



Fermi  
Gamma-ray Space Telescope

# ENERGY DISPERSION TUTORIAL (A.K.A. IRFs DO NOT BITE)

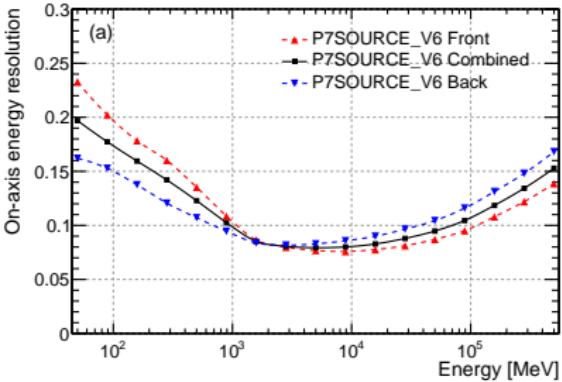
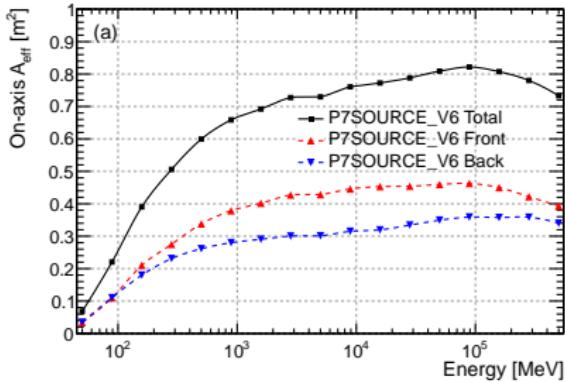
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Fermi Summer School 2012  
Lewes, 2012

# THE NAME OF THE GAME

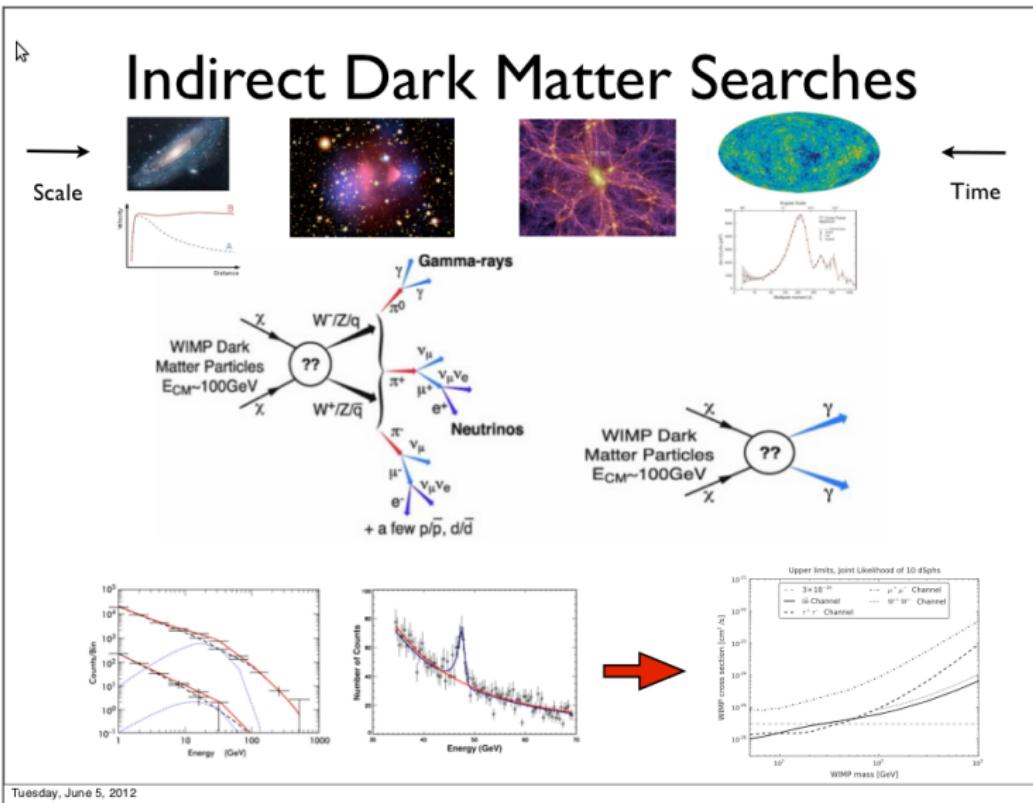
- ▶ Wouldn't it be nice if we could measure the photon energy perfectly?
- ▶ Evaluate the effect of the energy dispersion on a power-law input spectrum.
- ▶ Remember: *ScienceTools* ignore the energy dispersion by default!
- ▶ In real life this is really something you would do with *gtobbsim*:
  - ▶ all the dirty work done for you in the background;
  - ▶ you get the pointing history for free.
- ▶ This is really an excuse to dig into the IRF fits files:
  - ▶ it'll turn out to be handy if you have to create bracketing IRFs, for instance.

