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A COMPARATIVE ANALYSIS OF STUDIES ON HEAT TRANSFER AND FLUID FLOW IN MICROCHANNELS

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The extremely high rates of heat transfer obtained by employing microchannels makes them an attractive alternative to conventional methods of heat dissipation, especially in applications related to the cooling of microelectronics. A compilation and analysis of the results from investigations on fluid flow and heat transfer in micro- and mini-channels and microtubes in the literature is presented in this review, with a special emphasis on quantitative experimental results and theoretical predictions. Anomalies and deviations from the behavior expected for conventional channels, both in terms of the frictional and heat transfer characteristics, are discussed.

Among the novel methods for thermal management of the high heat fluxes found in microelectronic devices, microchannels are the most effective at heat removal. The possibility of integrating microchannels directly into the heat-generating substrates makes them particularly attractive, since thermal contact resistances may be avoided. The two important objectives in electronics cooling—minimization of the maximum substrate temperature and reduction of substrate temperature gradients—can be achieved by the use of microchannels.

A large number of recent investigations have undertaken to study the fundamentals of microchannel flow, as well as to compare the flow and heat transfer characteristics of microchannels with conventional channels. A comprehensive review of these investigations conducted over the past decade is presented here in concise tabular form.

Predictive correlations have also been proposed in the literature, based on experimental investigations on liquid and gas flow in microchannels. Various combinations of channel size, pitch, and substrate material have been considered. Generally, these correlations have been cast in the same forms as conventional relationships for larger-diameter tubes and channels, but have included modified coefficients. A comparative study of the correlations for single-phase flow is presented in this review.

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NOMENCLATURE

A_c	cross-sectional area	Pr	Prandtl number
A_s	surface area	ΔP	pressure drop
B	slot width	$q_{m,p}$	CHF based on heated channel area
c_a	acoustic velocity	Q	volumetric flow rate
C	coolant heat capacity	R_{th1D}	thermal resistance from 1-D analysis
D	diameter	R_{th3D}	thermal resistance from 3-D analysis
D_h	hydraulic diameter	Re	Reynolds number
f	friction factor	T_i	inlet temperature
G	mass velocity ($\equiv m$, mass flux)	t	tube wall thickness
h_{fg}	latent heat of vaporization	v	inlet velocity
H	height (depth) of microchannel	W	width of microchannel
k_c	coolant thermal conductivity	We	Weber number
L	length	x	distance from stagnation point
Nu	Nusselt number	μ	dynamic viscosity
Nu _{Gn}	Nusselt number (Gnielinski correlation)	ν	kinematic viscosity
P	channel pitch	ρ	density

REVIEW OF THE LITERATURE

Studies on microchannel flows in the past decade are categorized into various topics and summarized in Table 1. The literature survey extends over a wide range of topics such as measurement and estimation of friction factor and heat transfer in microchannels and small-diameter tubes, comparison with flow in conventional channels, investigation of single-phase, boiling, and two-phase flows in microchannels, minichannels, and small tubes, gas flow in microchannels, analytical studies on microchannel flows, and design and testing of microchannel heat sinks for electronics cooling. For each study, key descriptors of the cooling configuration and the primary observations are included.

QUANTITATIVE COMPARISONS

A comparative study of correlations for single-phase flow and heat transfer in microchannels proposed by various investigators is presented in this section. Correlations for friction factor and heat transfer in the laminar and turbulent regimes are compared and contrasted with conventional correlations for macrotubes and channels. Details of each of the studies discussed in this section are available in Table 1.

Friction Correlations

Correlations for friction factor have been proposed based on experiments with nitrogen and water as working fluids [3, 9, 36] in trapezoidal and rectangular channels and microtubes. Peng et al. [13] analyzed water flow in rectangular channels to obtain correlations for various combinations of the channel hydraulic diameter and channel pitch in rectangular channels for laminar and turbulent flow. A plot of the friction factor correlations proposed for laminar and turbulent flow in microchannels is shown in Figure 1. The graph shows the product of friction factor and Reynolds number ($f \cdot \text{Re}$) plotted against the Reynolds number. Conventional correlations are also included for comparison: the Blasius correlation ($f = 0.140 \text{Re}^{-0.182}$) is used for turbulent flow, while for

Table 1. Summary of Microchannel Studies in the Literature

Configuration/parameters	Nature of work	Observations/conclusions	Reference
Microchannel concepts and early work			
Rectangular cross section; water in silicon $W = 50 \mu\text{m}$; $H = 300 \mu\text{m}$; $Q = 4.7, 6.5, 8.6 \text{ cm}^3/\text{s}$	Experiments on integral heat sink for silicon integrated circuits	<ul style="list-style-type: none"> • Demonstrated use of microchannels for very high convective heat transfer in cooling integrated circuits (790 W/cm^2 at a substrate-to-coolant temperature difference of 71°C) 	Tuckerman & Pease (1981) [1]
Microchannels in cooling of integrated circuits	Microchannel fabrication and implementation details discussed	<ul style="list-style-type: none"> • Coolant selection, packaging/headering, microstructure selection, fabrication, and bonding discussed • Etching and precision-sawing compared; fabrication and advantages of "micropillars," using precision-sawing discussed • Expressions for coolant figure of merit (CFOM) provided: CFOM = $(k_c \rho C / \mu)^{0.25}$ for given coolant pressure and $(k_c \rho^2 C^2 / \mu)^{0.25}$ for given pumping power 	Tuckerman & Pease (1982) [2]
Trapezoidal; nitrogen in silicon and glass $W = 130\text{--}300 \mu\text{m}$, $H = 30\text{--}60 \mu\text{m}$, $D_h = 55\text{--}76 \mu\text{m}$	Friction factors measured and compared with Moody's chart values for commercial channels	<ul style="list-style-type: none"> • Friction factor for glass channels 3–5 times larger than smooth-pipe predictions • Flow transition occurred at $Re \approx 400$ • Correlations for friction factor: $f = (110 \pm 8)/Re$ $Re < 900$ $f = 0.165/(3.48 - \log Re)^{2.4} + (0.081 \pm 0.007)$ $900 < Re < 3,000$ 	Wu & Little (1983) [3]
As in [3]	Heat transfer experiments	<ul style="list-style-type: none"> • Correlation for Nusselt number in the turbulent regime: $f = (0.195 \pm 0.017)/Re^{0.11}$ $3,000 < Re < 15,000$ $Nu = 0.0022 Pr^{0.4} Re^{1.09}$ $Re > 3,000$ 	Wu & Little (1984) [4]
Rectangular; air in silicon $W = 0.13\text{--}0.25 \text{ mm}$, $H/W = 10$, $A_s = 47\text{--}63 \text{ cm}^2/\text{cm}^3$	Comparison of performance with conventional heat sinks, based on correlations	<ul style="list-style-type: none"> • Microstructured compact heat sinks attractive compared to conventional air circulation heat sinks 	Mahalingam & Andrews (1987) [5]
Rectangular; water in silicon $W = 50\text{--}600 \mu\text{m}$	Theoretical model for fully developed, developing flows Experiments	<ul style="list-style-type: none"> • Turbulent flow designs showed equivalent or better performance compared to laminar flow designs 	Phillips et al. (1989) [6]
Rectangular, <i>n</i> -propanol in silicon $A_c = 80\text{--}7,200 \mu\text{m}^2$	Microchannel applications discussed	<ul style="list-style-type: none"> • Channels with larger cross-sectional areas showed better agreement with theoretical predictions for the friction factor • Proposed $f = C/Re$, with C given as C-versus-Re graphs (laminar) • Applications of microchannels to electronics cooling, compact heat exchangers, heat shields, and fluid distribution systems discussed 	Prahler et al. (1990) [7] Hoopman (1990) [8]

