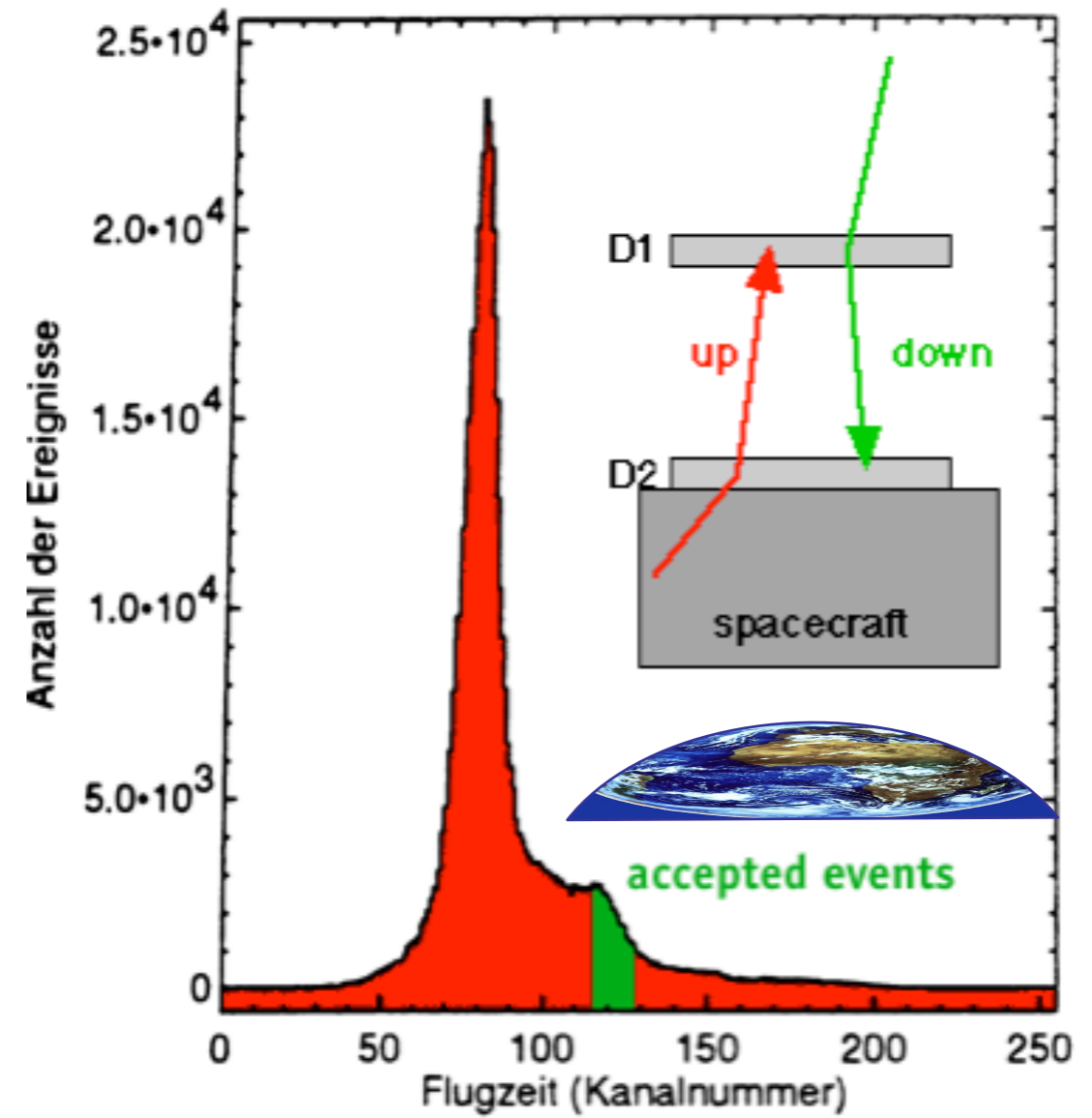
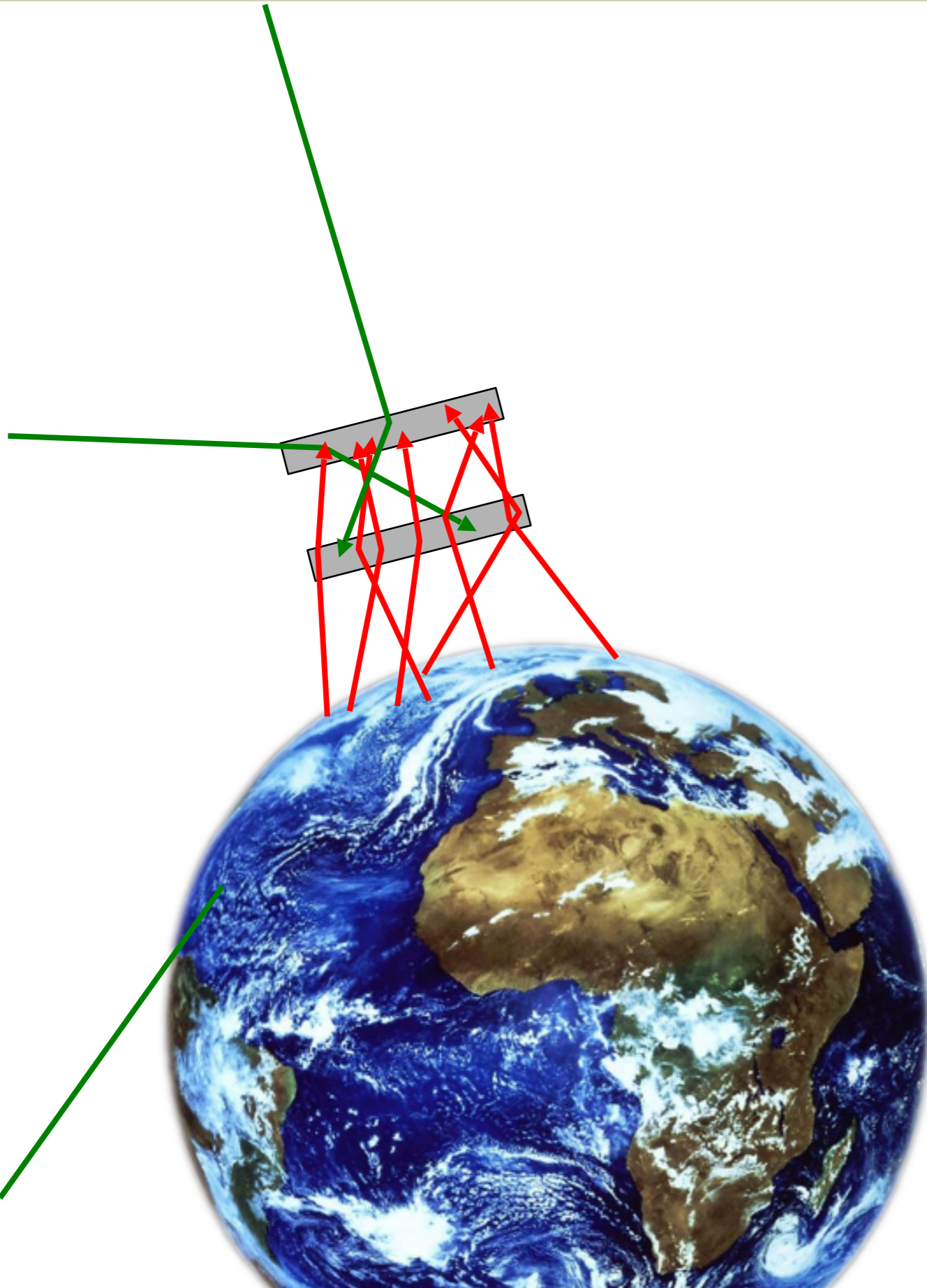


# option A : time-of-flight electronics

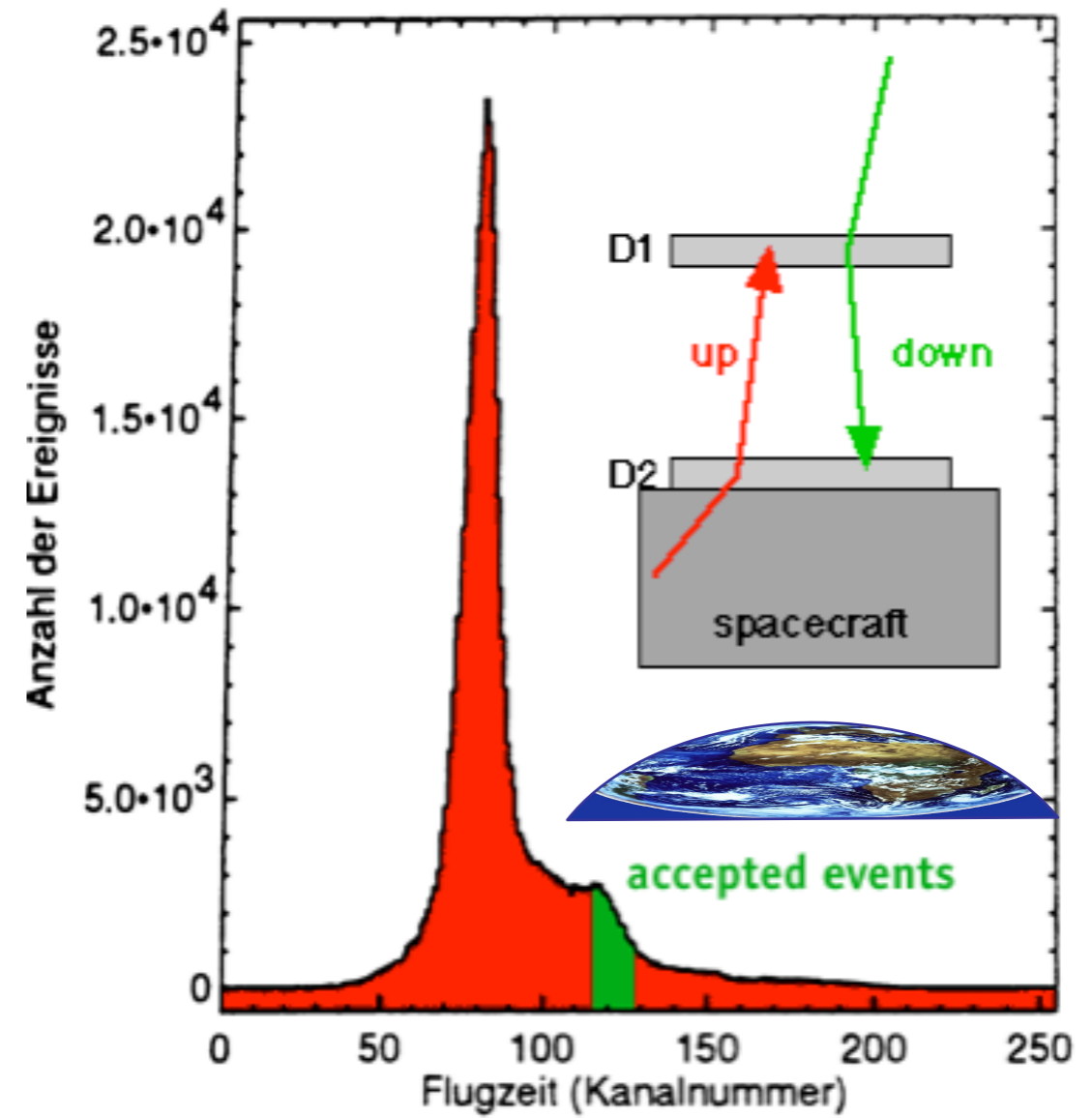
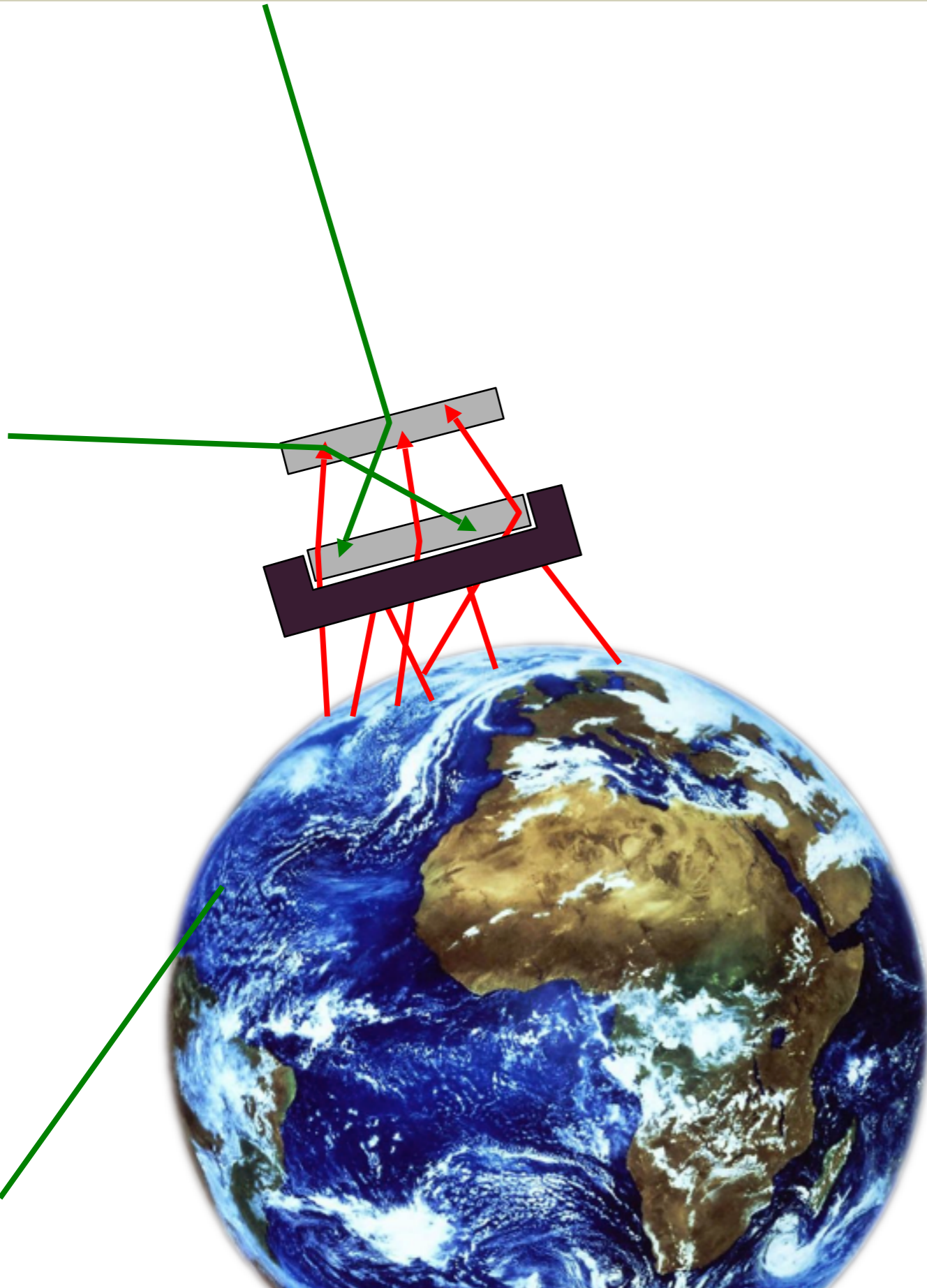