# BEFORE WE START . . .



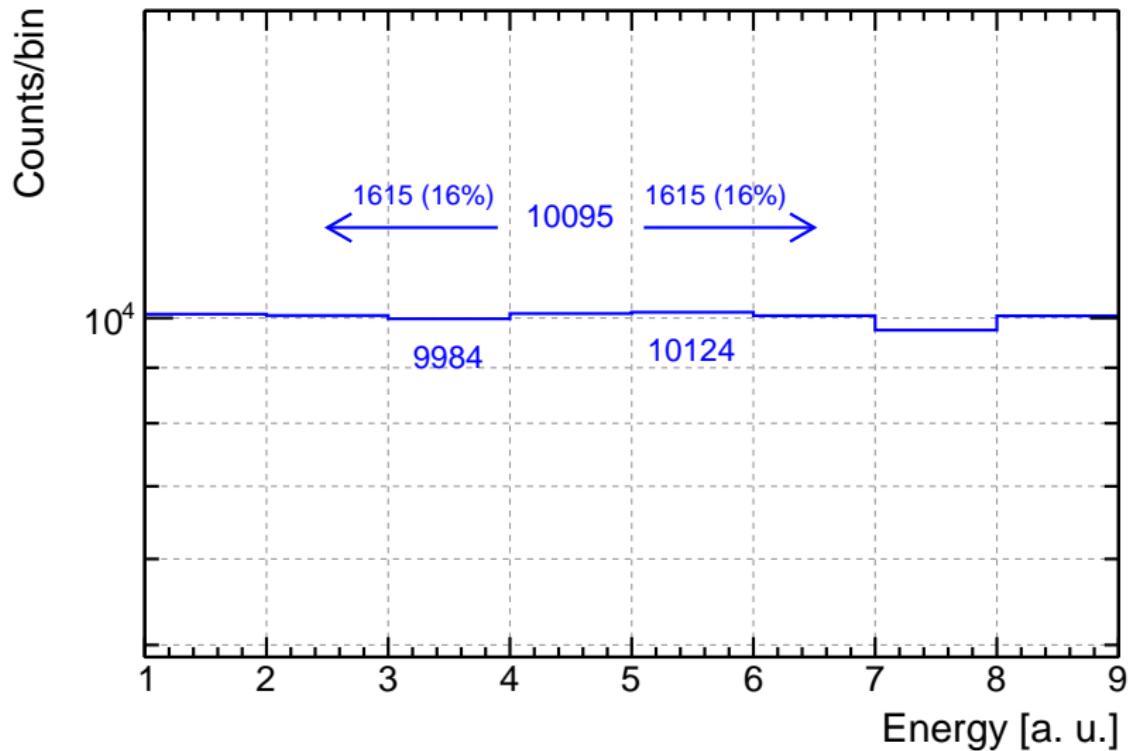
- ▶ Why the energy dispersion is (potentially) important?
  - ▶ IRFs are binned in true energy;
  - ▶ in real life you have bin-to-bin migrations in the count spectra.
- ▶ Why is the effect of the energy dispersion coupled to  $A_{\text{eff}}$ ?
- ▶ Shall we prefer low-energy or high-energy tails?
- ▶ Is the effect more pronounced for hard or steep spectra?
- ▶ Is the effect more pronounced for high or low energies?