(continued)

Table 1. (Continued)

Configuration/parameters	Nature of work	Observations/conclusions	Reference
Microtubes; nitrogen in silica $D = 3; 7; 10; 53; 81 \mu\text{m}$, $L = 24\text{--}52 \text{ mm}$	Experiments on friction and heat transfer	<ul style="list-style-type: none"> Correlations for friction factor and Nusselt number: Laminar ($\text{Re} < 2,000$) $f = 64/\text{Re}[1 + 30(\sqrt{D_{eff}})]^{-1}$ Turbulent ($2,500 < \text{Re} < 20,000$) $f = 0.140 \text{ Re}^{-0.182}$ Laminar $\text{Nu} = 0.000972 \text{ Re}^{1.17} \text{ Pr}^{1/3}$ Turbulent $\text{Nu} = 3.82 \times 10^{-6} \text{ Re}^{1.96} \text{ Pr}^{1/3}$ 	Choi et al. (1991) [9]
Rectangular; water in etched silicon $W = 1 \text{ mm}$, $H = 176\text{--}325 \mu\text{m}$, $L = 46 \text{ mm}$, $P = 2 \text{ mm}$	Experiments	<ul style="list-style-type: none"> Nusselt numbers higher than those predicted from analytical solutions for developing laminar flow 	Rahman & Gui (1993) [10]
Single-phase (liquid) experiments			
Rectangular; deionized water in stainless steel $W = 0.6 \text{ mm}$; $H = 0.7 \text{ mm}$, $T_i = 30\text{--}60^\circ \text{C}$, $v = 0.2\text{--}2.1 \text{ m/s}$	Experiments on single-phase forced convection	<ul style="list-style-type: none"> In single-phase convection, a steep increase in wall heat flux with the wall temperature 	Peng & Wang (1993) [11]
Rectangular; water; methanol in stainless steel $W = 0.2, 0.4, 0.6, 0.8 \text{ mm}$, $H = 0.7 \text{ mm}$, $T_i = 10\text{--}35^\circ \text{C}$ (water), $14\text{--}19^\circ \text{C}$ (methanol), $v = 0.2\text{--}2.1 \text{ m/s}$	Experiments on forced-convection flow and heat transfer	<ul style="list-style-type: none"> Heat flux for microchannels higher than for normal-size tube Heat transfer augmented as liquid temperature was reduced and as liquid velocity was increased Fully developed turbulent convection regime starts at $\text{Re} = 1,000\text{--}1,500$ Correlation for turbulent heat transfer $\text{Nu} = 0.00805 \text{ Re}^{4/5} \text{ Pr}^{1/3}$ Flow transition occurred for $\text{Re} = 200\text{--}700$ Correlations proposed (values for $C_{f,i}$, $C_{f,t}$ provided in Table 2) $f = C_{f,i}/\text{Re}^{1.98}$ laminar flow $f = C_{f,t}/\text{Re}^{1.72}$ turbulent flow 	Wang & Peng (1994) [12] Peng et al. (1994) [13]
Rectangular; water in stainless steel $D_h = 0.133\text{--}0.367 \text{ mm}$, $L = 50 \text{ mm}$, $H/W = 0.333\text{--}1$, $T_i = 22\text{--}44^\circ \text{C}$, $v = 0.25\text{--}12 \text{ m/s}$, $\text{Re} = 50\text{--}4,000$	Experiments on frictional behavior in laminar and turbulent flow	<ul style="list-style-type: none"> Correlations proposed (values for $C_{f,i}$, $C_{f,t}$ provided in Table 2) 	Peng et al. (1994) [14]
As in [13]	Experiments on forced-convection heat transfer characteristics	<ul style="list-style-type: none"> Fully turbulent convective conditions reached at $\text{Re} = 400\text{--}1,500$ Transition Re diminished with a reduction in microchannel dimension $\text{Nu} = C_{h,i} \text{ Re}^{0.62} \text{ Pr}^{1/3}$ laminar $\text{Nu} = C_{h,t} \text{ Re}^{0.8} \text{ Pr}^{1/3}$ turbulent (values for $C_{h,i}$, $C_{h,t}$ provided in Table 2) 	Peng & Peterson (1995) [15]
As in [12] except $T_i = 11\text{--}28^\circ \text{C}$ (water), $12\text{--}20^\circ \text{C}$ (methanol) $v = 0.2\text{--}2.1 \text{ m/s}$ (water), $0.2\text{--}1.5 \text{ m/s}$ (methanol)	Experiments on effect of thermofluid properties and geometry on convective heat transfer	<ul style="list-style-type: none"> Changes in flow regimes and heat transfer modes initiated at lower Re in microchannels compared to conventional channels Transition zone and heat transfer characteristics in laminar and transition flow influenced by liquid temperature, velocity, Re, and microchannel size 	Peng & Peterson (1996) [16]
As in [13]	Experiments on single-phase flow and heat transfer	<ul style="list-style-type: none"> Ratio of experimental to theoretical friction factor at critical Re plotted as a function of $Z (= \min\{H, W\} / \max\{H, W\})$ Correlations proposed $\text{Nu} = 0.1165(D_h/P)^{0.81}(H/W)^{-0.79} \text{ Re}^{0.62} \text{ Pr}^{0.33}$ laminar $\text{Nu} = 0.072(D_h/P)^{1.15}[1 - 2.421(Z - 0.5)^2] \text{ Re}^{0.8} \text{ Pr}^{0.33}$ turbulent 	Peng & Peterson (1996) [16]

