

SSC physical parameters from observations

(1)

Given the synchrotron peak energy ϵ_s , SSC peak energy ϵ_c (both in the observer's frame), the synchrotron peak luminosity L_s and the Compton dominance k of a source, find its physical configuration.

Recall: $\epsilon_s = \gamma_b^2 \frac{B\delta}{B_{crit}}$, $\epsilon_c = \gamma_b^4 \frac{B\delta}{B_{crit}}$, B in $B_{crit} = 4.41 \cdot 10^3$ G units

This is the slow cooling regime

Here I assumed that the electron index is $p > 2$ (neglect γ^p) and the peaks of the SED are formed at γ_b where cooling and escape times are equal. I also assumed Thomson scattering forms the SSC peak.

I will include numerical values as an example for the case of aligned HSP sources that are powerful:

$\epsilon_s = 2 \times 10^{-4}$, $\epsilon_c = 2 \times 10^5$, $L_s = 10^{46}$ erg/s, $k=1$

Find γ_b : $\gamma_b = \left(\frac{\epsilon_c}{\epsilon_s}\right)^{1/2} = 10^{4.5}$

Thompson condition requires $\frac{1}{\delta} \epsilon_s \gamma_b < 1 \Rightarrow \delta > \epsilon_s \left(\frac{\epsilon_c}{\epsilon_s}\right)^{1/2} = (\epsilon_s \epsilon_c)^{1/2} = 6.3$

From $\epsilon_s = \gamma_b^2 \frac{B\delta}{B_{crit}} = \frac{\epsilon_c}{\epsilon_s} \frac{B\delta}{B_{crit}} \Rightarrow \boxed{B\delta = \frac{\epsilon_s^2}{\epsilon_c} B_{crit} = 8.82 = C_1}$
In Gauss

Also $v_s = \frac{L_s}{4\pi c R^2 \delta^4} = k v_B = k \frac{B^2}{2 \cdot 8\pi} \Rightarrow B^2 R^2 \delta^4 = \frac{2}{k} \frac{L_s}{c} \Rightarrow$

$\boxed{BR\delta^2 = \left(\frac{2L_s}{kc}\right)^{1/2} = \frac{8.16 \cdot 10^{17}}{k^{1/2}} = C_2}$

Finally, (assuming escape time $\frac{\lambda R}{c}$)

$\gamma_b = \frac{3mc^2}{46\pi \frac{B^2}{8\pi} (1+k) R \lambda} = \frac{2.32 \cdot 10^{19}}{B^2 R (1+k)} \Rightarrow \boxed{BR = \frac{2.32 \cdot 10^{19}}{\lambda (1+k) \gamma_b} = \frac{7.34 \cdot 10^{14}}{(1+k) \lambda} = C_3}$

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We now have three equations with three unknowns

$$\left. \begin{aligned} B\delta &= C_1 \\ BR\delta^2 &= C_2 \\ B^2R &= C_3 \end{aligned} \right\} \begin{aligned} \delta &= \frac{C_1}{B} \\ BR \frac{C_1^2}{B^2} &= C_2 \Rightarrow \frac{R}{B} = \frac{C_2}{C_1^2} \\ B^2R &= C_3 \end{aligned} \left. \begin{aligned} B^3 &= \frac{C_3 C_1^2}{C_2} = \frac{7.34 \cdot 10^{14} \cdot 8.82^2 \sqrt{\text{K}}}{\lambda(1+z) \cdot 8.16 \cdot 10^{17}} \\ B^3 &= 7 \times 10^{-2} \frac{\sqrt{\text{K}}}{\lambda(1+z)} \end{aligned} \right\}$$

For our numerical example, $\lambda = 1$: $B = 0.33 \text{ G}$

$$\delta = \frac{C_1}{B} = 26.7 \quad R = \frac{C_3}{B^2} = 3.43 \cdot 10^{15} \text{ cm}$$

This corresponds to a variability light crossing time scale

$$\tau_{\text{var}} = \frac{R}{c\delta} = 4.28 \cdot 10^3 \text{ s} \sim 1 \text{ h}$$

If the jet has opening angle $\theta = \frac{1}{\Gamma} \cong \frac{1}{\delta}$, then ^{for} the blob to "fit" in the jet needs to be at a distance at least

$$d_{\text{min}} = R\Gamma = 6.5 \cdot 10^{16} \text{ cm}$$

The Poynting flux required is $L_B = \frac{B^2}{8\pi} \pi R^2 c \Gamma^2 = 3.42 \cdot 10^{42} \frac{\text{erg}}{\text{s}}$

The radiated power is $L_{\text{rad}} = L'_{\text{rad}} \Gamma^2 = \frac{L_S(\text{CHX})}{\delta^4} \Gamma^2 = \frac{L_S(\text{CHX})}{\delta^2 \gamma_b^2} = 1.4 \cdot 10^{43} \frac{\text{erg}}{\text{s}}$

The minimum relativistic electron power is $L_{e,\text{min}} = L_{\text{rad}}$

So the source is particle dominated at least by $\frac{L_{e,\text{min}}}{L_B} = 4.1$