measuring TOFs requires  
long baselines between D1 and D2  
=> low probability for coincidence  
=> **low efficiencies** (few % at most)

# option B : anticoincidence shield



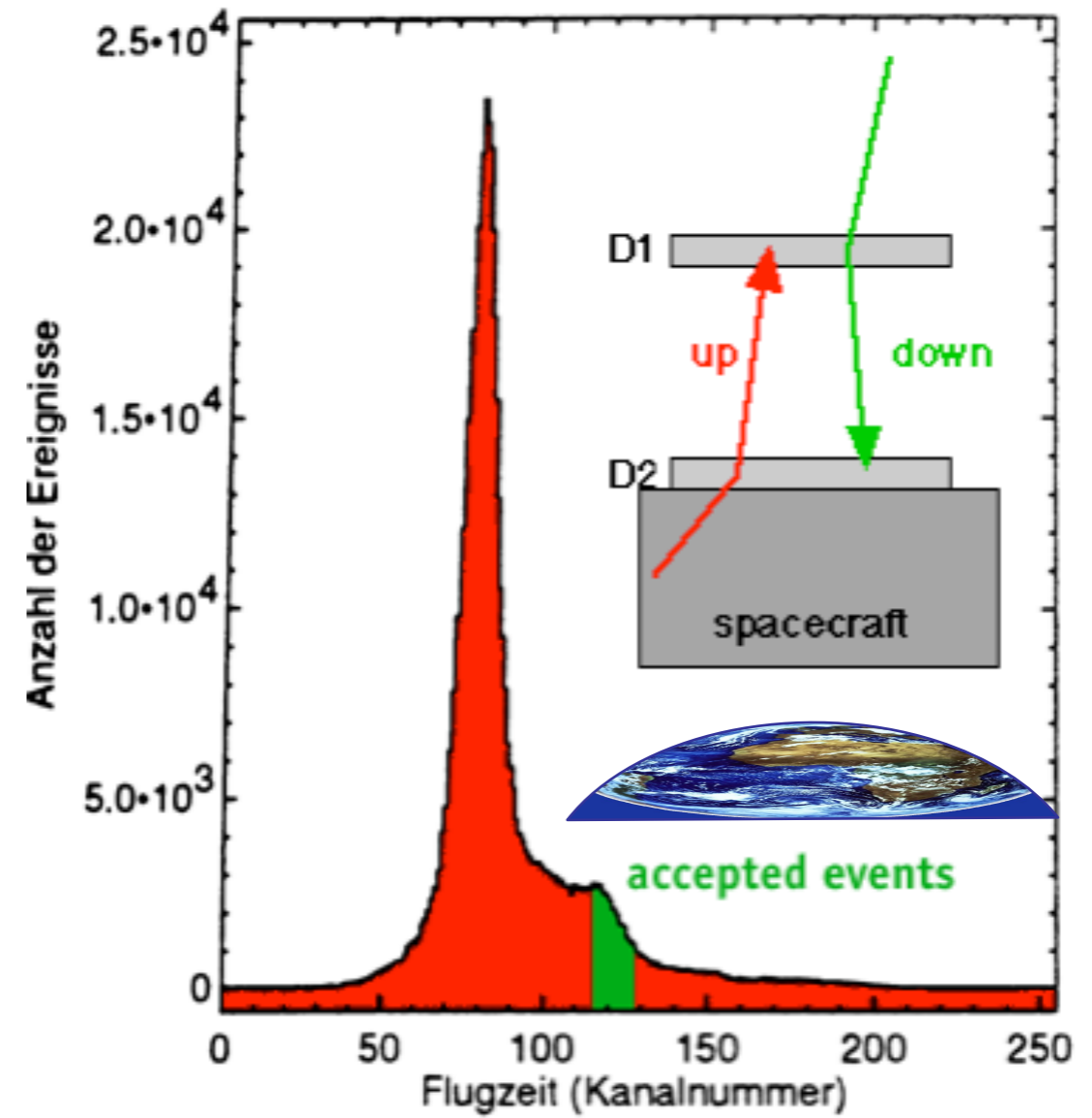
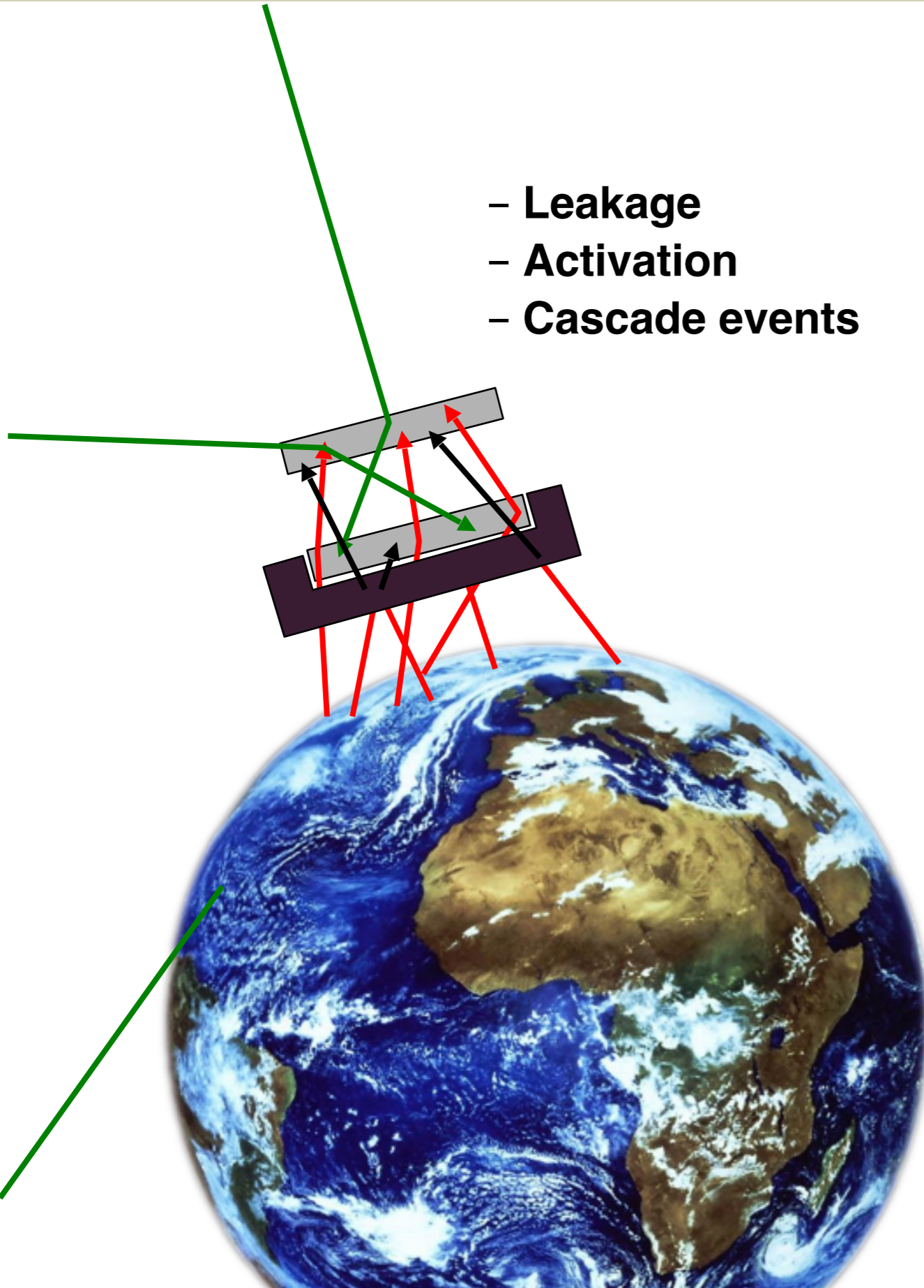
# option B : anticoincidence shield