Rectangular; water-methanol mixture in stainless steel $D_h = 0.133\text{--}0.367$ mm; $L = 50$ mm; $W = 0.1, 0.2, 0.3, 0.4$ mm; $H = 0.2, 0.3$ mm; $T_i = 14\text{--}36$ °C; $v = 0.04\text{--}3.8$ m/s; $Re = 6\text{--}3,500$	Experiments	<ul style="list-style-type: none"> Laminar heat transfer ceased for $Re \approx 70\text{--}400$ depending on flow conditions; fully developed turbulent heat transfer achieved at $Re = 200\text{--}700$, depending on D_h Transition Re reduced with a reduction in microchannel size $D_h/H/W$, and mixture mole fraction influenced heat transfer Heat transfer increased for smaller mole fractions of the more volatile component 	Peng & Peterson (1996) [17]
Rectangular; deionized water in silicon $W = 251$ μm , $H = 1,030$ μm , $D_h = 404$ μm , $L = 2.5$ cm, $Q = 5.47\text{--}118$ cm ³ /s	Experimental & theoretical study	<ul style="list-style-type: none"> Critical Re of 1,500 identified for onset of turbulence Analysis showed that flow and heat transfer performance could be improved by increasing H, and that for the same pressure drop and pumping power, thermal resistance was smaller for deeper channels 	Harms et al. (1997) [18]
Rectangular; FC-72 and transformer oil in stainless steel $H = 0.10\text{--}0.58$ mm; nozzle dimensions (mm): length = 35, $B = 0.146, 0.210, 0.234$, height = 12, $v = 0.54\text{--}8.45$ m/s $Re = 70\text{--}170$ (oil), 911–4807 (FC72)	Experiments in impingement on 2-D microchannels	<ul style="list-style-type: none"> Empirical correlation for Nusselt number for the two liquids: $Nu_x = 0.429 Re^{0.583} Pr^{1/3} (x/2H)^{0.349} (B/2H)^{-0.494}$ 	Zhuang et al. (1997) [19]
Circular; distilled water in copper $D = 0.102\text{--}1.09$ mm, $v < 18.9$ m/s, $Re = 2.6 \times 10^3\text{--}2.3 \times 10^4$, $Pr = 1.53\text{--}6.43$, $q' < 3.0$ MW/m ²	Experiments on turbulent single-phase flow	<ul style="list-style-type: none"> Nusselt numbers higher than those predicted by large-channel correlations Gnielinski [71] correlation modified for Nusselt number for turbulent flow in circular microchannels (f' from Filonenko, 1954): $Nu = Nu_{Gn}(1 + F)$, where $F = C Re[1 - (D/D_o)^2]$ $Nu_{Gn} = (f'/8)(Re - 1,000) Pr / [1 + 12.7(f'/8)^{1/2}(Pr^{2/3} - 1)]$ $C = 7.6 \times 10^{-3}$; $D_o = 1.164$ mm, $f = [1.82 \log(Re) - 1.64]^{-2}$ 	Adams et al. (1998) [20]
Non-circular; water in copper $D_h = 1.13$ mm, $Re = 3.9 \times 10^3\text{--}2.14 \times 10^4$, $Pr = 1.22\text{--}3.02$	Experiments on turbulent convection	<ul style="list-style-type: none"> Experimental Nusselt number well predicted by Nu_{Gn} $D_h \approx 1.2$ mm proposed as reasonable lower limit for applicability of standard Nusselt-type correlations to noncircular channels 	Adams et al. (1999) [21]
Rectangular laminar and transition flow	Dimensional analysis based on experimental data in the literature	<ul style="list-style-type: none"> Attempted to explain the observation that Nu may decrease with increasing Re in laminar regime and may remain unaffected in transition regime Proposed that Brinkman number may better correlate convective heat transfer 	Tso & Mahuliakar (1998, 1999) [22, 23]
Almost circular; water in aluminum $D_h = 0.73$ mm	Experiments	<ul style="list-style-type: none"> Laminar flow data found to correlate well using Brinkman number 	Tso & Mahuliakar (2000) [24]
Single-phase (liquid) models and optimization studies			
Triangular microgrooves channel angle 20–60	Analytical/numerical analysis	<ul style="list-style-type: none"> Friction factor–Reynolds number product strongly dependent on channel angle, contact angle, and dimensionless vapor–liquid interface flow number 	Ma et al. (1994) [25]

(continued)

Table 1. (Continued)

Configuration/parameters	Nature of work	Observations/conclusions	Reference
Microchannel plate-fin heat sink; air in copper, aluminum $W = 400, 500 \mu\text{m}$, $H = 2.5 \text{ cm}$, $\dot{Q} = 1\text{--}6$ liters/s Circular capillary channels $D = 8, 1\text{--}96 \mu\text{m}$, $0.76\text{--}4.7 \mu\text{m}$	Thermal resistance model, experiments, optimization Numerical study on the flow of superfluid Helium using a "two-fluid model"	<ul style="list-style-type: none"> Thermal resistance of microchanneled heat sink lower than for heat sinks employing direct air cooling, by a factor of more than 3 Existence of an optimum channel diameter for maximum mass flow rate indicated 	Kleiner et al. (1995) [26] Takamatsu et al. (1997) [27]
Rectangular: fluorocarbon in silicon $P = 100\text{--}1,000 \mu\text{m}$, $H = 150\text{--}200 \mu\text{m}$, $W = 56.6\text{--}113.4 \mu\text{m}$, $v = 0.1\text{--}1.0 \text{ m}^3/\text{s}$	Numerical analysis of manifold microchannel heat sinks	<ul style="list-style-type: none"> 3-D model showed close agreement with simple 1-D model at high inlet velocity Numerical results showed much weaker effect of W compared to analytical results L had almost no effect on thermal resistance and affected only pressure drop 	Copeland et al. (1997) [28]
Parallel plates at $25 \mu\text{m}$ separation dilute aqueous electrolyte $L = 10 \text{ mm}$ Parallel plates ($10 \times 20 \text{ mm}$) of P-type silicon and glass at $10\text{--}280 \mu\text{m}$ separation $\Delta P = 0\text{--}350 \text{ mbar}$ Rectangular, dilute aqueous electrolyte in silicon $H = 20 \mu\text{m}$, $W = 30 \mu\text{m}$, $L = 10 \text{ mm}$, $\Delta P = 2 \text{ atm}$, $T_i = 298 \text{ K}$, $q'' = 1.0 \times 10^5 \text{ W/m}^2$	Theoretical analysis incorporating effects of electric double layer field Experimental study and comparison with predicted volume flow rates Numerical analysis with effects of EDL and flow-induced electrokinetic field	<ul style="list-style-type: none"> EDL resulted in a reduced flow velocity than in conventional theory, thus affecting temperature distribution and reducing Re Higher heat transfer predicted without the double layer For solutions of high ionic concentration as well as for $D_h > a$ few hundred μm, EDL effect negligible EDL effect becomes significant for dilute solutions The EDL field and electrokinetic potential act against the liquid flow, resulting in higher friction coefficient, reduced flow rate, and reduced Nusselt number, for dilute solutions 	Mala et al. (1997) [29] Mala et al. (1997) [30] Yang et al. (1998) [31]
Rectangular (flat-plate micro heat exchangers)	Optimization study on microchannel shape Comparative analysis of jet impingement and microchannel cooling	<ul style="list-style-type: none"> Width of heat exchanger conduits may be optimized to reduce maximum temperature of the uniformly heated surface Thermal performance of jet impingement without any treatment of spent flow substantially lower than microchannel cooling, regardless of target dimension Microchannel cooling preferable for target dimensions smaller than $7 \times 7 \text{ cm}$ 	Bau (1998) [32] Lee and Vafai (1999) [33]
Microchannel cooling and jet impingement			
Gas flow			
Rectangular: helium in silicon $W = 52, 25 \mu\text{m}$, $H = 1.33 \mu\text{m}$, $L = 7, 500 \mu\text{m}$; inlet-to-outlet pressure ratio = $1.2\text{--}2.5$, $\text{Re} = (0.5\text{--}4) \times 10^{-3}$ As in [34] with pressure ratio = $1.6\text{--}4.2$, $\text{Re} = (1.4\text{--}12) \times 10^{-3}$	Flow rates measured and compared with theoretical model Experiments and comparison of mass flow with results from 2-D analysis with slip boundary condition	<ul style="list-style-type: none"> Mass flow–pressure relationship accurately modeled by including a slip flow boundary condition at the wall Discussions on nondimensional formulation and perturbation solution 	Arkilic et al. (1994) [34] Arkilic et al. (1997) [35]