NOTE: For $p > 2$ $L'_{\text{ins}} = \frac{Qmc^2}{p-2} \gamma_{\text{min}}^{2-p}$, $L'_{\text{rad}} = \frac{Qmc^2}{p-2} \gamma_b^{2-p}$

$$L'_{\text{ins}} = L'_{\text{rad}} \left(\frac{\gamma_{\text{min}}}{\gamma_b} \right)^{2-p} = \frac{L_S(\text{CHX})}{\delta^4} \left(\frac{\gamma_{\text{min}}}{\gamma_b} \right)^{2-p} \cdot L_{\text{ins}} = L'_{\text{ins}} \Gamma^2 = \frac{L_S(\text{CHX})}{\delta^4} \Gamma^2 \left(\frac{\gamma_{\text{min}}}{\gamma_b} \right)^{2-p}$$

Scaling and Paper form for B:

$$B^3 = \frac{C_3 C_1^2}{C_2} = \frac{C'}{1+K} \left(\frac{E_s}{E_c} \right)^{1/2} \frac{E_s^4}{E_c^2} B_{cr}^2 \frac{K^{1/2} C^{1/2}}{2^{1/2} L_s^{1/2}}$$

$$B = \text{const}^{1/3} \frac{E_s^{3/2}}{E_c^{5/6} L_s^{1/6}} \frac{K^{1/6}}{(1+K)^{1/3}}$$

$$\text{const} = \frac{C' C^{1/2} B_{cr}^2}{2^{1/2}}$$

$$C' = \frac{6 \pi m c^2}{6 \tau}$$

Substituting $E_s = \frac{h 10^{17}}{m c^2} \nu^{1/7}$, $E_c = \frac{10^{11}}{5.11 \cdot 10^5} E_{100}$, $L_s = (10^{45})^{1/5} L_{45}$
↑
100 GeV

We obtain $B = 5.01 \frac{\nu^{3/2}}{E_{100}^{5/6} L_{45}^{1/6}} \frac{K^{1/6}}{(1+K)^{1/3}} G$

Scaling and Paper form for δ

$$\delta = \frac{a}{B} = \frac{E_s^2}{E_c} B_{cr} \left[\frac{\nu^2}{C' C^{1/2} B_{cr}^2} \right]^{1/3} \frac{E_c^{5/6} L_s^{1/6}}{E_s^{3/2}} \frac{(1+K)^{1/3}}{K^{1/6}}$$

$$\delta = \left[\frac{\nu^2 B_{cr}}{C' C^{1/2}} \right]^{1/3} \frac{E_s^{1/2}}{E_c^{1/6}} \frac{1/6}{L_s} \frac{(1+K)^{1/3}}{K^{1/6}}$$

$$\delta = 29.42 \frac{\nu^{1/2}}{E_{100}^{1/6}} \frac{1/6}{L_{45}} \frac{(1+K)^{1/3}}{K^{1/6}}$$

Scaling and paper form for R

$$R = \frac{C_3}{B^2} = \frac{C'}{1+K} \frac{1}{\text{Conse}^{2/3}} \frac{\epsilon_s^{1/6} \epsilon_c^{10/6} L_s^{1/3}}{\epsilon_s^{3/2}} \frac{(L+K)^{2/3}}{K^{1/2}}$$

$$R = \frac{C'}{\text{Conse}^{2/3}} \frac{\epsilon_c^{7/6} L_s^{1/3}}{\epsilon_s^{5/2}} \frac{1}{(K(K+1))^{1/3}}$$

$$R = 5.95 \cdot 10^{13} \frac{\epsilon_{100}^{7/6} L_{45}^{1/3}}{V_{17}^{5/2}} \frac{1}{(K(K+1))^{1/3}} \text{ cm}$$

Finally

$$t_{\text{var}} = 67.4 \frac{\epsilon_{100}^{4/3} L_{45}^{1/6}}{V_{17}^3} \frac{1}{K^{1/6} (1+K)^{2/3}} \text{ s}$$

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Assuming a jet power of $10^{44.5}$ from the cavity work,

the Poynting flux is $\frac{3.42 \cdot 10^{42}}{10^{44.5}} = 10^{-2}$ of the jet power

the radiative efficiency is $\frac{L_{\text{rad}}}{L_{\text{jet}}} = \frac{1.4 \cdot 10^{43}}{10^{44.5}} = 4.4 \times 10^{-2}$

To obtain the lowest γ_{min} to which the injected power-law can extend we equate $L_{\text{inj}} = L_{\text{jet}} = 10^{44.5}$

$$L_{\text{jet}} = \frac{L_s}{\delta^4} r^2 \left(\frac{\gamma_{\text{min}}}{\gamma_b} \right)^{2-p} \Rightarrow \left(\frac{\gamma_{\text{min}}}{\gamma_b} \right)^{2-p} = \frac{L_{\text{jet}} \delta^2}{L_s}$$

$$\gamma_{\text{min}} = \gamma_b \left(\frac{L_{\text{jet}} \delta^2}{L_s} \right)^{\frac{1}{2-p}} \quad \text{for } p=2.5$$

$$\gamma_{\text{min}} = 10^{4.5} \left(\frac{10^{44.5} \cdot 26.7^2}{10^{46}} \right)^{-2} = 62$$

This corresponds to $v_{\text{min}} = 2 \times 10^6 B \gamma^2 r \approx 2 \times 10^{12} \text{ Hz}$