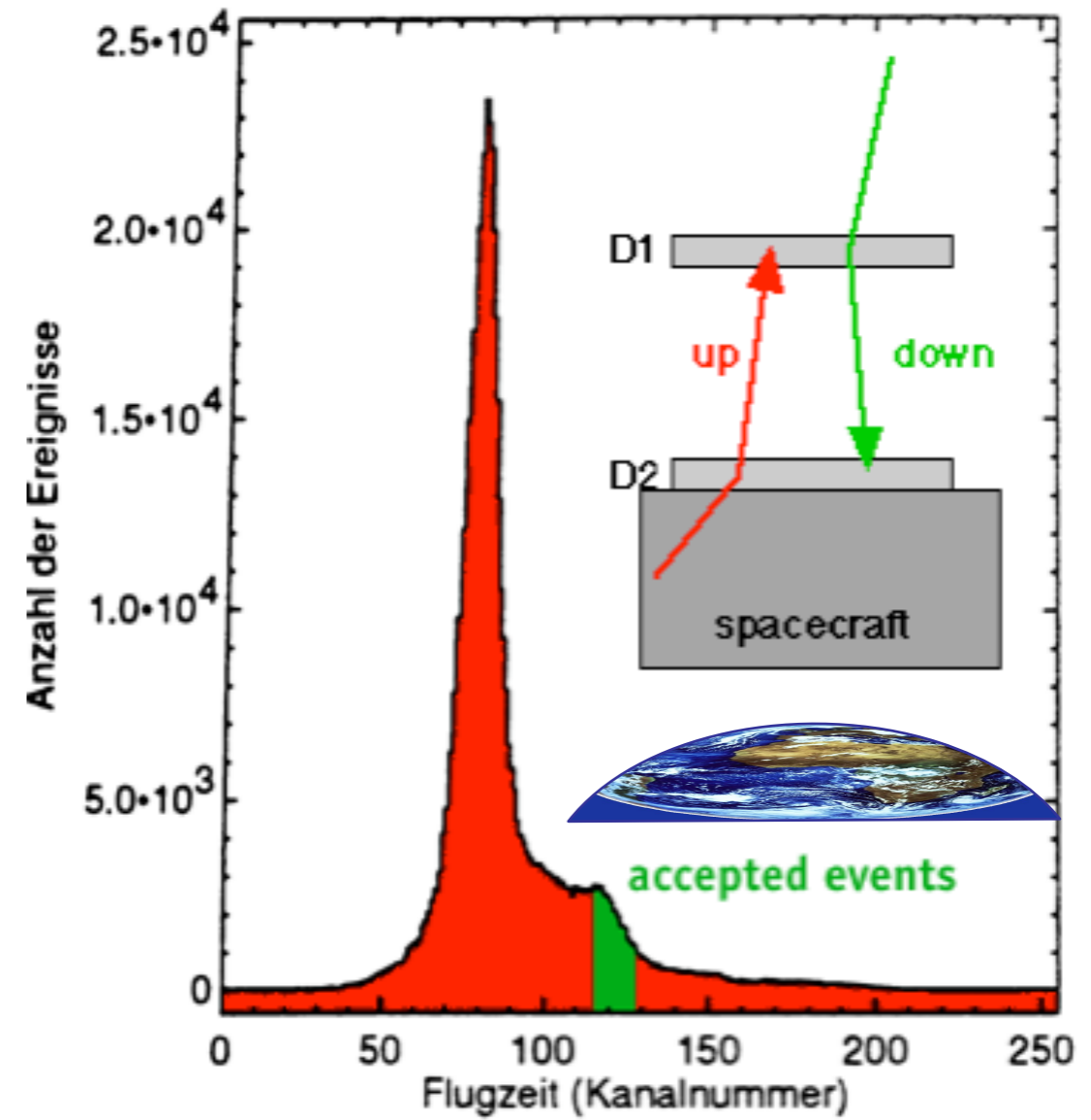
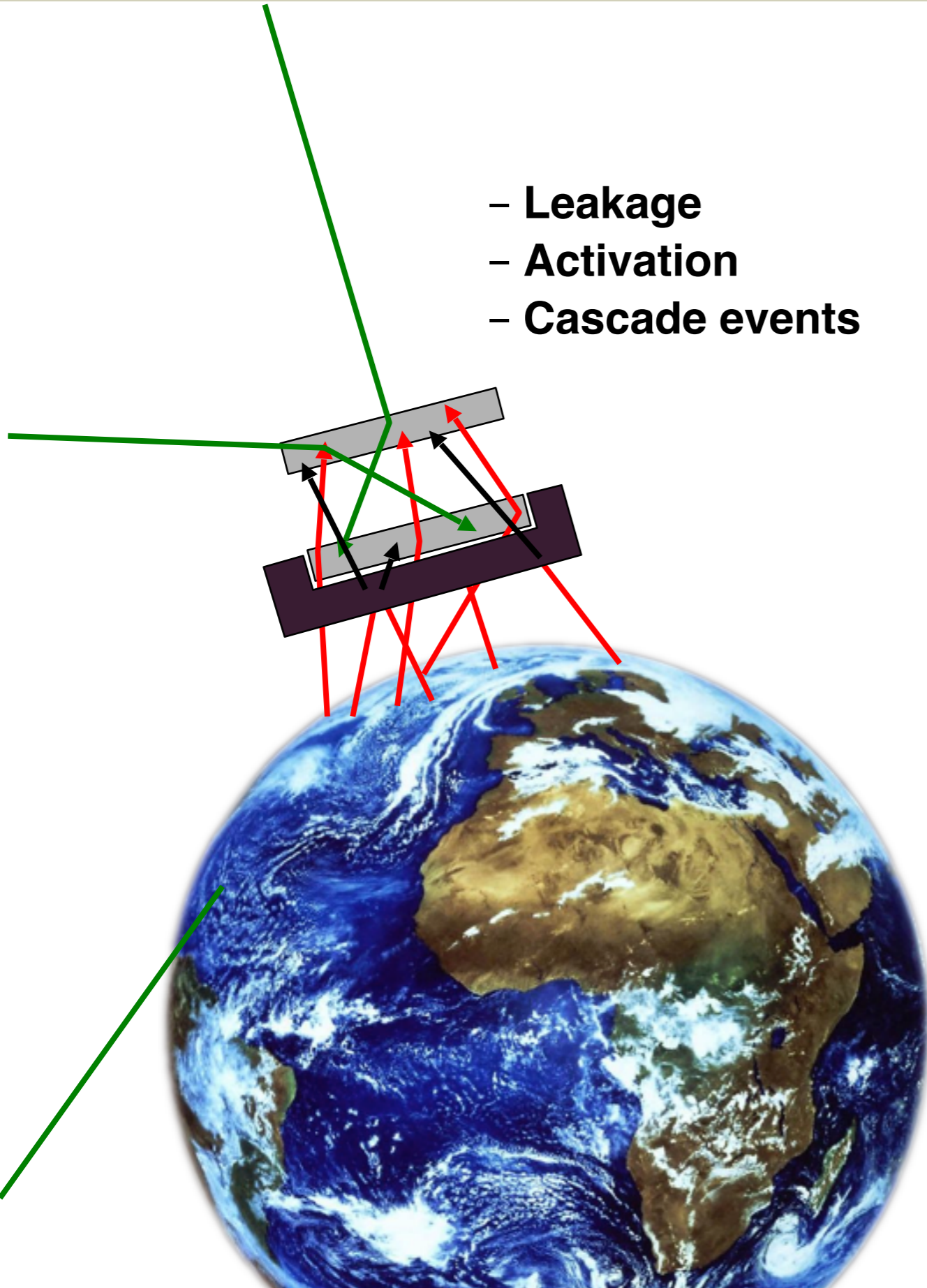
# option B : anticoincidence shield

- Leakage
- Activation
- Cascade events



# option B : anticoincidence shield

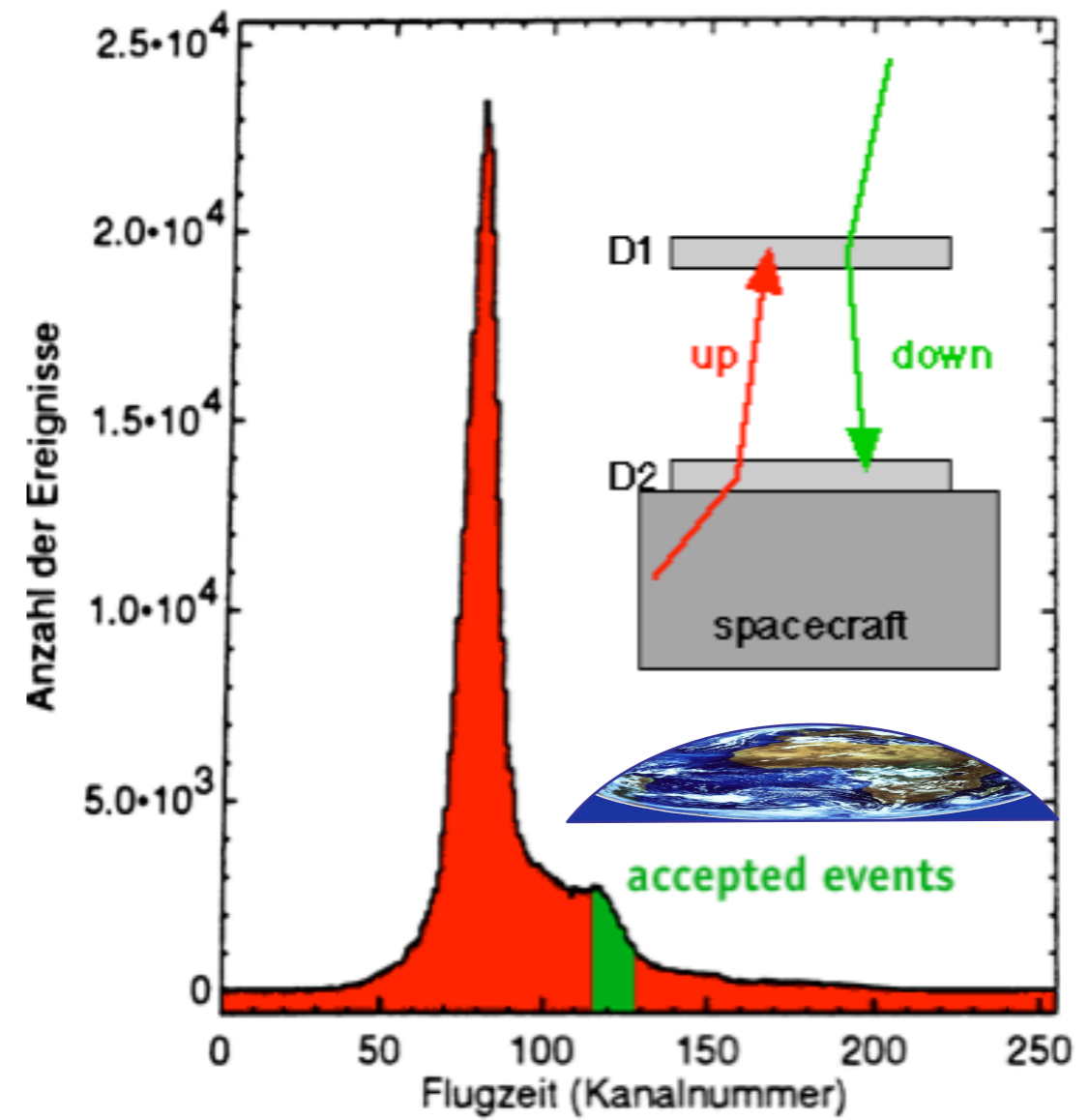
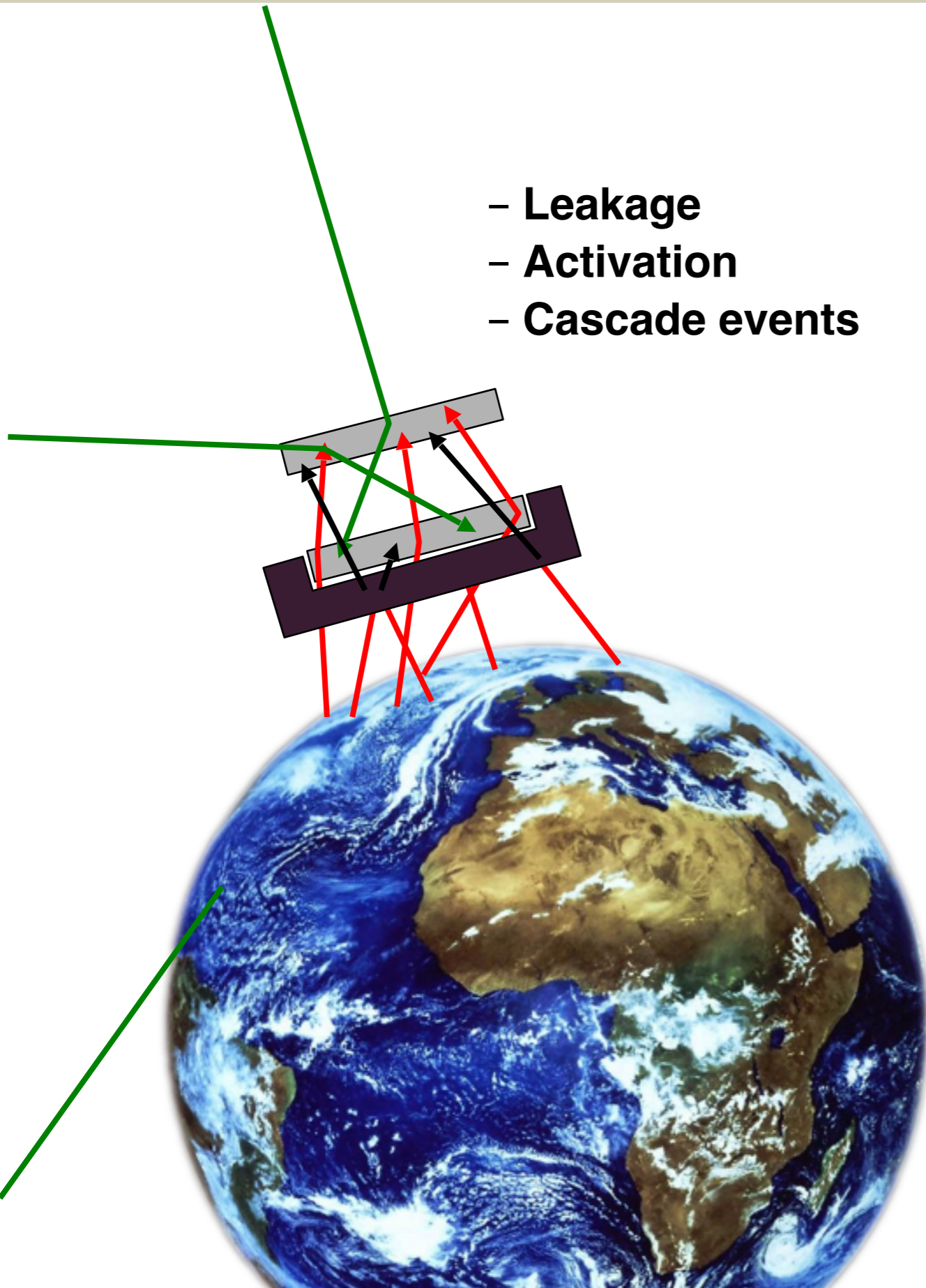
- Leakage
- Activation
- Cascade events