Microtubes; nitrogen and water in silica $D = 19, 52, 102 \mu\text{m}$; $\text{Pr} = 0.7-5$; $\text{Re} = 250-20,000$	Experiments; theoretical scaling analysis	<ul style="list-style-type: none"> Turbulent momentum and energy transport in the radial direction significant in the near-wall zone of a microtube Correlations proposed: <ul style="list-style-type: none"> $f = 50.13/\text{Re}$ (laminar, $\text{Re} < 2,000$) $f = 0.302/\text{Re}^{0.25}$ (transition, $2,000 < \text{Re} < 6,000$) $\text{Nu} = 0.007 \text{Re}^{1.2} \text{Pr}^{0.2}$ (turbulent, $6,000 < \text{Re} < 20,000$) Heat flux on the channel surface decreases with increase in Knudsen number and channel length in supersonic flow 	Yu et al. (1995) [36]
Rectangular $H = 0.5, 5 \mu\text{m}$; $H/W = 2.5, 5, 10, 20$ (subsonic); $5, 10, 20$ (supersonic)	Numerical study using direct simulation Monte Carlo technique	<ul style="list-style-type: none"> Nusselt number and friction coefficient substantially reduced for slip flows compared to continuum flows Effect of compressibility significant at high Re 	Mavriplis et al. (1995) [37]
Rectangular helium (as in [34]) helium and nitrogen, $D_{\text{He}} = 1.01 \mu\text{m}$, $L = 10.9 \text{ mm}$ (as in [38]) Gas flow	2-D numerical model, comparison with experiments in literature	<ul style="list-style-type: none"> Local Nusselt number increased with dimensionless length, due to compressibility 	Kavehpour & Faghri (1997) [39]
Rectangular; nitrogen, helium in silicon $W = 40 \mu\text{m}$, $H = 1.2 \mu\text{m}$, $L = 3 \text{ mm}$ (N_2), $W = 52 \mu\text{m}$, $H = 1.33 \mu\text{m}$, $L = 7.5 \text{ mm}$ (He)	Numerical solution with slip boundary condition	<ul style="list-style-type: none"> f-Re product not constant; dependent on Re Small velocities and high pressure gradients due to large wall shear stresses Comparisons with experiments of [34] 	Guo & Wu (1997) [40] Chen et al. (1998) [41]
3-D straight and spiral grooves	Numerical study on slip flow in long microchannels	<ul style="list-style-type: none"> Nonlinear pressure gradients along the microchannels due to density variations 	Niu (1999) [42]
Boiling in microchannels			
Circular: R-113 in copper $D = 2.45 \text{ mm}$ (mini), $510 \mu\text{m}$ (micro) $Q = 19-95 \text{ ml/min}$, ΔT : $10-32^\circ\text{C}$ As in [43]	Experiments on boiling & two- ϕ flow; boiling curves & CHF values obtained Pressure drop model developed; predictions compared to experiments	<ul style="list-style-type: none"> Microchannel yielded higher CHF (28% greater at $Q = 64 \text{ ml/min}$) than minichannel, with a larger ΔP (0.3 bar for micro, 0.03 bar for mini) Major contributor to pressure drop identified as the acceleration resulting from evaporation Compressibility effect important for microchannel when Mach number > 0.22 Channel erosion effects more predominant in microchannels than in minichannels 	Bowers & Mudawar (1994) [43] Bowers & Mudawar (1994) [44]
As in [43]	Experiments on boiling and two-phase flow Experiments on subcooled boiling of water	<ul style="list-style-type: none"> Single CHF correlation for mini- and microchannels developed: $q_{\text{me,p}}/(Gh_{\text{f,g}}) = 0.16 \text{ We}^{-0.19} (L/D)^{0.54}$ Nucleate boiling intensified and wall superheat for flow boiling smaller in microchannels than in normal-sized channels for the same wall heat flux No partial nucleate boiling observed in microchannels 	Bowers & Mudawar (1994) [45] Peng & Wang (1993) [11]

(continued)

Table 1. (Continued)