# WHY IS ENERGY DISPERSION IMPORTANT?

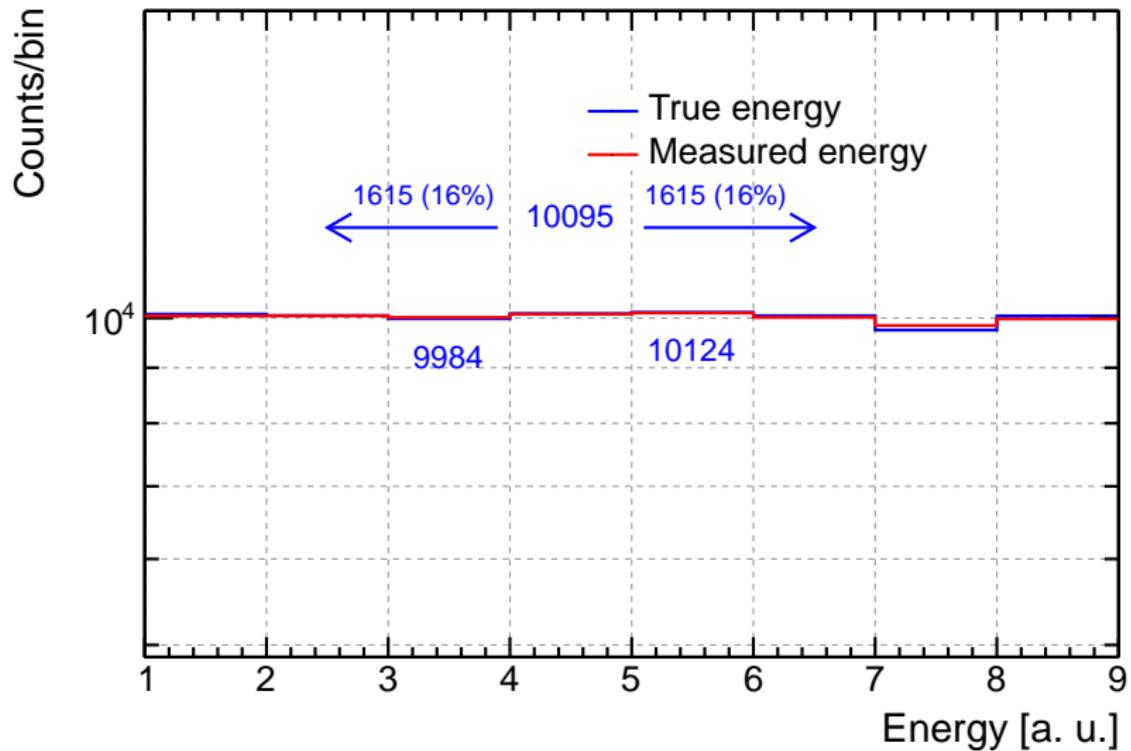


—Credits: Chris Anelli

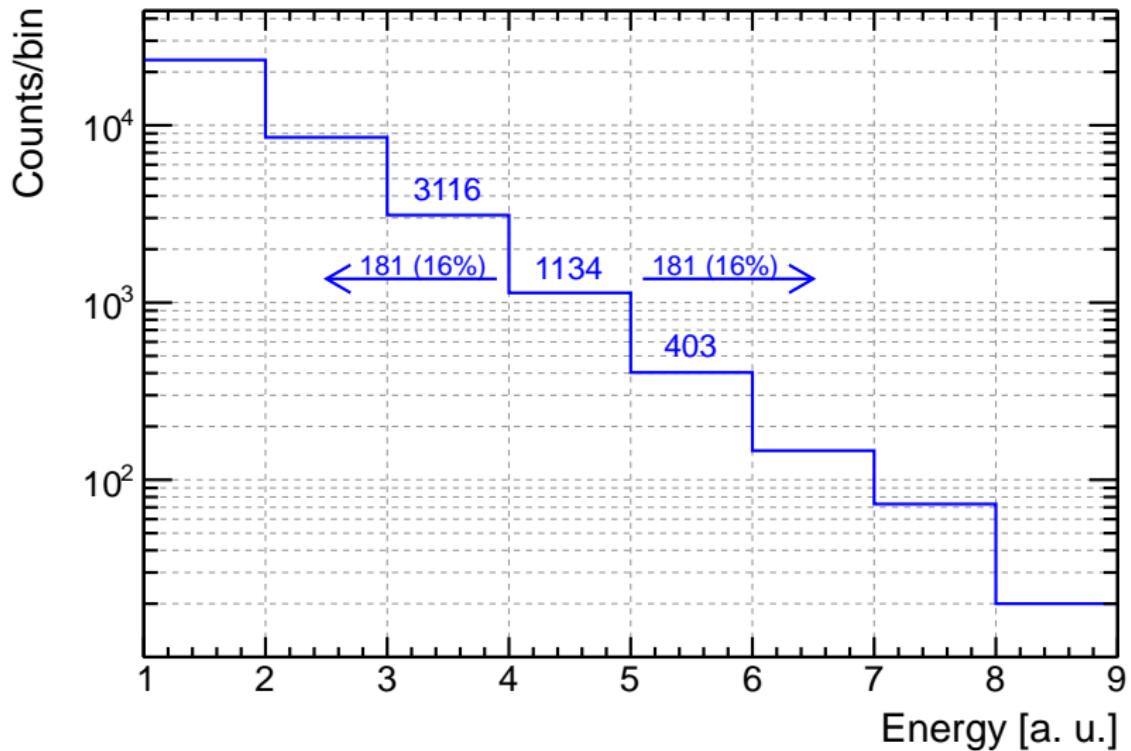
# WHAT HAPPENS WITH A BROADBAND SPECTRUM?



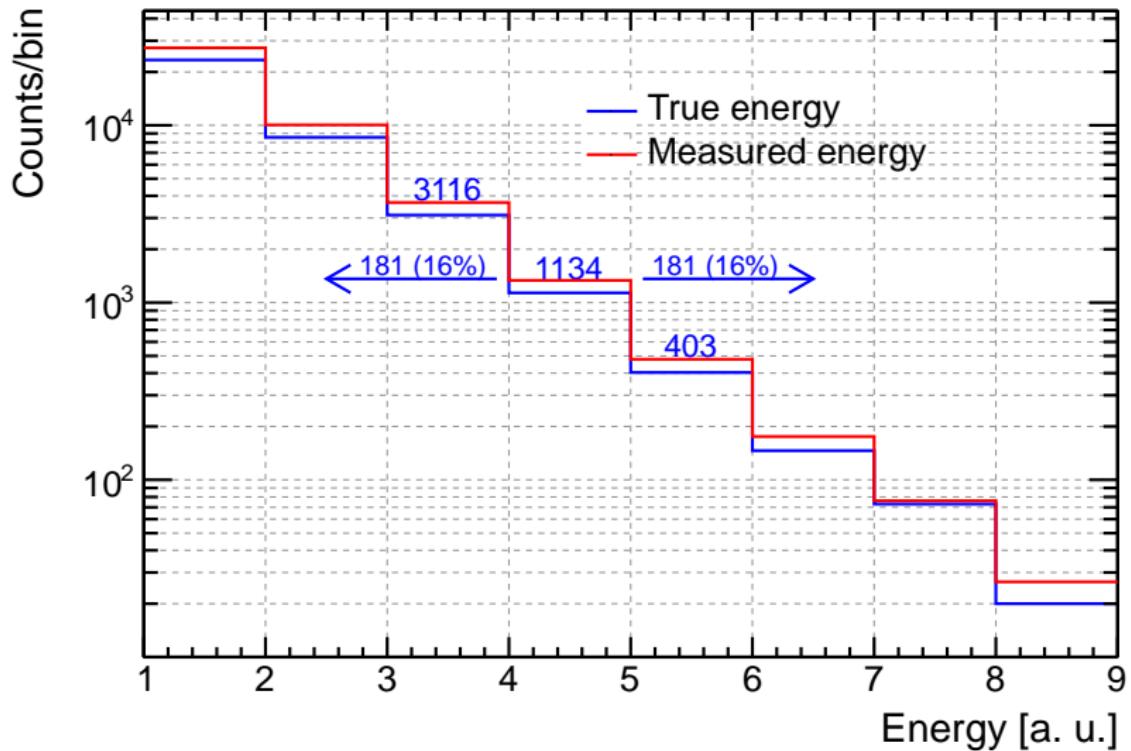
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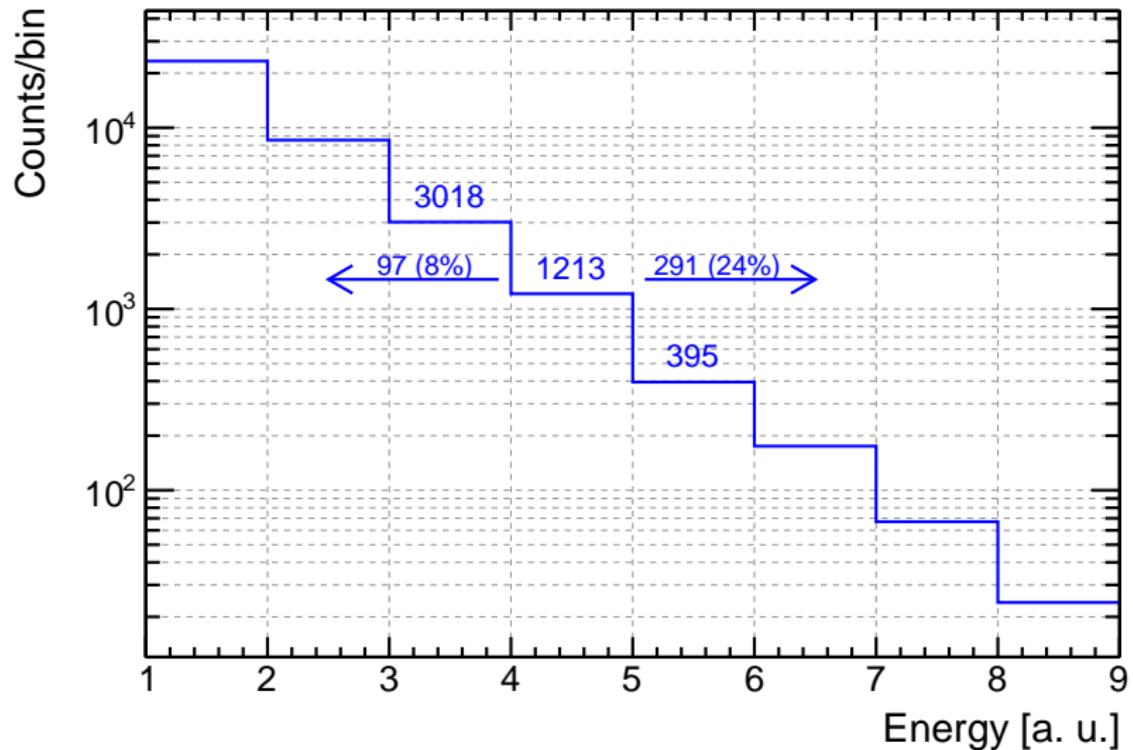
# WHAT IF THE SPECTRUM IS STEEP?