Compact solid state CTs

# option B : anticoincidence shield

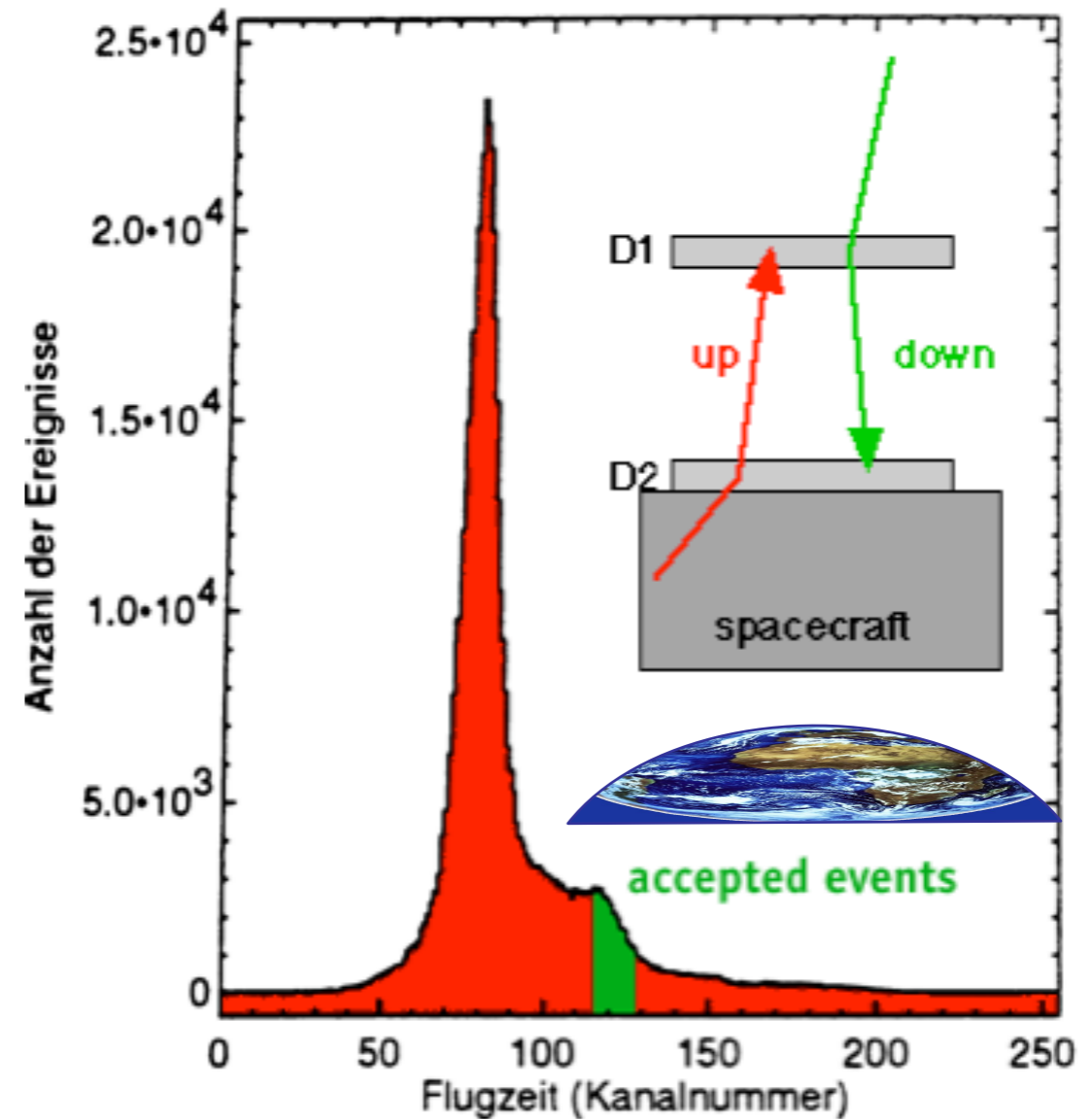
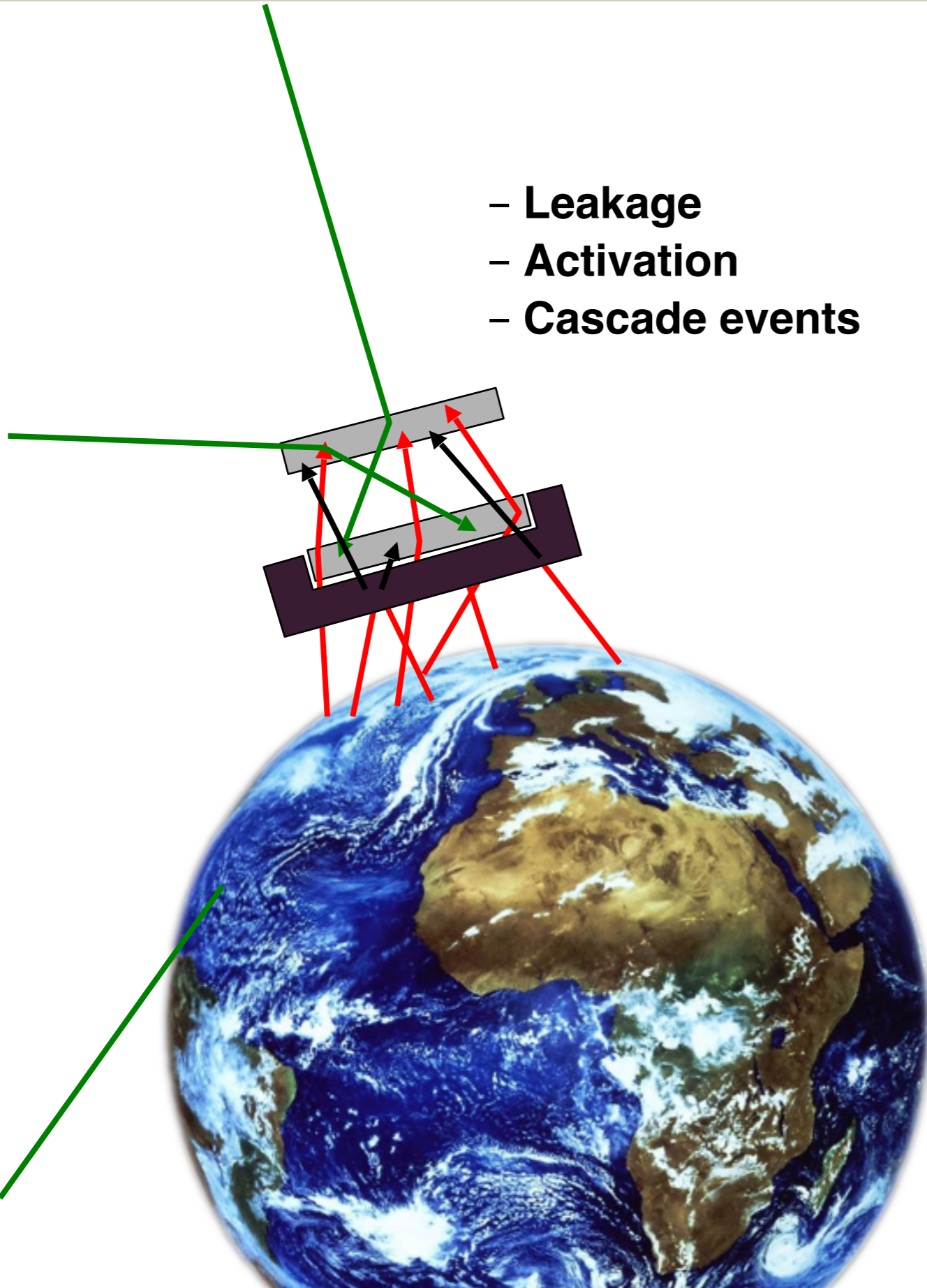
- Leakage
- Activation
- Cascade events