Configuration/parameters	Nature of work	Observations/conclusions	Reference
Rectangular; methanol in stainless steel $W = 0.2, 0.4, 0.6$ mm, $H = 0.7$ mm, $L = 45$ mm, $P = 2.4-4$ mm; $T_i = 14-19^\circ\text{C}$ (Subcooling: $45-50^\circ\text{C}$), $v = 0.2-1.5$ m/s	Experiments on boiling	<ul style="list-style-type: none"> Liquid velocity and subcooling do not affect fully developed nucleate boiling Greater subcooling increased velocity and suppressed initiation of flow boiling 	Peng et al. (1995) [46]
Rectangular; methanol-water mixture in stainless steel $W = 0.1, 0.2, 0.3, 0.4$ mm; $H = 0.2, 0.3$ mm; $L = 45$ mm; $D_h = 0.133-0.343$ mm; $v = 0.1-4.0$ m/s; $T_i = 18-27.5^\circ\text{C}$ (Subcooling: $38-82^\circ\text{C}$)	Experiments on flow boiling in binary mixtures	<ul style="list-style-type: none"> Heat transfer coefficient at onset of flow boiling and in partial nucleate boiling greatly influenced by concentration, microchannel/substrate dimensions, flow velocity, and subcooling These parameters had no significant effect on heat transfer coefficient in the fully nucleate boiling regime Mixtures with small concentrations of methanol augmented flow boiling heat transfer 	Peng et al. (1996) [47]
V-shaped; water and methanol in stainless steel Groove angle: $30-60^\circ$; $D_h = 0.2-0.6$ mm v (water) = $0.31-1.03$ m/s v (methanol) = $0.12-2.14$ m/s	Experiments on flow boiling	<ul style="list-style-type: none"> Heat transfer and pressure drop were affected by flow velocity, subcooling, D_h, and groove angle No bubbles observed in microchannels during flow boiling, unlike in conventional channels Experiments indicated an optimum D_h and groove angle Analytical expression developed for the evaporating film profile 	Peng et al. (1998) [48]
V-shaped Circular and rod bundle; water in copper $D = 1.17, 1.45$ mm, $D_h = 1.131$ mm $m = 250-1,000$ kg/m ² s Exit pressure = $344-1,043$ kPa Inlet pressure = $407-1,204$ kPa $T_i = 49-72.5^\circ\text{C}$	Analysis of microgrooves with nonuniform heat input Analysis of axial flow of evaporating thin film Experiments on CHF in flow of subcooled water	<ul style="list-style-type: none"> Used perturbation method to solve the axial flow of an evaporating thin film through a V-shaped microchannel with tilt CHF found to increase monotonically with increasing mass flux or pressure CHF depends on the channel cross section geometry, and increases with increasing D 	Ha & Peterson (1996) [49] Ha & Peterson (1998) [50] Roach et al. (1999) [51]
Circular; water in stainless steel $D = 2.5$ mm, $t = 0.25$ mm, $v = 10-40$ m/s	Experiments on subcooled flow boiling of water under high heat fluxes	<ul style="list-style-type: none"> Experimental data did not match predictions from CHF correlations in the literature 	Celata et al. (1993) [52]
Rectangular; water and R141b in copper $W = 1, 2, 3$ mm, $H/W < 3$, $m = 50, 200, 300$ kg/m ² s	Experiments on flow boiling in narrow channels of planar heat exchanger elements	<ul style="list-style-type: none"> Boiling curves and variations of heat transfer coefficient with local and average heat fluxes obtained 	Mertz et al. (1996) [53]

<p>Rectangular; FC-72 in fibreglass $W = 5$ mm, $H = 2.5$ mm, heated length = 101.6 mm, $v = 0.25\text{--}10$ m/s, $Re = 2,000\text{--}130,000$ Subcooling at outlet = 3, 16, 29°C As in [54]</p>	<p>CHF experiments on long channels; flow visualization</p> <p>Theoretical model for CHF; data analysis</p>	<ul style="list-style-type: none"> • Propagation of vapor patches resembling a wavy vapor layer along the heated wall at the critical heat flux • Length and height of vapor patch found to increase along flow direction, and decreased with increasing subcooling and velocity • Effect of periodic distribution of vapor patches idealized as a sinusoidal interface with amplitude and wavelength increasing in flow direction 	<p>Sturgis & Mudawar (1999) [54]</p> <p>Sturgis & Mudawar (1999) [55]</p>
Two-phase flow			
<p>Rectangular; R124 in copper $W = 0.27$ mm, $H = 1.0$ mm, $D_h = 425$ mm, $Re_{Dh} = 100\text{--}750$; $q'' < 40$ W/cm² Rectangular; R124 in copper $W = 270$ μm, $H = 1,000$ μm, $L = 2.052$ cm, $D_h = 425$ μm; inlet subcooling: 5–15°C; $\dot{Q} = 35\text{--}300$ ml/min</p>	<p>Experiments on microchannel heat exchanger</p> <p>Experiments with two microchannel patterns (parallel and diamond)</p> <p>Visual observation of flow patterns and pattern maps</p>	<ul style="list-style-type: none"> • Nusselt number (~5 to 12) showed an increase with Reynolds number in single-ϕ flow, but was approximately constant in two-ϕ flow • Heat transfer coefficient and pressure drop found to be functions of flow quality and mass flux, in addition to the heat flux and surface superheat • Heat transfer coefficient decreased by 20–30% for an increase in exit vapor quality from 0.01 to 0.65 • Bubbly, churn, slug, slug-annular, and annular flow patterns observed 	<p>Cuta et al. (1996) [56]</p> <p>Ravigururajan (1998) [57]</p> <p>Triplett et al. (1999) [58]</p>
<p>Circular and semitriangular; air–water mixture in glass $D = 1.1, 1.45$ mm, $D_h = 1.09, 1.49$ mm, $v(\text{air}) = 0.02\text{--}80$ m/s, $v(\text{water}) = 0.02\text{--}8$ m/s (superficial velocity) As in [58]</p>	<p>Frictional pressure drops measured and compared with various two-ϕ friction models</p> <p>Experiments, flow visualization</p>	<ul style="list-style-type: none"> • Models and correlations overpredicted channel void fraction and pressure drop in annular flow pattern • Annular flow interface momentum transfer and wall friction in microchannels significantly different from those in larger channels • Tube diameter influences the superficial gas and liquid velocities at which flow transitions take place, due to combined effect of surface tension, hydraulic diameter and aspect ratio 	<p>Triplett et al. (1999) [59]</p> <p>Coleman & Garimella (1999) [60]</p>
Design and testing			
<p>Rectangular; water in silicon</p>	<p>Numerical solution for temperature field; comparison with experiments [61]</p>	<ul style="list-style-type: none"> • Design algorithm developed for selection of heat exchanger dimensions • Expression for maximum pumping power obtained as function of channel geometry 	<p>Weisberg et al. (1992) [62]</p>

(continued)

Table 1. (Continued)