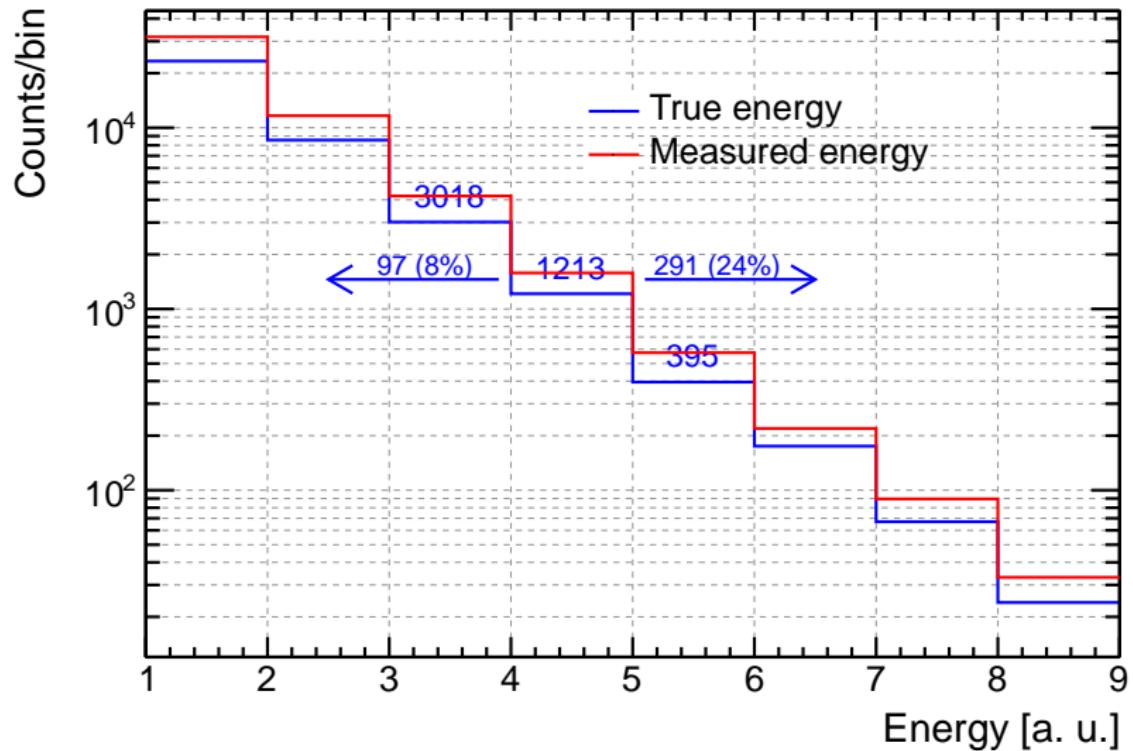
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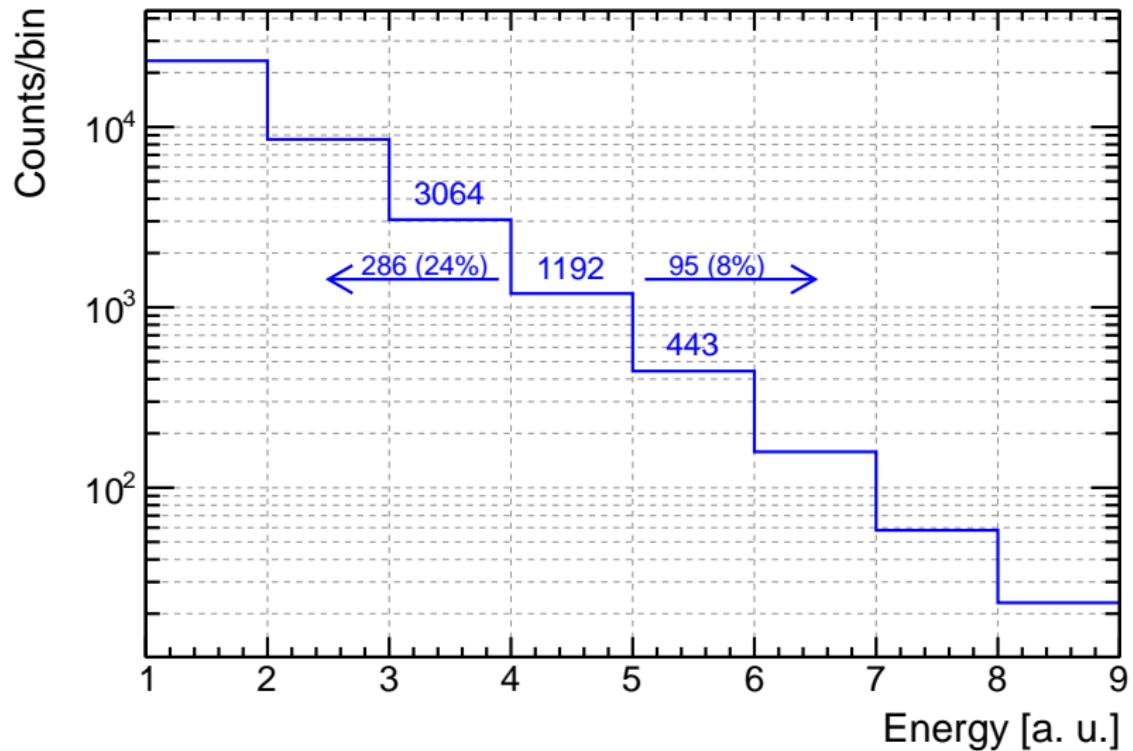
# WHAT IF THE ENERGY DISPERSION IS ASYMMETRIC?