Compact solid state CTs  
Higher Compton efficiency

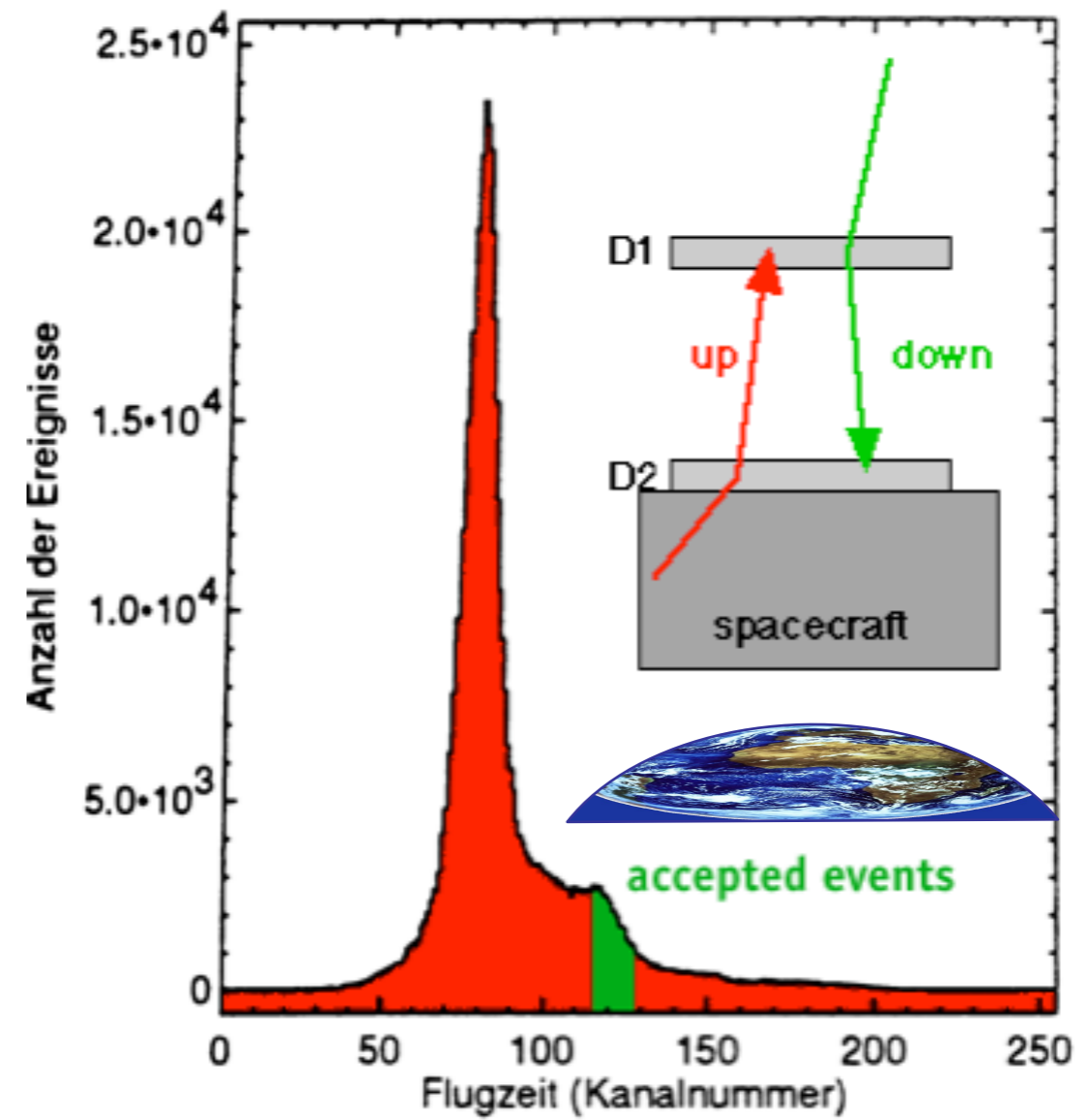
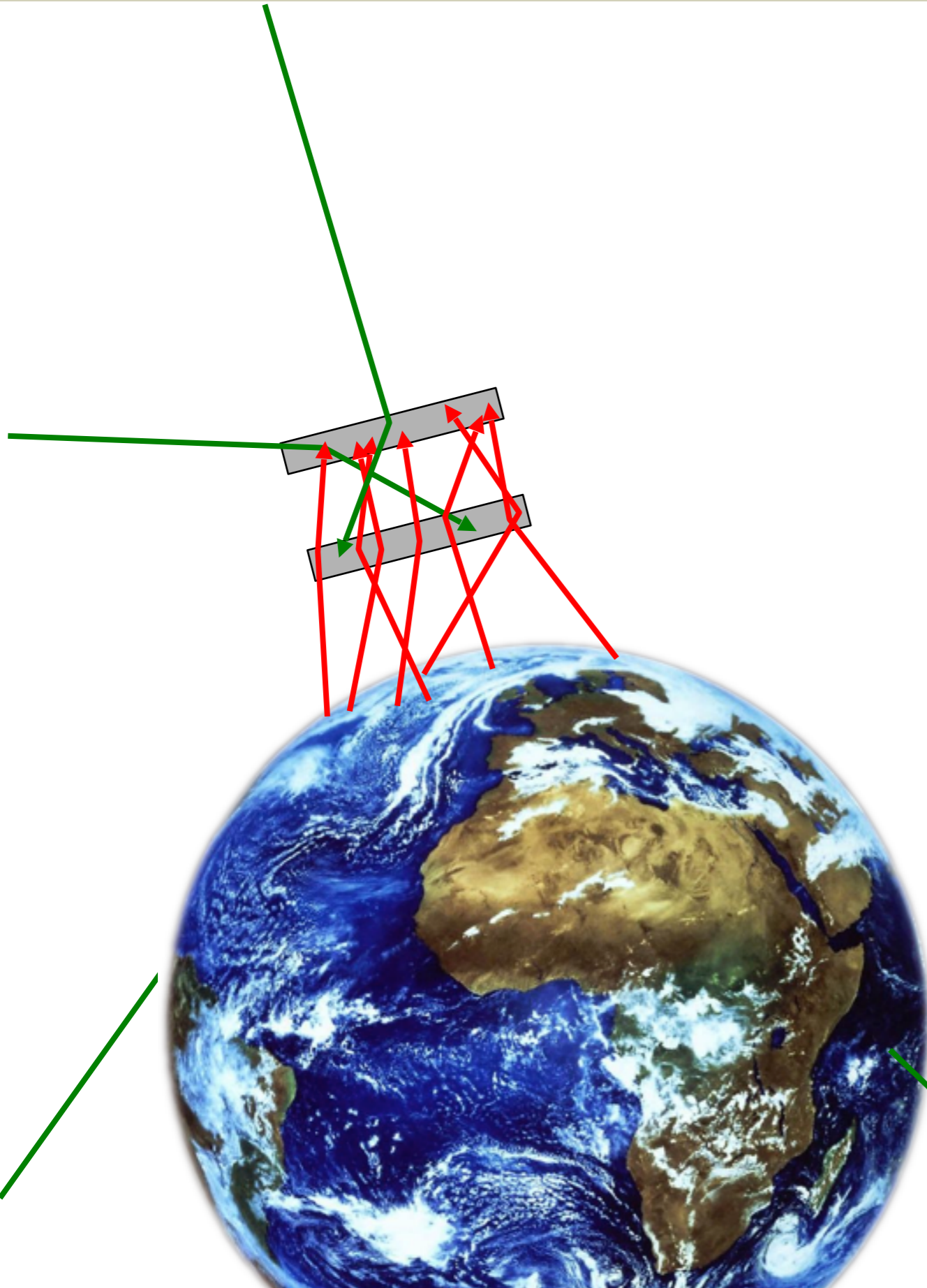
# option B : anticoincidence shield

- Leakage
- Activation
- Cascade events

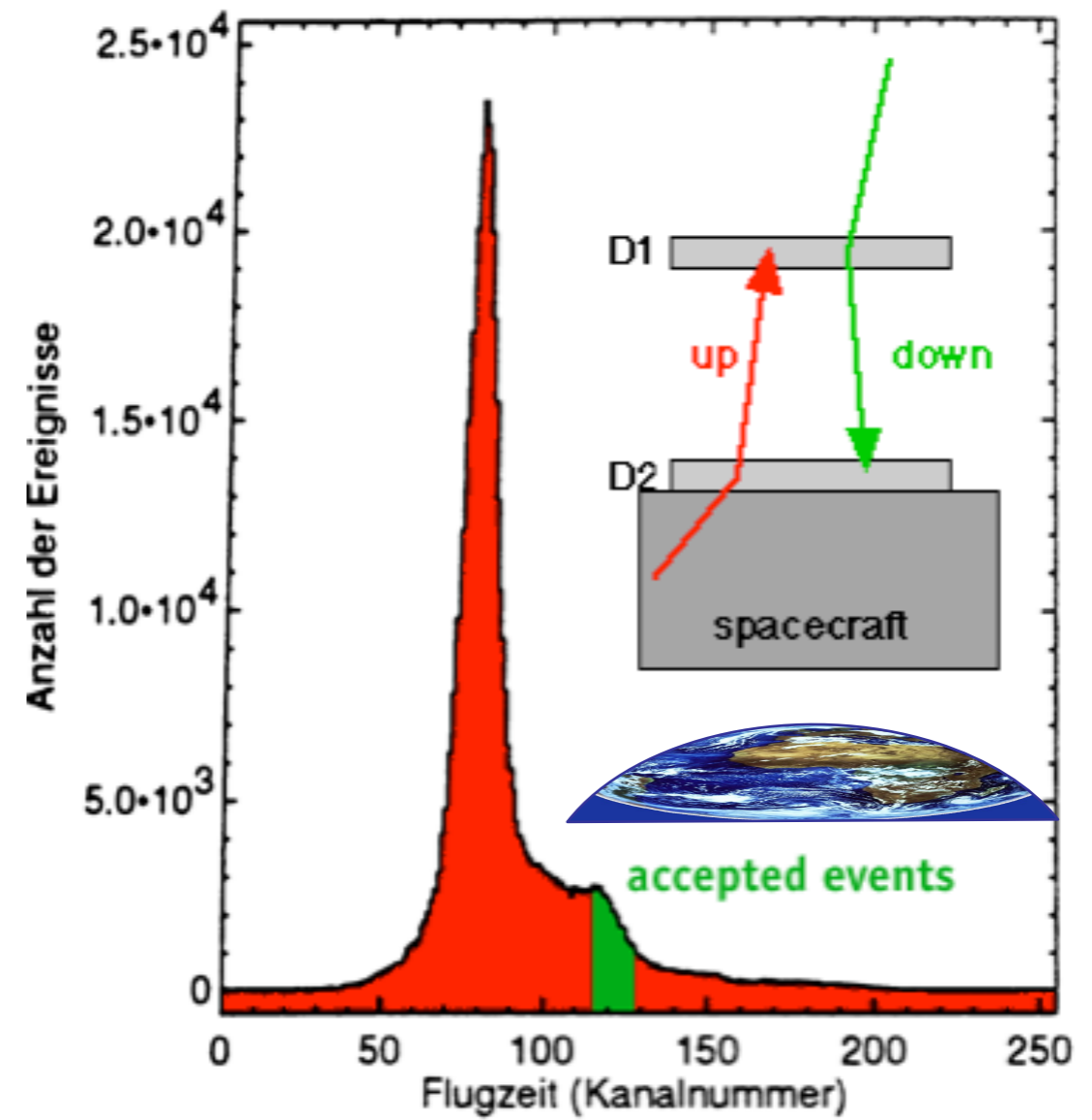
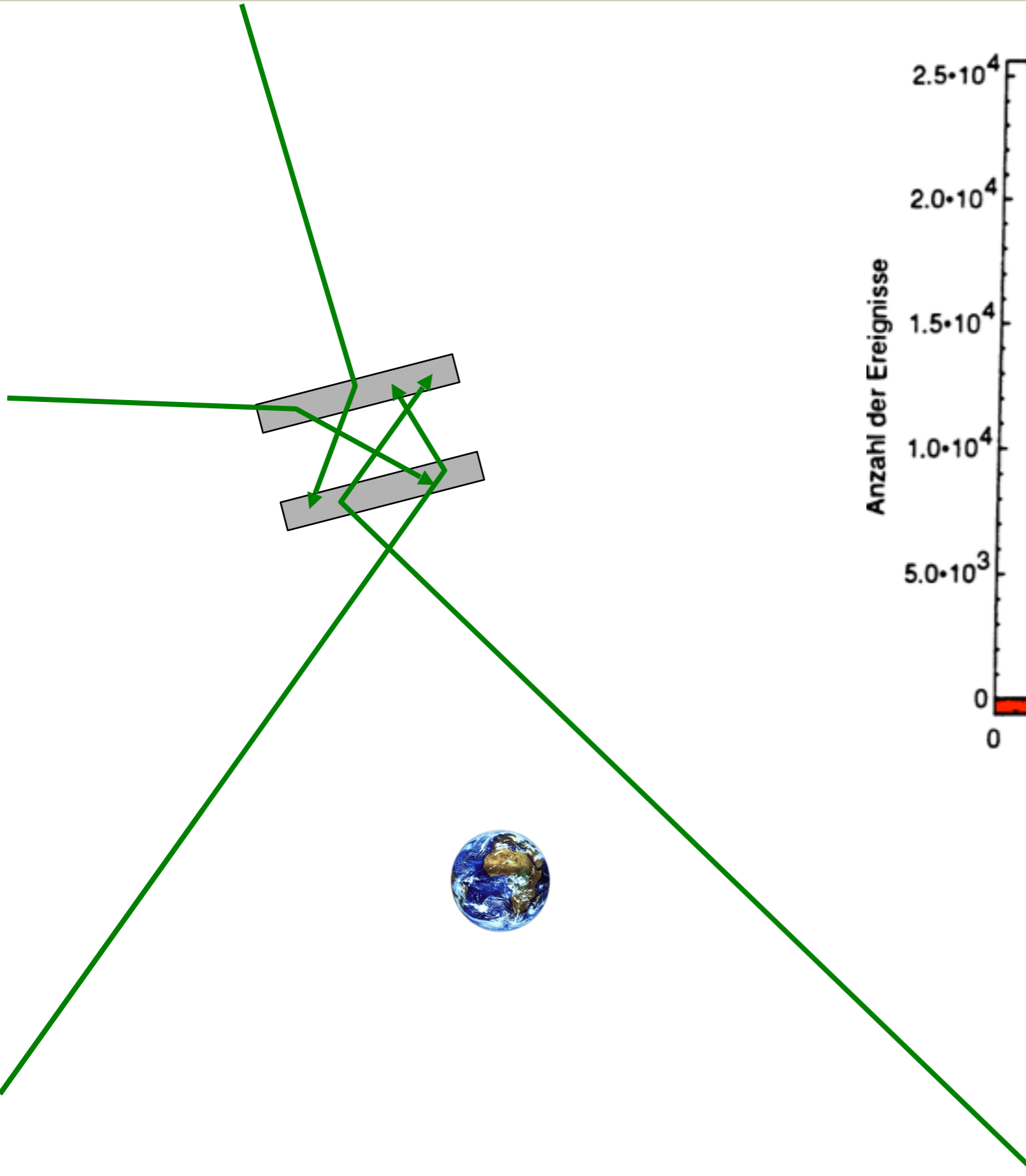