Configuration/parameters	Nature of work	Observations/conclusions	Reference
Almost rectangular; water in copper 0.5×12 mm, 0.125×12 mm, $Q = 0.47\text{--}5$ gpm	Design and testing, microchannel heat exchanger for laser diode arrays	<ul style="list-style-type: none"> Thermal resistance due to solder bond estimated 	Roy & Avanic (1996) [63]
Rectangular and almost triangular; air in copper, aluminum	Parametric studies and experiments of air impingement in microchannels	<ul style="list-style-type: none"> Thermal resistance model developed Parametric studies to determine influence of static pressure, pumping power and geometric parameters on thermal resistance Rectangular geometry had the lowest thermal resistance 	Aranyosi et al. (1997) [64]
Rectangular; diamond-shaped and hexagonal; water in silicon	3-D numerical model; optimization for reducing thermal resistance	<ul style="list-style-type: none"> Power densities of $230\text{--}350$ W/cm² dissipated with a temperature rise of 3.5 C, and a pumping power of about 1 W per chip Parameter "heat spread effect" defined: 	Perret et al. (1998) [65]
Rectangular; water; FC72 in copper	Experiments on micro heat sink for power multichip module; 3-D and 1-D thermal resistance models	<ul style="list-style-type: none"> Power densities of $230\text{--}350$ W/cm² dissipated with a temperature rise of 3.5 C, and a pumping power of about 1 W per chip Parameter "heat spread effect" defined: 	Gillot et al. (1998) [66]
Rectangular; water; FC72 in copper $W = 230, 311$ μ m, $H = 730, 3,040$ μ m, Q (ml/min) = $1,350$ (water, $1-\phi$) and 30 ($2-\phi$); $2,000$ (FC 72 $1-\phi$) and 300 ($2-\phi$)	Experiments on single and two-phase micro heat exchangers for cooling transistors	<ul style="list-style-type: none"> Two-phase heat exchanger provided lower thermal resistance and pressure drop compared to single-phase heat exchangers 	Gillot et al. (1999) [67]
Rectangular; air in copper $W = 800$ μ m, $H = 50$ mm, $Q = 140$ m ³ /h	Experiments and thermal resistance model	<ul style="list-style-type: none"> Pressure drop found to have large deviation from predicted values at high air flow rates Cooling capacity $\approx 1,700$ W at heat flux ≈ 15 W/cm² 	Yu et al. (1999) [68]
Measurement techniques			
Triangular; water in silicon $W = 28\text{--}182$ μ m, $Q = 0.01\text{--}1,000$ μ l/min	Optical flow measurements using microscope	<ul style="list-style-type: none"> Measured flow rates in good agreement with theoretical values for laminar flow through triangular channels 	Richter et al. (1997) [69]
Rectangular; water in glass $W = 300, H = 30, L = 25$ mm	Particle image velocimetry	<ul style="list-style-type: none"> Results agreed well with analytical solutions for Newtonian flow in rectangular channels 	Meinhart et al. (1999) [70]

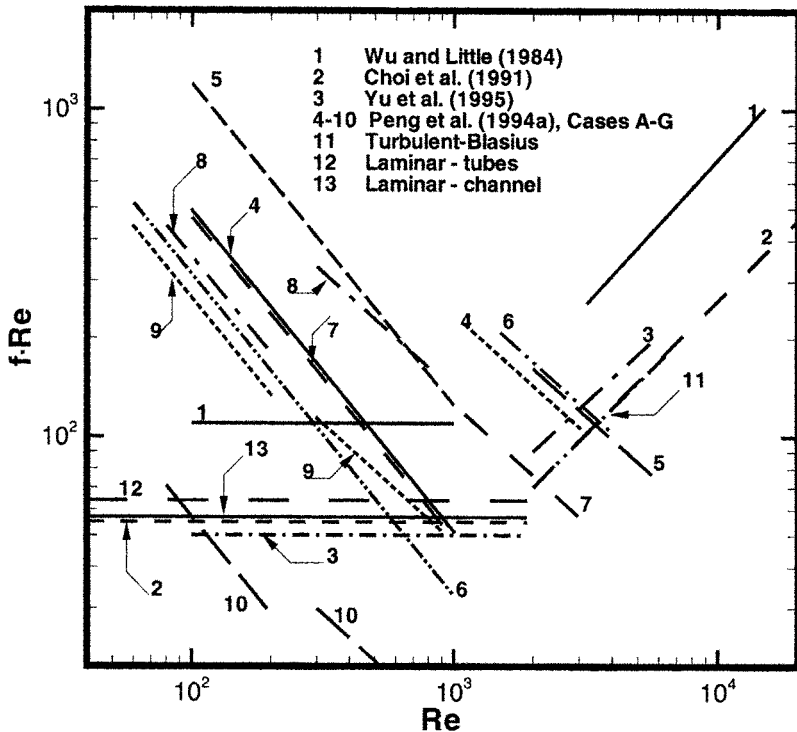


Figure 1. Friction-factor predictions from the literature for microchannels and conventional channels, in the laminar and turbulent regimes.

laminar flow, circular-pipe ($f = 64/Re$) and square-channel predictions ($f = 57/Re$) are shown. The $f \cdot Re$ product is independent of Reynolds number for laminar flow in conventional channels. In the turbulent regime, the friction factor is almost independent of the Reynolds number ($f \cdot Re$ increases linearly with Re). The predictions in the literature for microchannels may be analyzed with greater ease by considering the laminar and turbulent regions separately.

Predictions of $f \cdot Re$ in the laminar regime are shown in Figure 2. The correlations of Wu and Little [3], Choi et al. [9], and Yu et al. [36] predict constant values of $f \cdot Re$, with the magnitude of this product being greater than for conventional channels in Wu and Little (110), and lower in Choi et al. [55] and Yu et al. [50], respectively. Predictions from Peng et al. [13] for water flow in rectangular microchannels (see Table 1 for details) show an altogether different trend: in all cases, $f \cdot Re$ decreases with an increase in the Reynolds number. For cases A, B, and C from Peng et al. (in which $D_h \geq 267 \mu\text{m}$), the laminar regime extends to $Re \approx 700$, whereas for cases E, F, and G ($D_h < 200 \mu\text{m}$), the onset of turbulence occurs as early as $Re = 300$. (As in the original work, the laminar plots for cases A–D are extended until $Re = 1,000$). While the slopes of the curves for all test cases are identical ($Re^{-0.98}$), the magnitude of $f \cdot Re$ is highest for the largest microchannels (D_h) and lowest for the smallest.

The friction correlations in the turbulent regime are compared with conventional correlations in Figure 3. Predictions for nitrogen flow from Choi et al. agree very well

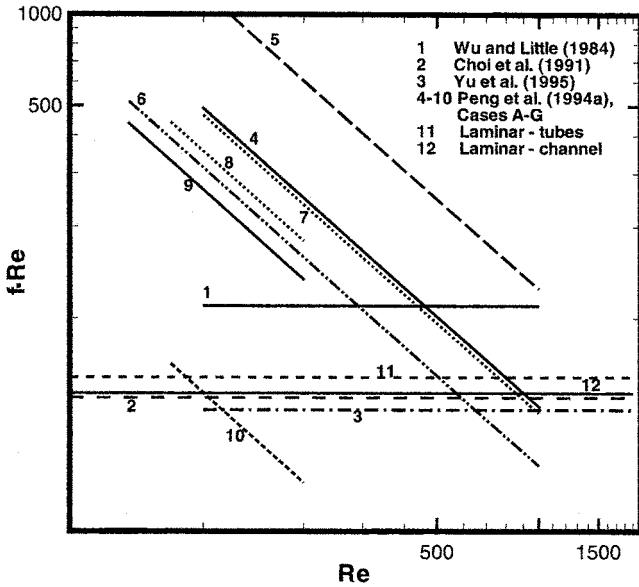


Figure 2. Friction-factor predictions in the laminar regime.