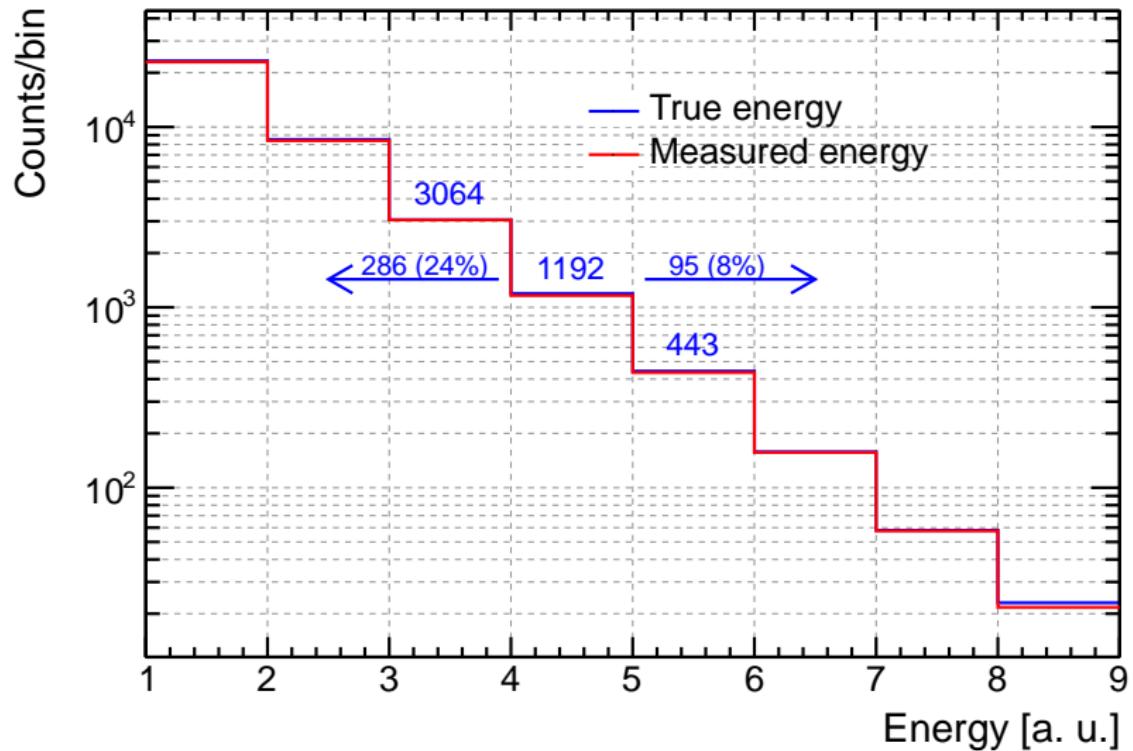
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# BACK ON TRACK

- ▶ Basic steps:

1. generate the input spectrum (energies follow a power-law distribution with  $\Gamma = 2$ , directions are isotropic in the upper hemisphere in the LAT frame);
2. fold the input spectrum with the effective area (use P7SOURCE::FRONT only);
3. smear the input spectrum with the LAT energy resolution;
4. compare the count spectra in true and measured energy.

- ▶ I'll use python as a pseudo-code, but you're welcome to use whatever you're more comfortable with!

- ▶ Some snippets of code on

<http://www-glast.stanford.edu/cgi-bin/viewcvs/users/lbaldini/fermiSummerSchool2012/macro/>  
(but you shouldn't use them).

# GENERATING THE INPUT SPECTRUM FOR LAZY PEOPLE

Let ROOT do the dirty work behind the scenes.