Compact solid state CTs  
Higher Compton efficiency  
They require **AC shields** often **more massive than the instrument itself**

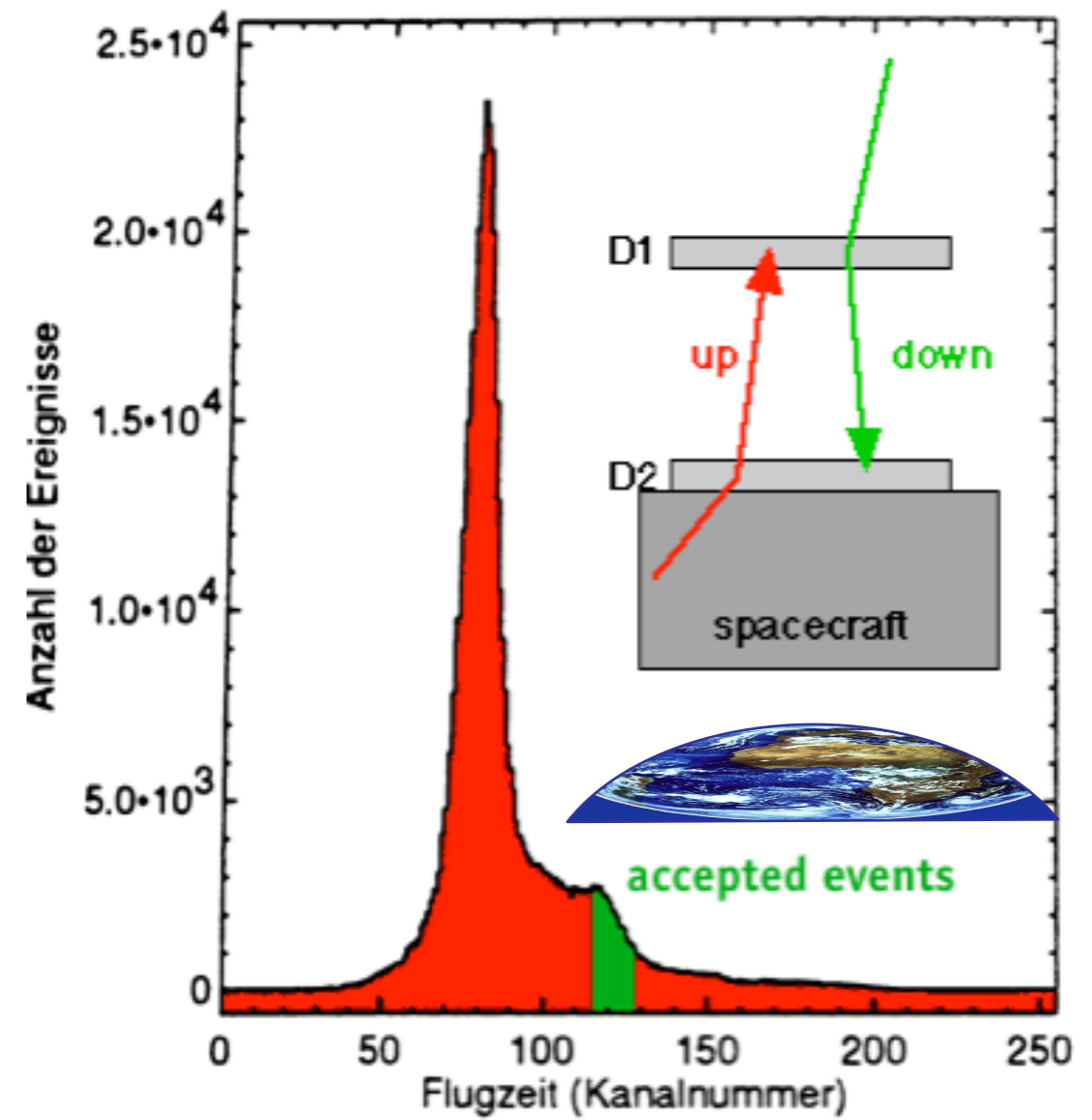
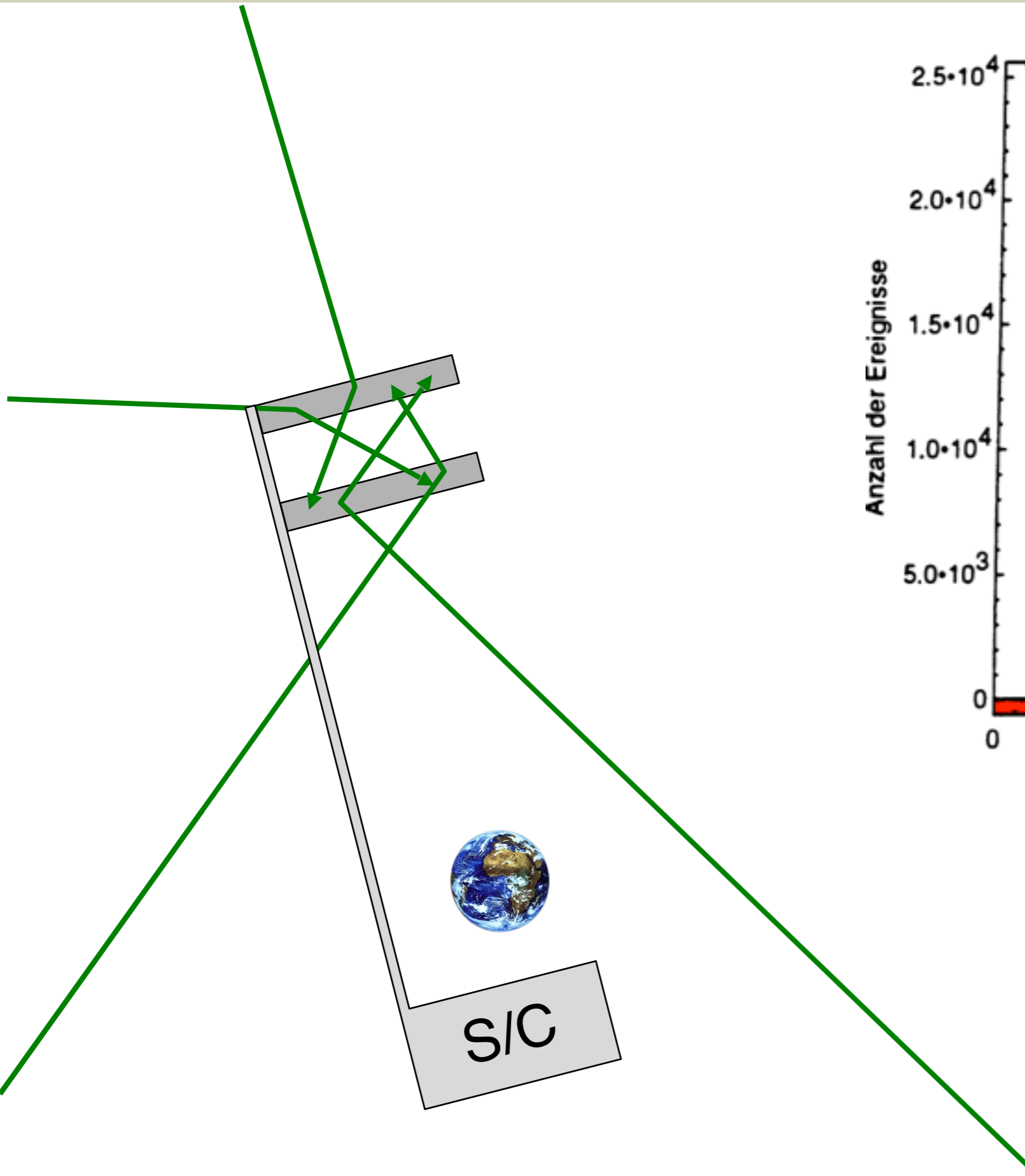
# option C : no bkg from external passive mass