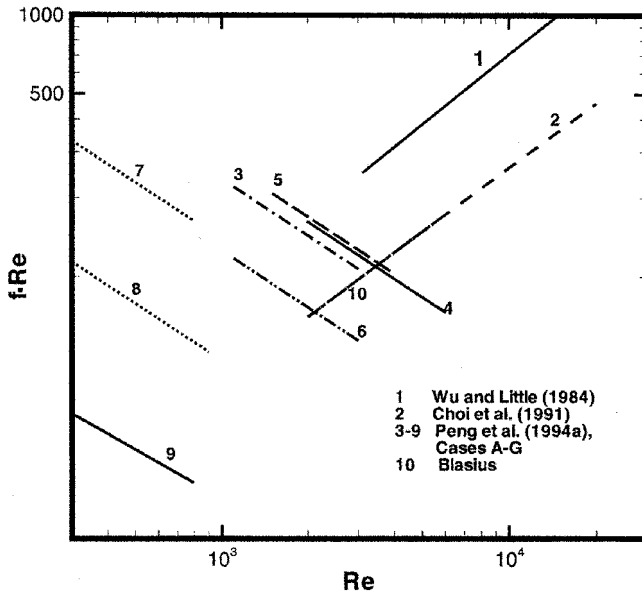


Figure 3. Friction-factor predictions in the turbulent regime.

Table 2. Microchannel configurations and coefficients from Peng et al. [13, 14]

Case	W (mm)	H (mm)	L (mm)	D_h (mm)	H/W	Re_{cr}	$C_{f;l}$	$C_{f;t}$	$C_{h;l}$	$C_{h;t}$
A	0.4	0.3	50	0.343	0.75	700	44800	34200	0.058	0.0134
B	0.3	0.3	50	0.3	0.1	700	109000	38600	0.0384	0.00726
C	0.4	0.2	50	0.267	0.5	700	28600	40400	0.0426	0.0166
D	0.3	0.2	50	0.24	0.667	400	42600	18200	0.0472	*
E	0.2	0.2	50	0.2	1	200	32400	20100	0.0468	0.00696
F	0.3	0.1	50	0.15	0.333	200	24200	6920	0.0104	0.00483
G	0.2	0.1	50	0.133	0.5	200	5200	1820	0.0285	0.00939

with conventional results; the Wu and Little correlation is similar to these two in its trend of variation, but the predicted values are much higher in magnitude. The correlations of Peng et al. [13] for water flow again exhibit a very different trend: in all cases, $f \cdot Re$ decreases with an increase in Reynolds number (as $Re^{-0.72}$), in contradiction to conventional correlations. The onset of turbulence is also seen to occur much earlier for the microchannels studied by Peng et al. Another observation of interest in Figure 3 is that the $f \cdot Re$ values predicted by Peng et al. decrease in magnitude as the channel hydraulic diameter decreases; the drop in $f \cdot Re$ with D_h is very steep when D_h becomes smaller than 200 μm .

Heat Transfer Correlations

Correlations for the average Nusselt number in microchannels in terms of the Reynolds and Prandtl numbers have been proposed in the literature for laminar and turbulent regimes, based on experiments with a range of fluid–substrate combinations, channel dimensions, and configurations, as summarized in Table 1.

Heat transfer correlations for nitrogen flow [4, 9, 36], water flow in rectangular microchannels [12, 14, 16], water flow in circular microchannels [20, 36], and methanol in rectangular channels [12] are considered for comparison. Figure 4 shows a composite plot of predicted values of $Nu/Pr^{0.33}$ as a function of Reynolds number from the correlations in these studies. Conventional channel correlations are also included for comparison: the Dittus-Boelter correlation for turbulent flow in conventional channels, and for laminar flow, $Nu_{D_h} = 1.86(Re_{D_h} Pr)^{0.33}(D_h/L)^{0.33}$; a sample set of parameters ($L = 50$ mm and $D_h = 0.24$ mm) is used to compute values from this correlation. A significant amount of scatter is seen in these plots, as was true for predictions of friction factor, with the predictions of Choi et al. [9] and Yu et al. [36] being among the highest in the turbulent regime. All predictions reflect an increase in Nusselt number with increasing Re .

The heat transfer correlations are again considered separately in the laminar and turbulent regimes in Figures 5 and 6, respectively. The end of the laminar regime was identified to be at quite different Reynolds numbers in the studies considered, as noted with the friction factor predictions. The dependence of Nusselt number on Reynolds number is stronger in all the microchannel predictions when compared to conventional results, as indicated by the steeper slopes of the former; Choi et al. [9] predict the strongest variation of Nusselt number with Re . Also, the predictions for all cases from Peng et al. fall below those for a conventional channel.

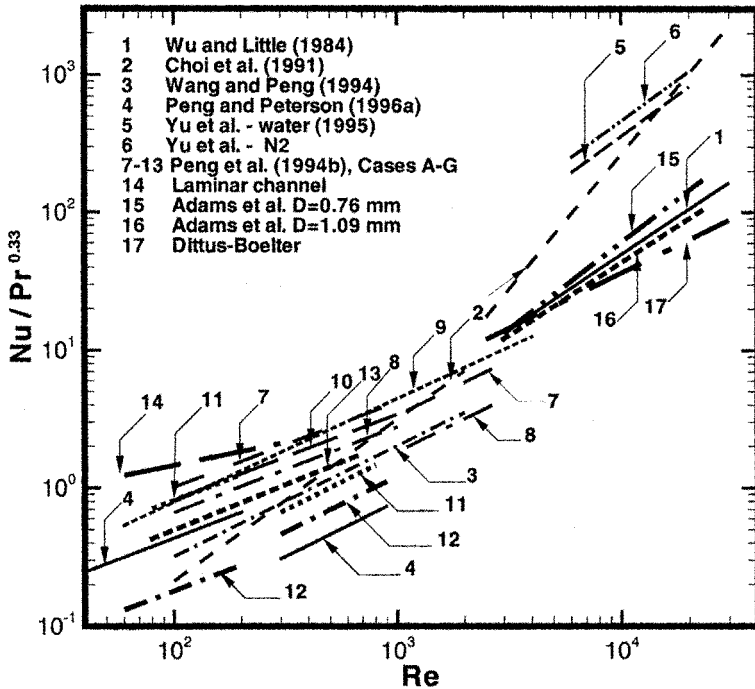


Figure 4. Heat transfer predictions from the literature for microchannels and conventional channels, in the laminar and turbulent regimes.