Generate a power law with pyroot

```
import ROOT
E_MIN = 10.
E_MAX = 1000000.
GAMMA = 2

POWER_LAW = ROOT.TF1('fpl', 'x**(-[0])', E_MIN, E_MAX)
POWER_LAW.SetParameter(0, GAMMA)

def plRoot():
    return POWER_LAW.GetRandom()
```

# GENERATING THE INPUT SPECTRUM

IF YOU'RE BRAVE—OR YOU DON'T HAPPEN TO HAVE ROOT INSTALLED

Use the inverse transform method:

$$u = \text{random}(0, 1) \quad x = F^{-1}(u) \quad (1)$$

The cumulative distribution for a truncated power law is

$$F(x) = \frac{E_{\min}^{(1-\Gamma)} - x^{(1-\Gamma)}}{E_{\min}^{(1-\Gamma)} - E_{\max}^{(1-\Gamma)}}, \quad (2)$$

(verify it!) therefore you can use:

$$x = \left[ E_{\min}^{(1-\Gamma)} - u(E_{\min}^{(1-\Gamma)} - E_{\max}^{(1-\Gamma)}) \right]^{\frac{1}{(1-\Gamma)}} \quad (3)$$

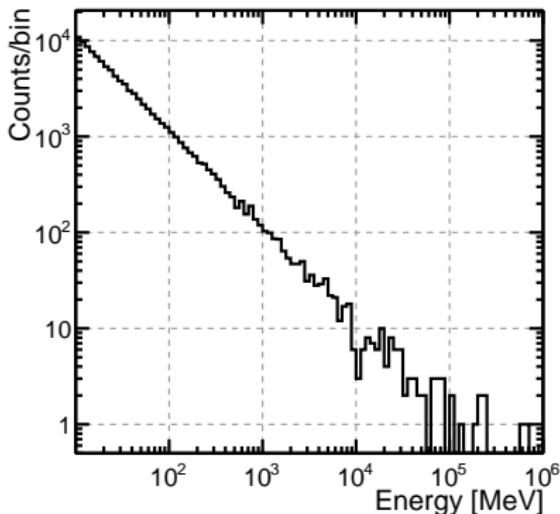
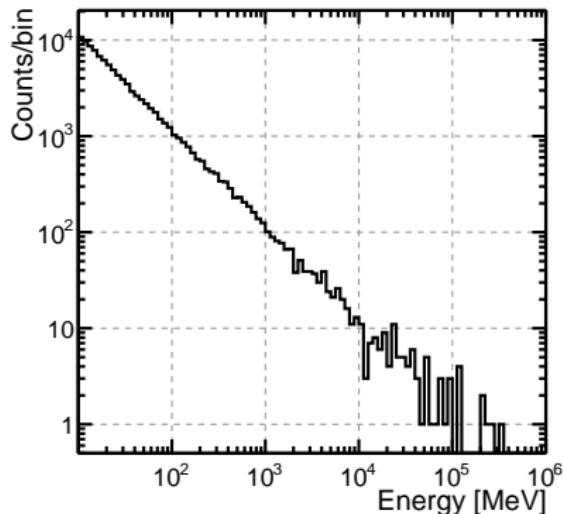
Generate a power law by the inverse transform method

```
import random
E_MIN = 10.
E_MAX = 1000000.
GAMMA = 2

def plPlain():
    u = random.random()
    x = (E_MIN**((1-GAMMA))-u*(E_MIN**((1-GAMMA))-E_MAX**((1-GAMMA))))**((1.)/(1-GAMMA))
    return x
```

# TESTING THE INPUT SPECTRUM

You should get something along these lines:



# GENERATING RANDOM DIRECTIONS

Standard recipe for a random point on the sphere:

$$\begin{aligned} u &= \text{random}(0, 1) \\ v &= \text{random}(0, 1) \\ \begin{cases} \phi = 2\pi u \\ \theta = \cos^{-1}(2v - 1) \end{cases} \end{aligned} \tag{4}$$

Here we don't bother about  $\phi$ ; also, the IRFs are tabulated in  $\cos \theta$ ; finally we are only interested in the upper hemisphere, therefore all you need is:

$$\cos \theta = \text{random}(0.2, 1) \tag{5}$$

(why 0.2?)

# IRFs: THE $A_{\text{eff}}$ TABLES

A quick look at an aeff fits file...

```
import pyfits

hdulist = pyfits.open(AEFF_FILE_PATH)
print hdulist.info()
```

...and the output should be:

```
Filename: /glast/builds/redhat5-i686-32bit-gcc41-Optimized/ScienceTools/...
No.    Name          Type      Cards   Dimensions   Format
0     PRIMARY       PrimaryHDU     10      ()           uint8
1     EFFECTIVE AREA BinTableHDU    58      1R x 5C   [60E, 60E, 32E, 32E, 1920E]
2     PHI_DEPENDENCE BinTableHDU   61      1R x 6C   [18E, 18E, 8E, 8E, 144E, 144E]
3     EFFICIENCY_PARAMS BinTableHDU  38      4R x 1C   [6E]
```

- ▶ Three main pieces:
  1. effective area table;
  2.  $\phi$  dependence correction;
  3. livetime correction.
- ▶ Effectively we only care about 1.
- ▶ And you should have your file in

`$INST_DIR/irfs/caldb_release/CALDB/data/glast/lat/bcf/ea/aeff_P7SOURCE_V6_front.fits`

# IRFs: THE $A_{\text{eff}}$ TABLES

A quick look at the effective area table...

```
import pyfits

hdulist = pyfits.open(AEFF_FILE_PATH)
hdu = hdulist['EFFECTIVE AREA']
print 'Header data unit %s' % hdu.name
for (i, column) in enumerate(hdu.columns):
    print '%3d\t%20s\t[%6s]\t%s\n' %\
        (i, column.name, column.unit, column.format)
```

...and the output should be:

```
Header data unit EFFECTIVE AREA
 0          ENERG_LO  [      Mev] 60E
 1          ENERG_HI  [      Mev] 60E
 2          CTHETA_LO  [      ] 32E
 3          CTHETA_HI  [      ] 32E
 4          EFFAREA   [      m2] 1920E
```

► So we get:

- the energy binning (60 bins);
- the  $\cos \theta$  binning (32 bins);
- the actual effective area values ( $60 \times 32 = 1920$  values).