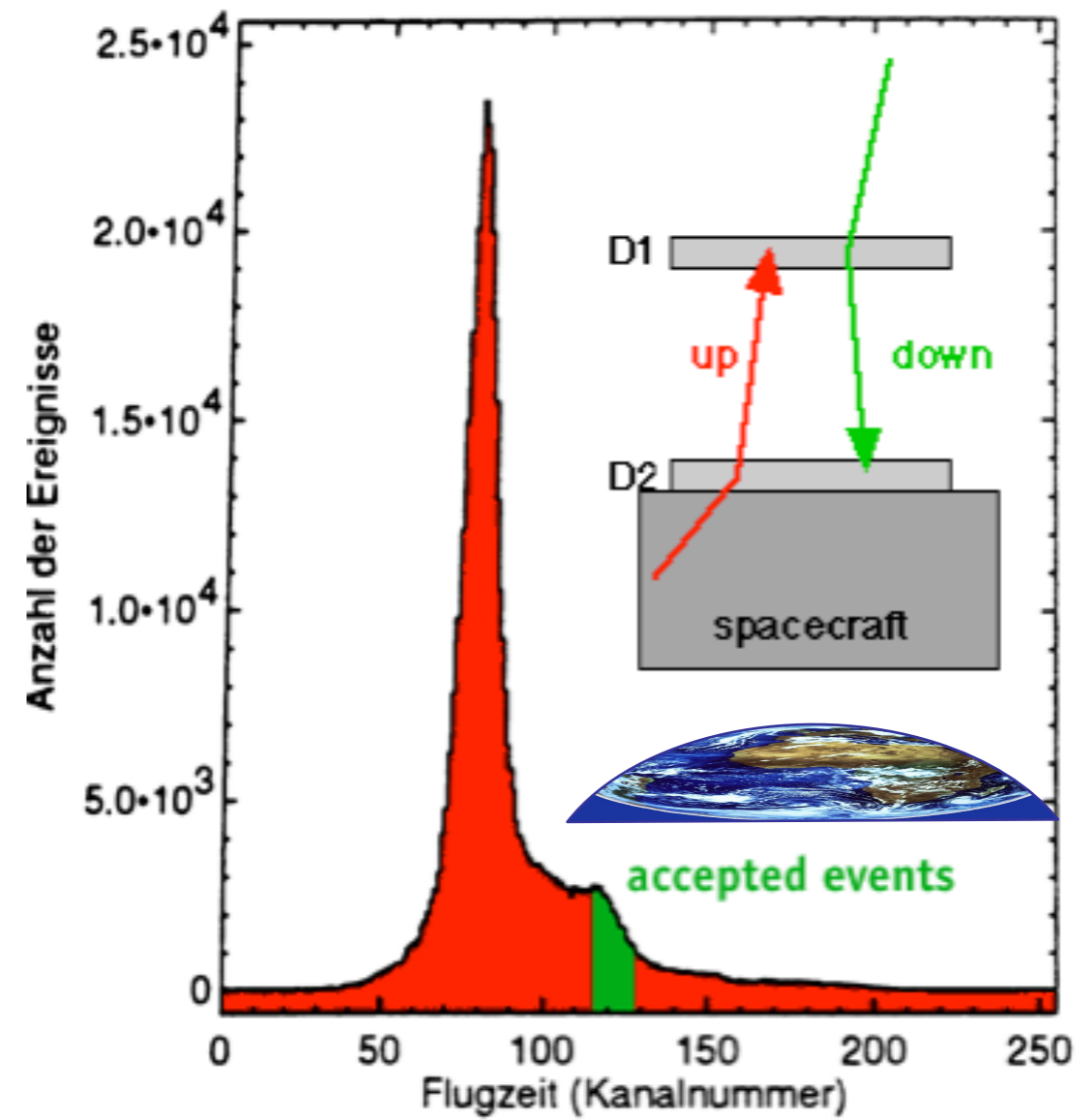
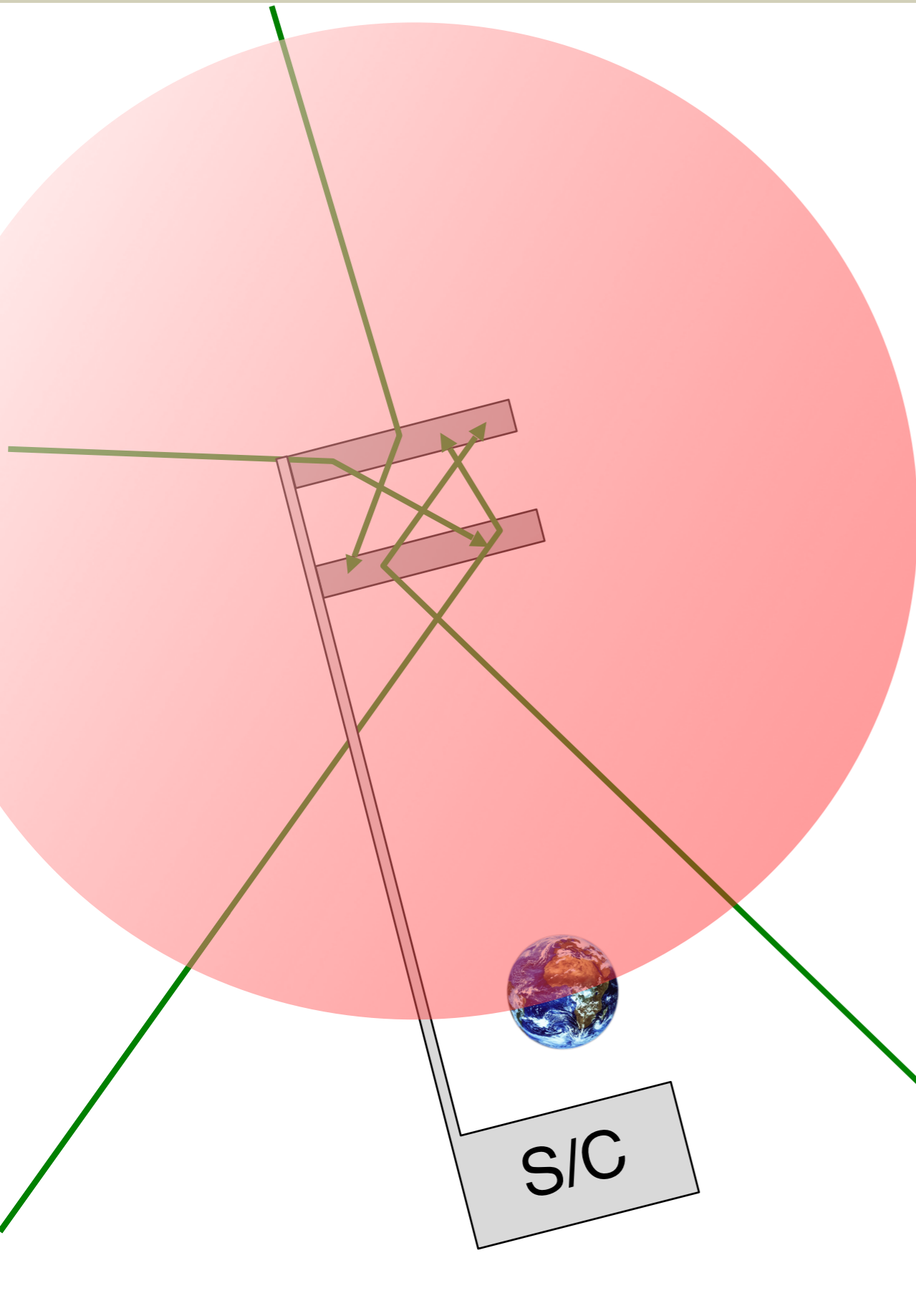
# option C : no bkg from external passive mass



# option C : no bkg from external passive mass



# option C : no bkg from external passive mass







# Which orbit? LEO or HEO?

- Low Earth Orbit

- Advantages

- Reduced CR background from geomagnetic shielding
  - Reduced prompt CR contamination
  - Reduced instrument and s/c activation
  - Note: want low inclination (i.e. near 0 deg)
    - » Maximizes geomagnetic screening, i.e. minimizes CR-induced background
    - » Improves livetime by avoiding SAA
- Increased payload mass at lower launch cost

- Disadvantages

- Strong atmospheric  $\gamma$ -ray background
- Earth occults  $\sim 1/3$  of the sky

- High Earth Orbit

- Advantages

- Reduced atmospheric  $\gamma$ -ray background
- Increased FOV (nearly  $4\pi$  possible)

- Disadvantages

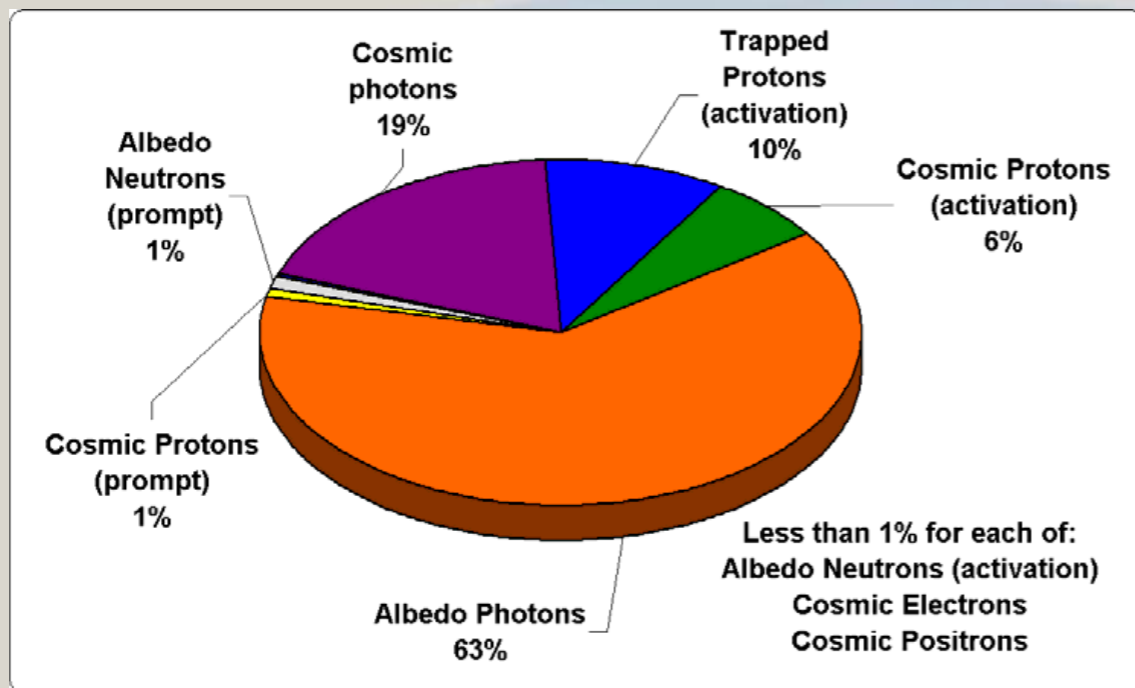
- Increased CR background
  - Increased prompt CR contamination
    - » High trigger rate, data volume
    - » More CRs to identify and reject
  - Increased instrument and s/c activation
- Decreased payload mass and/or higher launch cost

What's the right choice?



# Example trade study for Si ACT

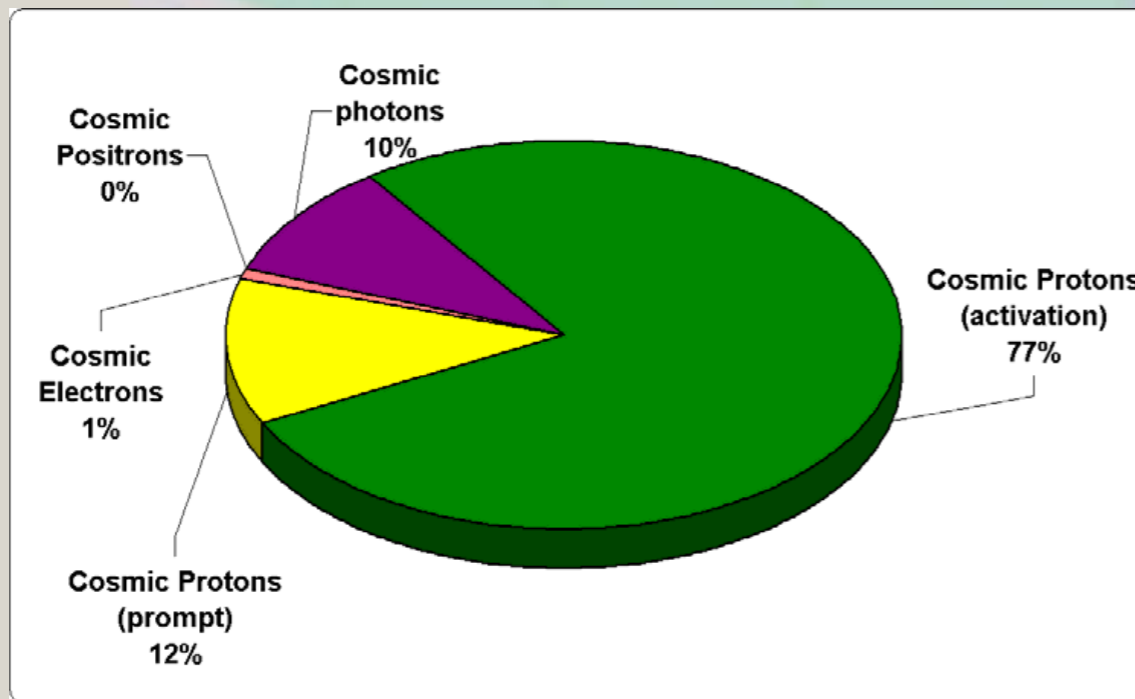
## Low Earth Orbit vs. High Earth Orbit Background



**LEO: 550 km, 8° inclination**  
 Horizon cut 92.5°  
 Contributions for: E = 847 +/- 22.75 keV  
 ARM radius 1°

3σ sensitivity over 10<sup>6</sup> sec,  
 at 3% FWHM brdn. 847keV line,  
 on-axis plain wave:

$$2.37 \cdot 10^{-6} \gamma \cdot \text{cm}^{-2} \cdot \text{s}^{-1}$$



**HEO: 40,000 km**

Horizon cut 10°  
 Contributions for: E = 847 +/- 22.75 keV  
 ARM radius 1°

3σ sensitivity over 10<sup>6</sup> sec,  
 at 3% FWHM brdn. 847keV line,  
 on-axis plain wave:

$$2.97 \cdot 10^{-6} \gamma \cdot \text{cm}^{-2} \cdot \text{s}^{-1}$$



# Backgrounds

- Beware of self-activity
  - Are lanthanum halides good choices for Compton calorimeter?
  - $\text{LaBr}_3$ ,  $\text{LaCl}_3$ 
    - Fast scintillator, good energy resolution ( $\sim 4\%$  at 1 MeV), high stopping power
    - Hot