In the turbulent regime (Figure 6), the predictions of all investigators with the exception of Peng et al. [14] and Peng and Peterson [16] fall above the conventional channel values. In particular, Adams et al. [20] and Wu and Little [4] lie in one group. The predictions of Choi et al. [9] and Yu et al. [36] are also somewhat comparable, and lie in a different group. It may be noted that these groups are not divided by fluid type (since both groups include results for nitrogen and water) or by microchannel dimensions. The rectangular microchannels of different dimensions (Table 1) considered in Peng et al. [14] exhibit a large variation in predicted Nusselt numbers. In all these results, as well as for Peng and Peterson [16], the predicted values lie below those from the Dittus-Boelter correlation. Turbulent heat transfer predictions for case D are not included in this comparison since it appears that the value of $C_{h,t}$ for this case may have been erroneously listed in [14] as 0.0926, and instead, should have been 0.00926. This latter value would more closely match other values for $C_{h,t}$, and would also result in the predictions for case D lying in the same group as cases A, B, and C.

CONCLUSION

A comparative study of the results of investigations in the literature on flow and heat transfer in microchannels has been compiled in tabular form, under various research topics. Correlations for single-phase friction factor and Nusselt number proposed by various investigators based on their experiments have been compared and contrasted

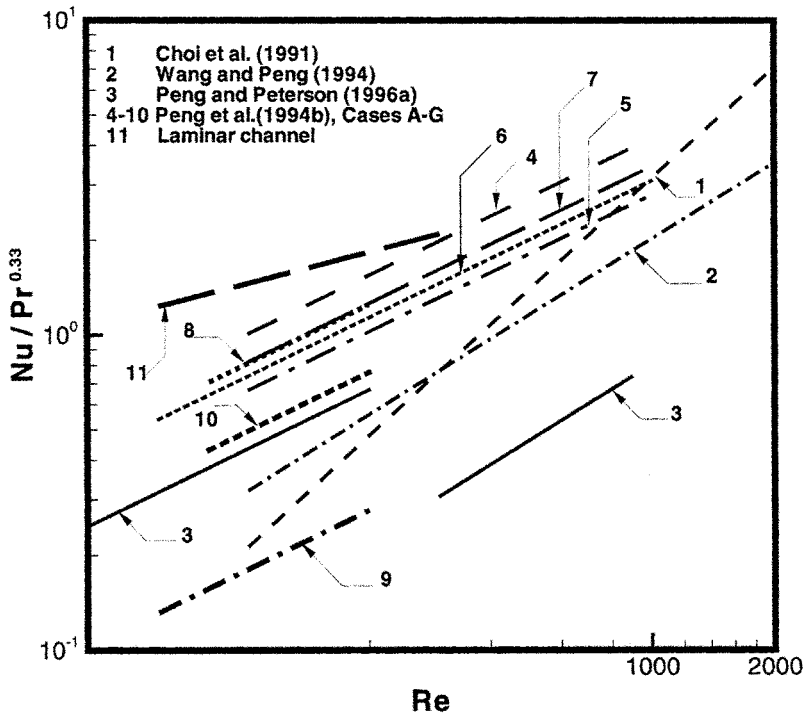


Figure 5. Heat transfer predictions in the laminar regime.

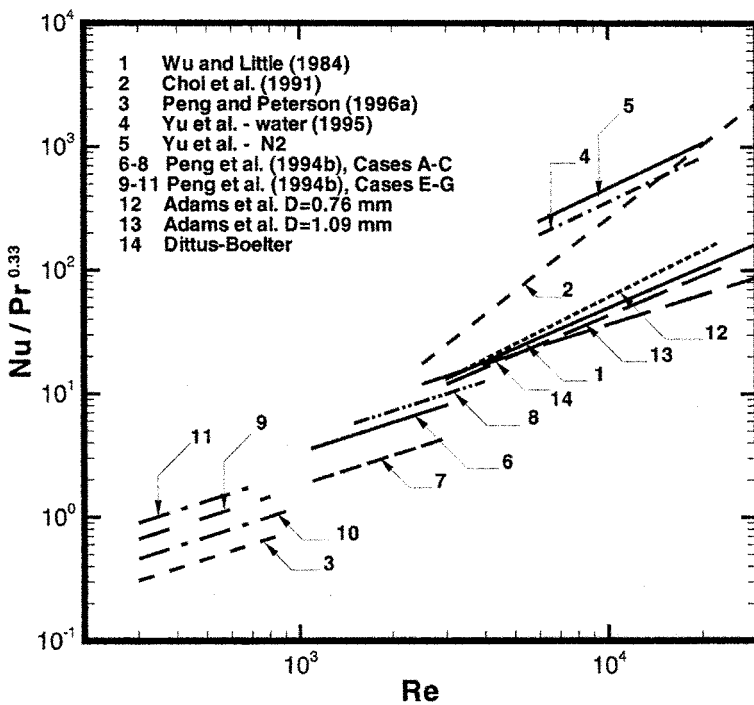


Figure 6. Heat transfer predictions in the turbulent regime.

with conventional correlations for larger, conventional tubes and channels in the laminar and turbulent flow regimes. A number of working fluid and substrate combinations, and shapes and configurations of the microchannels are included in this comparison. Little agreement is seen between the predictions of different investigators. The results are also not seen to be distinguished by fluid or substrate type or by microchannel dimensions and shapes.

The comparative study presented here points to differences between the flow and heat transfer in microchannels and that in channels of conventional sizes. However, the information in the literature thus far does not point to unequivocal trends of variation or reasons for such trends. There is no evidence that continuum assumptions are violated for the microchannels tested, most of which have hydraulic diameters of 50 μm or more. As such, analyses based on Navier-Stokes and energy equations would be expected to model the phenomena observed adequately, as long as the experimental conditions and measurements are identified and simulated correctly. The discrepancies in predictions may very well be due to entrance and exit effects, differences in surface roughness in the different microchannels investigated, nonuniformity of channel dimensions, nature of the thermal and flow boundary conditions, and uncertainties and errors in instrumentation, measurement, and measurement locations. Given the diversity in the results in the literature, a reliable prediction of the heat transfer rates and pressure drops in microchannels is not currently possible for design applications such as microchannel heat sinks. There is a clear need for additional systematic studies which carefully consider each parameter influencing transport in microchannels.

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