# IRFs: THE $A_{\text{eff}}$ TABLES

A quick look at the effective area table values...

```
import pyfits

hdulist = pyfits.open(AEFF_FILE_PATH)
data = hdulist['EFFECTIVE AREA'].data[0]
print data.field('ENERG_LO')
print data.field('ENERG_HI')
print data.field('CTHETA_LO')
print data.field('CTHETA_HI')
print data.field('EFFAREA')
# Retrieve the value of the i-th energy bin and the j-th cos theta bin.
# (You might or not need the ``flatten()``, here).
nebins = 60
i = 10
j = 22
print data.field('EFFAREA').flatten()[i + nebins*j]
```

The output is too long but you can imagine how it looks.

# IRFS: PARSING THE $A_{\text{eff}}$ TABLE...

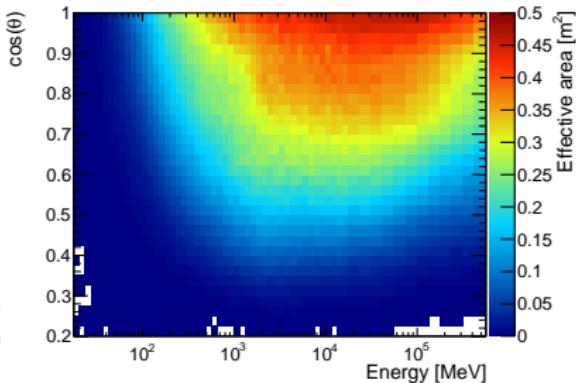
... AND STICKING IT INTO A ROOT HISTOGRAM

Read an effective area table and create a TH2F

```
# Open the file and retrieve the data.  
hdulist = pyfits.open(AEFF_FILE_PATH)  
data = hdulist['EFFECTIVE AREA'].data[0]  
  
# Create the arrays for the histogram binning.  
xbins = numpy.hstack((data.field('ENERG_LO'), data.field('ENERG_HI')[-1]))  
nx = len(xbins) - 1  
ybins = numpy.hstack((data.field('CTHETA_LO'), data.field('CTHETA_HI')[-1]))  
ny = len(ybins) - 1  
  
# Create and fill the histogram.  
h = ROOT.TH2F('aeff', 'Effective area', nx, xbins, ny, ybins)  
for i in xrange(nx):  
    for j in xrange(ny):  
        aeff = data.field('EFFAREA').flatten()[i + nx*j]  
        h.SetBinContent(i + 1, j + 1, aeff)
```

- You can use your favorite language/framework/data structure to store the  $A_{\text{eff}}$  values (and retrieve them!).

# TESTING THE $A_{\text{eff}}$ REPRESENTATION



- ▶ Your effective area table should look more or less like this.
  - ▶ Note that the bin width is not constant in  $\log(E)$ —in case you are tempted to find the bin corresponding to a given  $E$  with one division;
  - ▶ If you use ROOT you can rely on the `ROOT::TH2F::Interpolate()` method.

# ENERGY DISPERSION RECAP

First we define the *scaled energy deviation*:

$$x = \frac{1}{S_D(E, \theta)} \frac{(E' - E)}{E} \quad (6)$$

where

$$\begin{aligned} S_D(E, \theta) = & c_0 (\log_{10} E)^2 + c_1 (\cos \theta)^2 + c_2 \log_{10} E + c_3 \cos \theta + \\ & + c_4 \log_{10} E \cos \theta + c_5. \end{aligned} \quad (7)$$

Then, in each energy/angle bin we fit the distribution of the scaled deviation with four piecewise Rando functions:

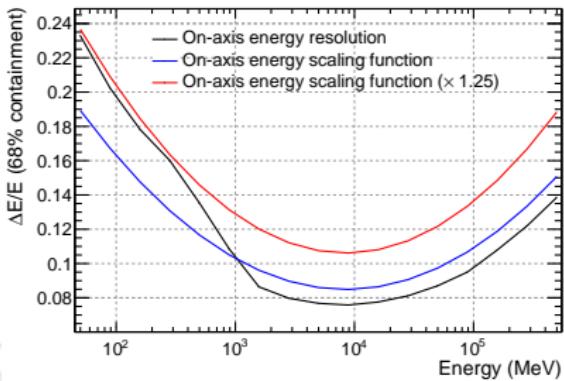
$$R(x, x_0, \sigma, \gamma) = N \exp \left( -\frac{1}{2} \left| \frac{x - x_0}{\sigma} \right|^{\gamma} \right). \quad (8)$$

All the parameters are stored in the fits files, so in principle you can reverse-engineer that (and you're very welcome to do so).

# A CHEAP PROXY FOR THE ENERGY DISPERSION

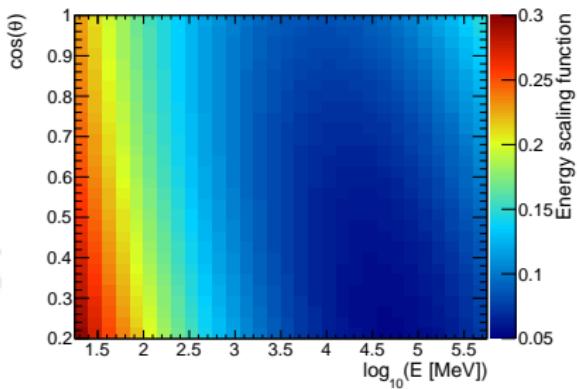
- ▶ Remember: the scaling function is defined with the goal of making the energy dispersion as independent as possible from energy and angle:
  - ▶ if the energy dispersion was Gaussian...
  - ▶ ... and the prescaling function captured all the richness of the energy dispersion itself...
  - ▶ then we wouldn't need to perform the fitting step: the scaled energy deviation would be  $\text{Norm}(0, 1)$  and the value of the scaling function would be the energy resolution.
- ▶ For the purpose of this exercise (at least to start with) we'll cheat (twice):
  - ▶ take the energy scaling function as a proxy for the energy resolution;
  - ▶ assume the energy dispersion is Gaussian.
- ▶ The parameters of the energy scaling function (7) are stored in the edisp fit file as the first 6 numbers of the 'EDISP\_SCALING\_PARAMS' HDU.