## Nal vs LaBr Compton Calorimeter

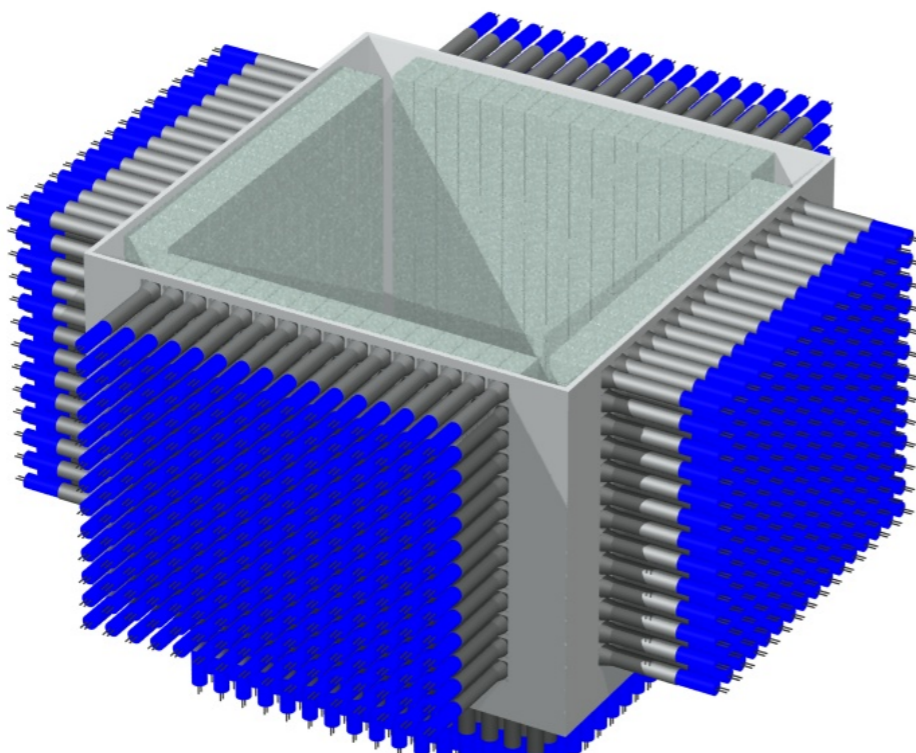
Bernard Philips

Code 7654

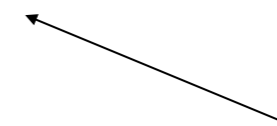
Naval Research Laboratory

- Study performed for ACT, GRIPS

## Calorimeter



- NaI crystals are standard parts
- Frame will also support stack of silicon
- Need slots for cables from/to silicon detectors
- Area inside calorimeter  $\sim 45 \text{ cm} \times 45 \text{ cm}$



Note size



# Self activity or induced activation

- Beta-gamma decays look just like signal
  - e.g. La self-activity for large instrument creates many kHz of nasty bkg

## Lanthanum Activation

- Lanthanum is 99.91%  $^{139}\text{La}$ , and 0.09%  $^{138}\text{La}$
- $^{138}\text{La}$  decays with 2 different decay schemes:- 788.7 keV gamma
  - 1438.8 keV gamma and a beta with 205 keV endpoint
- The activity is 1.8 Bq/cm<sup>3</sup> for LaCl<sub>3</sub> and 1.62 Bq/cm<sup>3</sup> for LaBr<sub>3</sub>
- For 5 cm thickness, have ~30 000 cm<sup>3</sup> calorimeter.
- ~50 000 Bq of activity within the instrument for LaBr<sub>3</sub>!

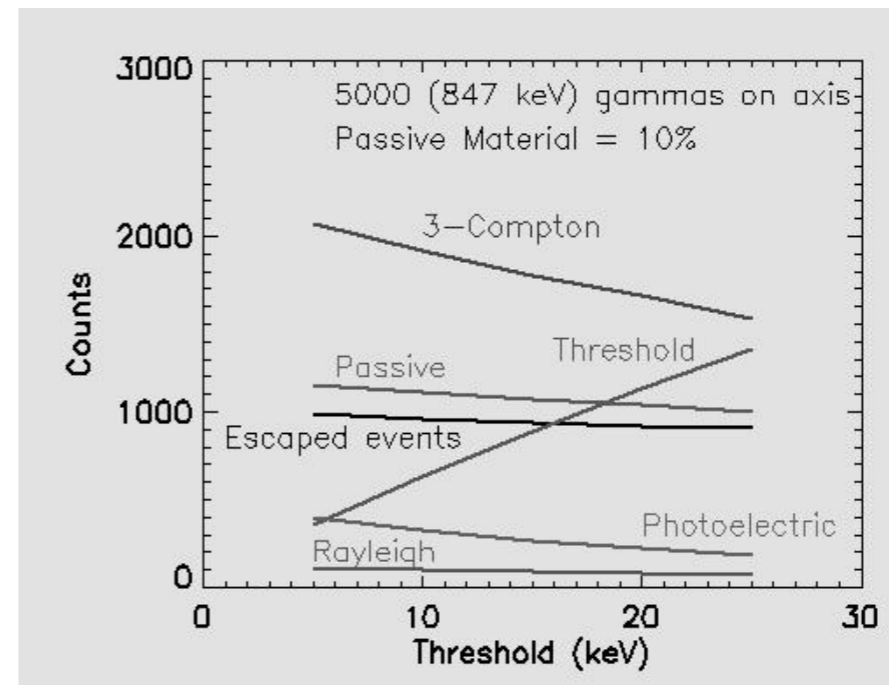
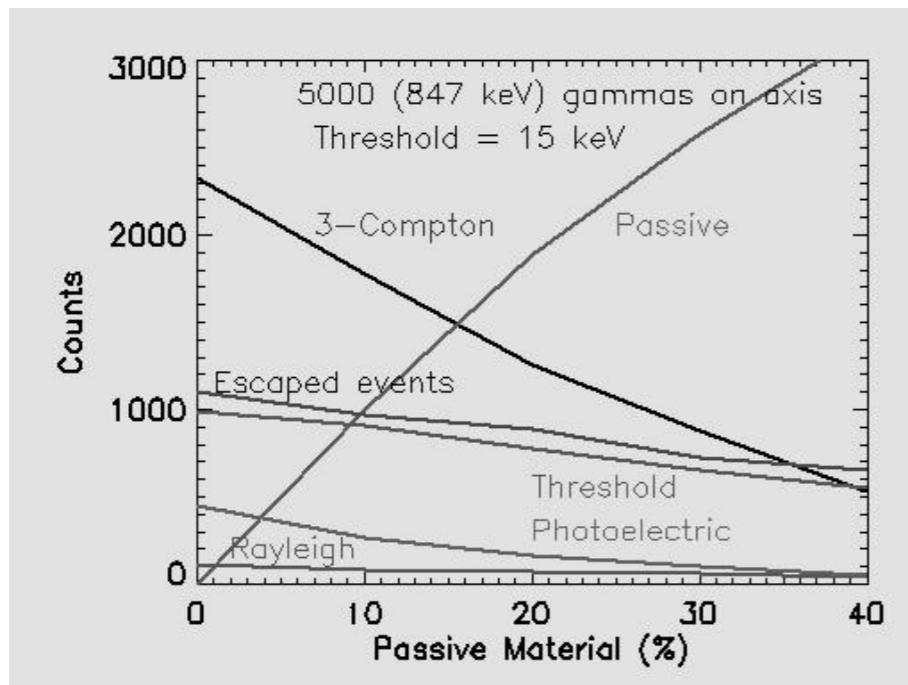
		Hits in Silicon								
		0	1	2	3	4	5	6	7	8
Hits in scintillator	0	16187	331	114	38	18	8	3	1	0
	1	27814	782	324	124	43	10	5	1	1
	2	22573	911	303	107	30	6	1	0	0
	3	9801	462	124	27	6	3	0	0	0
	4	2500	118	24	8	1	0	0	0	0
	5	442	21	4	1	0	0	0	0	0
	6	53	4	1	0	0	0	0	0	0
	7	5	1	0	0	0	0	0	0	0
	8	1	0	0	0	0	0	0	0	0

- We modeled the activity and logged the different types of events
- There are ~3500 coincidences/second between silicon and calorimeter from self activity!
- Lanthanum halides probably not the way to go for large instruments



# Passive material is bad

## Sensitivity Improvement with New Technologies



- Current simulations result in about 2-4% effective area
- This is  $\leq 10\%$  of the potential events that could be used
- Clearly worth effort to substantially improve this performance

**Reduce passive material**  
**Reduce thresholds**

18 August 2005

ACT Team Meeting

- Recall Fermi LAT TKR passive material
  - Even after deleting W, trays are ~50% Si and ~50% passive Al-composite
    - Don't forget also that not all of Si is active



# “Requirements” for high-res Compton tele

---

- Low Z scatterer
  - Minimizes Doppler broadening (most important below MeV)
  - Minimizes MCS of recoil electron, if tracking
- High Z absorber
  - Good stopping power to absorb scattered gamma (and minimize multi-Compton)
- High efficiency
  - Proper scatterer and absorber to give highest possible efficiency
  - Compact (as possible) to maximize geometric cross section for interaction
- Excellent energy resolution
  - Well matched with  $d^3x$
- Fine position resolution
  - Well matched with  $dE$ 
    - Thumb:  $\sim 1$  mm and  $\sim$ few keV are commensurate
- Low-power electronics
  - Preserve intrinsic  $dE$ ,  $d^3x$  of detectors while staying within power budget
- Minimal passive mass within detection volume
  - Interactions can be missed in passive material, and kill Compton performance
  - Minimize structural supports, co-located electronics
  - Minimize detector guard rings



# Backup material

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# Scaling with target material

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## Dependence of photon interaction (mass attenuation coefficient) on material composition

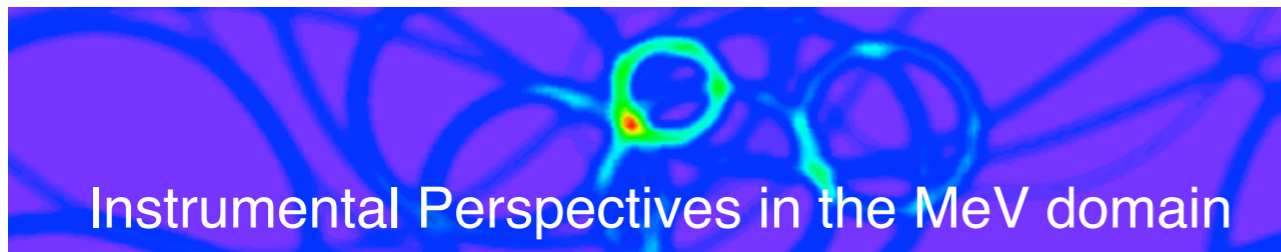
- **Photoelectric effect** (photons see atomic shells)
  - Increases strongly with Z
  - Absorption edges (especially at K shell)
- **Compton scattering** (photons see individual electrons)
  - Scales with electron density (number of electrons per gram)
    - High in hydrogen due to lack of neutrons
    - Only varies by 20% in other elements
- **Pair production** (photons see nucleus)
  - Increases strongly with Z (approximately as  $Z^2$ )
  - High energy limit ( $\gg m_e c^2$ ):

$$\sigma = \frac{7}{9} 4\alpha r_0^2 Z_{nucl} (Z_{nucl} + 1) \ln \left( \frac{183}{3\sqrt{Z_{nucl}}} \right)$$





# Slides borrowed from ... etc.



Instrumental Perspectives in the MeV domain

## GAMMA CUBE

(LE – GLAST)

A scintillation tracker

R. Chipaux, P. Laurent, F. Lebrun, R. Terrier

Instrument Options in the MeV range

What made progress so slow ?

Recent R&D projects towards a future MeV mission



all sky Compton imager

design considerations for Compton Telescopes  
the asCi choice - detector and mission concept  
performance estimates  
one more thing

Peter von Ballmoos, IRAP Toulouse

Peter von Ballmoos, IRAP Toulouse



## Thick Silicon Compton Imager for ACT

### The Nuclear Compton Telescope

A balloon-borne gamma-ray spectrometer, polarimeter, and imager

Steve Boggs

for the NCT collaboration

Bernard Phlips  
Jim Kurfess  
Eric Wulf  
Elena Novikova  
Neil Johnson

18 August 2005

**Simulations of a Si-based  
Advanced Compton Telescope  
(ACT)**

Elena I. Novikova<sup>1</sup>, Eric A. Wulf<sup>1</sup>, Bernard F. Phlips<sup>1</sup>,  
James D. Kurfess<sup>1</sup>, Andreas Zoglauer<sup>2</sup>,  
Georg Weidenspointner<sup>3</sup>, R.Marc Kippen<sup>4</sup>

<sup>1</sup>NRL, Washington, DC, USA  
<sup>2</sup>UC Berkeley, CA, USA  
<sup>3</sup>CESR, Toulouse, France  
<sup>4</sup>LANL, Los Alamos, NM, USA

October 26, 2005      IEEE NSS N19-5 Puerto Rico      1