# HOW MUCH OF A CHEAT IS CHEATING?



- ▶ A little bit, but not too much;
  - ▶ the general trend is well reproduced;
  - ▶ we do underestimate the energy resolution at low energy;
  - ▶ (this is actually the most interesting part, so you might as well multiply the scaling function by  $\sim 1.25$ );
  - ▶ and remember we are underestimating the tails (you might even multiply by  $\sim 1.5$ ).

# TESTING THE EDISP REPRESENTATION



- Your energy scaling function should look more or less like this.

# NOW WE'RE READY TO GO!

poor man's attempt at putting everything together

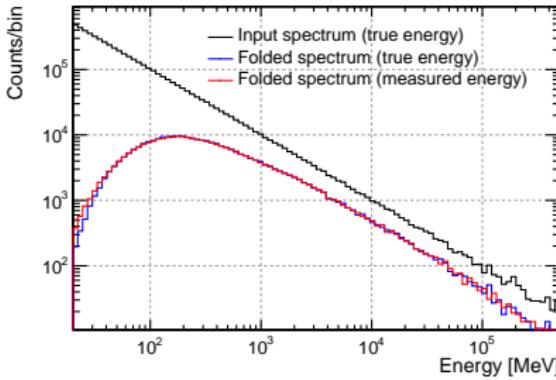
```
import random

MAX_EFF_AREA = 0.5 # put the maximum of the effective area, here.

for i in xrange(NUM_EVENTS):
    # Extract energy and direction
    energy = ... # use your function, here
    costheta = random.uniform(0.2, 1)
    # Convolve with the effective area.
    aeff = ... # retrieve your effective area at the right energy/angle.
    ... # do your book-keeping if needed.
    if random.random() < aeff/MAX_EFF_AREA:
        eres = ... # get the value of the scaling function.
        measuredEnergy = energy*random.gauss(1, 1.5*eres)
        ... # more book-keeping.
```

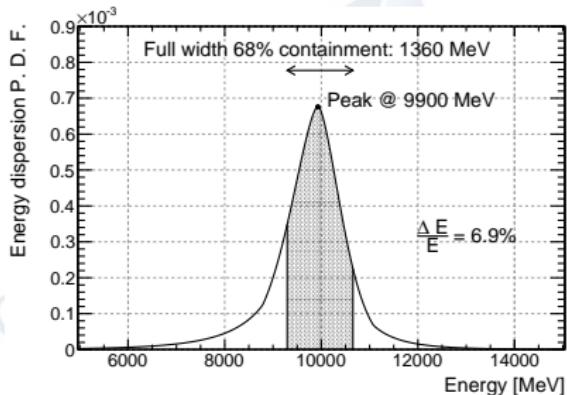
Fill in the blanks...

# AND FINALLY...



- ▶ The difference between blue and red is what you're after:
  - ▶ plot the ratio of the two histograms.
- ▶ You can change the input spectral index!
  - ▶ Things get better or worst with harder spectra? Why?
- ▶ And you can change the energy scale too!
  - ▶ Can you explain analytically what happens above a few GeV?

# ONE LAST THING



- ▶ It would be totally awesome if you plotted the actual energy dispersion, say at 10 GeV on axis;
  - ▶ (or at your favorite energy/angle)!

# SOME DIRECTIONS...

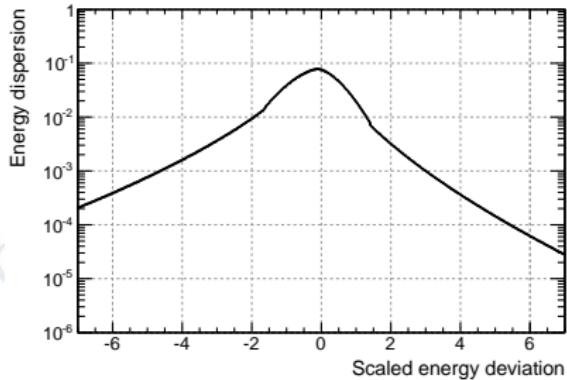
$$R(x, x_0, \sigma, \gamma) = N \exp \left( -\frac{1}{2} \left| \frac{x - x_0}{\sigma} \right|^{\gamma} \right) \quad (9)$$

$$D(x) = \begin{cases} N_L R(x, x_0, \sigma_L, \gamma_L) & \text{if } (x - x_0) < -\tilde{x} \\ N_I R(x, x_0, \sigma_I, \gamma_I) & \text{if } (x - x_0) \in [-\tilde{x}, 0] \\ N_r R(x, x_0, \sigma_r, \gamma_r) & \text{if } (x - x_0) \in [0, \tilde{x}] \\ N_R R(x, x_0, \sigma_R, \gamma_R) & \text{if } (x - x_0) > \tilde{x} \end{cases} \quad (10)$$

$\tilde{x}$	$\gamma_L$	$\gamma_I$	$\gamma_r$	$\gamma_R$
1.5	0.6	1.6	1.6	0.6

Find in the FITS files: the normalization  $N_r = N_I$  (NORM), the centroid position  $x_0$  (BIAS), the two core scales  $\sigma_r$  (RS1) and  $\sigma_I$  (LS1) and the two tail scales  $\sigma_R$  (RS2) and  $\sigma_L$  (LS2).

# SOME DIRECTIONS...



- And then you have to un-prescale...